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A NEW GENERATION OF SIMULTANEOUS FITS TO LHC DATA USING DEEP LEARNING

DIS2022 - SANTIAGO DE COMPOSTELA

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• Introduction:

Simultaneous fits of PDFs and physics parameters at the LHC

- The SimuNET methodology
 - (based on S. Iranipour, MU, arXiv:2201.07240)
 - ➡ The approach
 - ➡ Application to the determination of PDFs and SMEFT coefficients

Conclusions and outlook

INTRODUCTION

EXTRACTING PHYSICS PARAMETERS FROM LHC DATA

✓ Abundance of precise LHC data allows to extract information on SM and BSM parameters & nonperturbative objects such as PDFs with unprecedented precision



$$\chi^{2} = \frac{1}{N_{\text{dat}}} \sum_{i=1}^{N_{\text{dat}}} (T_{i}(\{\theta\}, \{c\}) - D_{i}) \operatorname{cov}_{ij}^{-1} (T_{j}(\{\theta\}, \{c\}) - D_{j}))$$

$$T_{i}(\{\theta\}, \{c\}) = PDFs(\{\theta\}, \{c\}) \otimes \hat{\sigma}_{i}(\{c\})$$
(B)SM parameters: $\boldsymbol{\alpha}_{s}(M_{z}), M_{w}, \theta_{w}$, SMEFT WCs....

Parameters determining PDFs at initial scale

✓ In a PDF fit typically

$$T_i(\{\theta\}) = \text{PDFs}(\{\theta\}, \{c = \bar{c}\}) \otimes \hat{\sigma}_i(\{c = \bar{c}\})$$

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or - for example in case of $\pmb{\alpha}_{s}(\mathsf{M}_{z)}$ -

 $T_i(\{c\}) = \text{PDFs}(\{\bar{\theta}\}, \{c\}) \otimes \hat{\sigma}_i(\{c\})$

SIMULTANEOUS FITS FOR SM PARAMETERS

Given the strong correlation between PDFs of the proton and α_s , a non simultaneous determination of α_s along with the PDFs from LHC processes might yield misleading results





PDG 2017

80

82

84

86

88

m_w [GeV]



- Correlation of PDFs and the EW parameters or m_t weaker than in the case of α_s , but the very high accuracy which is sought suggests that the effect of simultaneous determination is not negligible
- Similar considerations for fits of polarised/ unpolarised PDFs, proton/nuclear PDFs or PDFs and FFs (universal fits)

SIMULTANEOUS FITS FOR PDFS AND SMEFT

- Interplay between PDFs and new physics not negligible [Carrazza et al, 1905.05215][Greljo et al, 2104.02723] [Liu, Sun, Gao, 2201.06586] [CMS collaboration 2111.10431]
- Crucial to assess to what degree new physics might be absorbed in PDFs and how PDF treatment affects new physics bounds
- DY @ Run I+II: effect visible but within PDF uncertainties.
- DY @ HL-LHC not accounting for interplay leads to overconstrained bound [Greljo et al, 2104.02723]
- CMS analysis on inclusive jet cross section points to a nonnegligible interplay [CMS collaboration, 2111.10431]



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These results point towards the need of new generation of global fits, in which all ingredients that enter theoretical predictions are treated consistently.

Greljo et al, 2104.02723

THE SIMUNET METHODOLOGY

SIMUNET: A DEEP-LEARNING BASED SIMULTANEOUS FIT

- The idea: take a PDF fit based on NNPDF4.0 methodology [R. Stegeman's talk] and make dependence of observables on physics parameters {c_i} explicit via fast interface before computing the loss function (e.g. adding SMEFT corrections, expanding observables in terms of SM precision parameters)
- Perform minimisation of loss function over
 - $\hat{\theta} = \theta \bigcup \{c_i\}$

by adding new layer to the deep neural network used in NNPDF4.0

 Can expand dependence on c_i beyond linear terms in T (up to generic power in polynomial expansion) by adding nontrainable edges



- Case study at higher energy: EW oblique corrections in high-mass NC and CC Drell-Yan tails.
- W and Y parametrise the self-energy of gauge bosons and are powerful probes of quark-lepton contact interactions that produce effects that grow with energy [Torre et al, 2008.12978]





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- We performed a similar analysis as in Torre et al, now with emphasis on PDF and their interplay with bounds on oblique operators [Greljo, Iranipour, Kassabov, Madigan, Moore, Rojo, MU, Voisey: 2104.02723]
 [Iranipour, MU, Voisey: 2201.07240]
- Settings:
 - PDF fit based on DIS (~3000 data points), Drell-Yan on-shell and low-mass data from ATLAS, CMS and LHCb (~600 data points)
 - + Run I and II ATLAS and CMS high mass NC Drell-Yan data (~300 data points)
 - ➡ SM predictions at NNLO QCD + NLO EW

Exp.	\sqrt{s} (TeV)	Ref.	$\mathcal{L}~(\mathrm{fb}^{-1})$	Channel	$1\mathrm{D}/2\mathrm{D}$	$n_{ m dat}$	$m_{\ell\ell}^{ m max}~({ m TeV})$
ATLAS	7	[117]	4.9	e^-e^+	1D	13	[1.0,1.5]
ATLAS (*)	8	[83]	20.3	$\ell^-\ell^+$	2D	46	[0.5,1.5]
CMS	7	[118]	9.3	$\mu^-\mu^+$	$2\mathrm{D}$	127	[0.2, 1.5]
CMS (*)	8	[84]	19.7	$\ell^-\ell^+$	1D	41	[1.5,2.0]
CMS (*)	13	[119]	5.1	$e^-e^+,\mu^-\mu^+ \ \ell^-\ell^+$	1D	$\begin{array}{c} 43,\ 43\\ 43\end{array}$	[1.5, 3.0]
Total						270 (313)	





Greljo et al, arXiv:2104.02723

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 - ➡ SM predictions at NNLO QCD + NLO EW
 - SMEFT corrections added via local K-factors





• In the previous analysis, a scan in the (W,Y) parameter space was made [Greljo, Iranipour, Kassabov, Madigan, Moore, Rojo, MU, Voisey: 2104.02723]



 $T = f_1(\hat{W}) \otimes f_2(\hat{W}) \otimes \hat{\sigma}(\hat{W})$

THE SIMUNET ANALYSIS

• With SimuNET we can do a truly simultaneous fit, rather than a scan in benchmark point and it does not have limit in number of parameters that can be fitted alongside PDFs at the initial scale!

[Iranipour, MU, Voisey: 2201.07240]



Linear dim-6 operator

$$T(\hat{\theta}) = \Sigma(\{c_n\}) \cdot L^0(\theta) = T^{\text{SM}}(\theta) \cdot \left(1 + \sum_{n=1}^N c_n R_{\text{SMEFT}}^{(n)}\right)$$

$$T^{\rm SM}(\theta) = \Sigma^{\rm SM} \cdot L^0(\theta)$$

S. Iranipour, MU - arXiv: 2201.07240

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SMEFT

Observabl

Hidden Hidden Convolution Input PDF SMSMEFT layer 1 layer 2 Observable Observable layer flavours step $h_{1}^{(1)}$ $h_2^{(1)}$ $h_1^{(2)}$ f_3 $h_{3}^{(1)}$ $h_2^{(2)}$ f_4 $h_{3}^{(2)}$ $h_4^{(1)}$ TSM f_5 $\ln x$ $h_{5}^{(1)}$ f_6 $h_{20}^{(2)}$ $h_{25}^{(1)}$ SMObservable T^{SM} S. Iranipour, MU - arXiv: 2201.07240 $-c_{NN}$

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$$T^{\rm SM}(\theta) \,=\, \Sigma^{
m SM} \,\cdot\, L^0(\theta)$$

Quadratic dim-6 operator

$$T(\hat{\theta}) = T^{\text{SM}}(\theta) \cdot \left(1 + \sum_{n=1}^{N} c_n R_{\text{SMEFT}}^{(n)} + \sum_{1 \le n \le m \le N} c_{nm} R_{\text{SMEFT}}^{(n,m)}, \right)$$
$$C_n C_m$$

RESULTS: DRELL-YAN DATA @RUN1 AND RUN2



Distribution of W & Y best fits over MC reps with fixed SM PDFs (baseline)

Distribution of W & Y best fits over MC reps with PDFs fitted alongside them

RESULTS: DRELL-YAN DATA @RUN1 AND RUN2



Distribution of W & Y best fits over MC reps with fixed SM PDFs (baseline)

Distribution of W & Y best fits over MC reps with PDFs fitted alongside them

Same comparison for quark-antiquark luminosity

RESULTS: DRELL-YAN DATA @RUN1 AND RUN2



✓ Simultaneous analysis confirms results of previous study based on scan on benchmark points in the SMEFT space: with current data effect is not-negligible but small compared to PDF uncertainties

✓ Methodology able to find flat direction in W-Y parameter space

✓ To eliminate it, need Drell-Yan charged current data

RESULTS: DRELL-YAN DATA @HL-LHC

• Add HL-LHC projections for both NC and CC in PDF fit

$$\sigma_{i}^{\text{hllhc}} \equiv \sigma_{i}^{\text{th}} \left(1 + \lambda \delta_{\mathcal{L}}^{\text{exp}} + r_{i} \delta_{\text{tot},i}^{\text{exp}} \right) , \qquad i = 1, \dots, n_{\text{bin}}$$
$$\delta_{\text{tot},i}^{\text{exp}} \equiv \left(\left(\delta_{i}^{\text{stat}} \right)^{2} + \sum_{j=1}^{n_{\text{sys}}} \left(f_{\text{red},j} \delta_{i,j}^{\text{sys}} \right)^{2} \right)^{1/2}$$





RESULTS: DRELL-YAN DATA @HL-LHC



✓ Simultaneous analysis confirms results of previous study based on scan on benchmark points in the SMEFT space: at HL-LHC the effect of interplay becomes important as WCs bounds broaden and PDFs change significantly once SMEFT effects allowed in theory predictions entering PDF fit

 \checkmark Stress-tested and shown robustness with closure tests

- While huge progress made in determining key ingredients of theoretical predictions from the data, PDFs, α_s, SMEFT WCs coefficients, it is not yet evident how to combine all these partial fits into a global interpretation of the LHC data
- Time to wok on new generation of global fits, in which all ingredients that enter theoretical predictions are treated consistently.
- SimuNET methodology based on an extension of the NNPDF4.0 NN architecture, allows the addition of an extra layer to simultaneously determine PDFs alongside an arbitrary number of physics parameters that enter predictions.
- Proof-of-concept on simultaneous determination of W and Y oblique parameters and PDFs from DIS+DY fit
- Lots of exciting avenues being explored:
 - \blacktriangleright Determination of PDFs and $\alpha_{\rm s}$ (Stegeman et al)
 - Determination of PDFs and SMEFT coefficients in the top sector (Kassabov, Madigan, Mantani, Moore, Morales, Rojo, MU)
 - Systematic study of new physics contamination in PDF fits
 - Determination of PDFs and electroweak parameters

THANK YOU FOR YOUR ATTENTION!

EXTRA MATERIAL

FAST INTERFACE FOR THEORETICAL PREDICTIONS

 $T_I = \Sigma_I \cdot L^0$

 $\label{eq:star} \textbf{T}_{l} \text{:} \end{tabular} T_{l} \text{:} \end{tabular}$

$$\mathbf{\Sigma}_{I} = \mathbf{\Sigma}_{I} \begin{pmatrix} \alpha_{s}(M_{Z}) = 0.118\\ \sin^{2}\theta_{W}(M_{Z}) = 0.23\\ M_{W} = 80.4 \text{ GeV}/c^{2}\\ M_{Z} = 91.2 \text{ GeV}/c^{2}\\ \vdots\\ c_{i} = 0 \end{pmatrix}$$

$$\Sigma_I(c) = [\hat{\sigma}(c) \otimes \Gamma(c)]_I$$

Depending on parameter to fit alongside PDFs, its dependence might appear in either/both σ and Γ

 Need to include a fast interface to parameter dependence to the fast interface of initial scale PDFs

$$T(\hat{\theta}) = T^{\text{SM}}(\theta) \cdot \left(1 + \sum_{n=1}^{N} c_n R_{\text{SMEFT}}^{(n)} + \sum_{1 \le n \le m \le N} c_n c_m R_{\text{SMEFT}}^{(n,m)}, \right)$$
$$R_{\text{SMEFT}}^{(n)} \equiv \left(\mathcal{L}_{ij}^{\text{NNLO}} \otimes d\widehat{\sigma}_{ij,\text{SMEFT}}^{(n)} \right) / \left(\mathcal{L}_{ij}^{\text{NNLO}} \otimes d\widehat{\sigma}_{ij,\text{SM}} \right), \quad n = 1 \dots, N$$

$$R_{\rm SMEFT}^{(n,m)} \equiv \left(\mathcal{L}_{ij}^{\rm NNLO} \otimes d\widehat{\sigma}_{ij,{\rm SMEFT}}^{(n,m)} \right) / \left(\mathcal{L}_{ij}^{\rm NNLO} \otimes d\widehat{\sigma}_{ij,{\rm SM}} \right) , \quad n,m = 1 \dots, N$$

FAST INTERFACE FOR THEORETICAL PREDICTIONS

 $T_I = \Sigma_I \cdot L^0$

T_I: Theoretical prediction for experimental measurement I in a fit **\Sigma_I**: Pre-computed FK-tables (including both partonic cross section and evolution kernel from Q_0 to Q

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- 1. Take data, make theoretical predictions accounting for operator in partonic cross section with fixed SM PDFs.
- 2. Compute chi2 as a function of WCs (Wilson Coefficients)
- 3. Minimise chi2 and find best-fit and C.L.s of WCs
- 4. Extract bounds

 $T = f_{1,\text{SM}} \otimes f_{2,\text{SM}} \otimes \hat{\sigma}_{\text{BSM}}$

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INTERPLAY @ RUN I AND RUN II





- With current data, PDFs are moderately affected by inclusion of non-zero W and Y coefficients in the fit, mostly quark-antiquark luminosity within uncertainties
- Broadening of individual bounds on W and Y once SMEFT PDFs are used (i.e. PDFs that have been fitted with consistent values of W and Y) is not negligible, but still within PDF uncertainties
- If SMEFT PDFs are used in determining bounds from ATLAS search same mild broadening (larger than PDF uncertainties)



INTERPLAY @ HL-LHC

- Compare Wilson coefficients bounds from HL-LHC projections assuming SM PDFs (that include NC+CC data) to the bounds on the same Wilson coefficients obtained from a simultaneous fit of PDFs and Wilson coefficients
- Not accounting for interplay (using PDFs as a black box) leads to over-constrained bounds
- PDFs do absorb effect of new physics in this case!



$\underline{\mathsf{PDFs}\,\mathsf{AND}\,\alpha_{\mathsf{S}}}$

- → PDFs and α_s strongly correlated (PDF evolution with the scale and hard cross sections)
- Cleanest determinations of α_s from processes that do not require knowledge of the PDFs
- A determination of α_s jointly with the PDFs has advantage that it is driven by the combination of many experimental measurements from several different processes.





Ball et al, 1110.2483

- → Early determinations involve a scan over α_s and ignored PDF and α_s correlation in the fit
- → Recent simultaneous determination of PDF and α_s using correlated replica method
- → Many determination of α_s from analyses of specific LHC processes have been published recently (from tt~, Z and W production, jets)
- \Rightarrow How reliable are such partial determination of α_s ?

PDFS AND NEW PHYSICS

- In principle low-scale physics is separable from high-scale physics, BUT the complexity of the LHC environment might well intertwine them.
- PDFs are low-scale quantities extracted from experimental data at all scales, without considering any potential high-scale contamination due to new physics.
- (SM)EFT fits are performed by assuming a priori that PDFs are SM-like.



Ball et al, arXiv:2109.02653

Ethier et al, arXiv: 2105.00006

PDFS AND NEW PHYSICS

- From the point of view of PDF fits:
 - How to make sure that new physics effects are not inadvertently fitted away in a PDF fit?
- From the point of view of SMEFT fits:
 - Should I make sure I am using a clean set of PDFs in a SMEFT analysis? How to define it? Is it enough?
 - How would the bounds change if I was consistently using PDFs that include in the fit the same operators that I am fitting?

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PDFS AND NEW PHYSICS





