





## Parton Distributions in the high-precision LHC era

**Juan Rojo** VU Amsterdam & Theory group, Nikhef

**The Zurich Phenomenology Workshop 2018 (ZPW2018) Flavours: Light, Heavy, and Dark** University of Zurich 17/01/2017

## Anatomy of hadronic collisions

In high-energy **hadron colliders** the collisions involve **composite particles** (protons) with internal substructure (quarks and gluons): the LHC is actually a quark/gluon collider!



Calculations of **cross-sections** in hadron collisions require the combination of **perturbative cross-sections** with **non-perturbative parton distribution functions (PDFs)** 

## Why we need better PDFs?

Dominant TH unc for M <sub>W</sub> measurements at LHC							ATLA	S 2017		
Channel	$m_{W^+} - m_{W^-}$ [MeV]	Stat. Unc.	Muon Unc.	Elec. Unc.	Recoil Unc.	Bckg. Unc.	QCD Unc.	EW Unc.	PDF Unc.	Total Unc.
$\begin{array}{l} W \to e\nu \\ W \to \mu\nu \end{array}$	-29.7 -28.6	17.5 16.3	0.0 11.7	4.9 0.0	0.9 1.1	5.4 5.0	0.5 0.4	0.0 0.0	24.1 26.0	30.7 33.2
Combined	-29.2	12.8	3.3	4.1	1.0	4.5	0.4	0.0	23.9	28.0



## The PDF fitting landscape

April 2017	NNPDF3.0	MMHT2014	CT14	HERAPDF2.0	CJ15	ABMP16
Fixed Target DIS	<ul> <li>✓</li> </ul>	<ul> <li></li> </ul>	V	×	<ul> <li></li> </ul>	<ul> <li>✓</li> </ul>
JLAB	×	×	×	×	<ul> <li></li> </ul>	×
HERA I+II	<ul> <li></li> </ul>	<ul> <li>Image: A set of the set of the</li></ul>	<ul> <li>✓</li> </ul>	<ul> <li></li> </ul>	<ul> <li></li> </ul>	~
HERA jets	×	<ul> <li>Image: A second s</li></ul>	×	×	×	×
Fixed Target DY	<ul> <li></li> </ul>	<ul> <li>Image: A second s</li></ul>	<ul> <li></li> </ul>	×	<ul> <li></li> </ul>	<ul> <li></li> </ul>
Tevatron W,Z	<ul> <li></li> </ul>	<ul> <li>Image: A second s</li></ul>	<ul> <li></li> </ul>	×	<ul> <li></li> </ul>	<ul> <li></li> </ul>
Tevatron jets	<ul> <li></li> </ul>	<ul> <li>Image: A second s</li></ul>	<ul> <li></li> </ul>	×	<ul> <li></li> </ul>	×
LHC jets	<ul> <li></li> </ul>	<ul> <li>Image: A set of the set of the</li></ul>	<ul> <li></li> </ul>	×	×	×
LHC vector boson	<ul> <li>✓</li> </ul>	<ul> <li></li> </ul>	<ul> <li>✓</li> </ul>	×	×	<ul> <li>✓</li> </ul>
LHC top	<ul> <li>✓</li> </ul>	×	×	×	×	<ul> <li>✓</li> </ul>
Stat. treatment	Monte Carlo	Hessian Δχ² dynamical	Hessian Δχ² dynamical	Hessian Δχ²=1	Hessian Δχ²=1.645	Hessian Δχ²=1
Parametrization	Neural Networks (259 pars)	Chebyshev (37 pars)	Bernstein (30-35 pars)	Polynomial (14 pars)	Polynomial (24 pars)	Polynomial (15 pars)
HQ scheme	FONLL	TR'	ΑСΟΤ-χ	TR'	ΑСΟΤ-χ	FFN (+BMST)
Order	NLO/NNLO	NLO/NNLO	NLO/NNLO	NLO/NNLO	NLO	NLO/NNLO

### Ubiali, DIS2017

ZPW2018, Zurich, 17/01/2018

### The Structure of the Proton in the LHC Precision Era

Jun Gao<sup>a</sup>, Lucian Harland-Lang<sup>b</sup>, Juan Rojo<sup>c,d</sup>

<sup>a</sup>Institute of Nuclear and Particle Physics, Shanghai Key Laboratory for Particle Physics and Cosmology, School of Physics and Astronomy, Shanghai Jiao Tong University, Shanghai, China <sup>b</sup>Department of Physics and Astronomy, University College London, WC1E 6BT, United Kingdom <sup>c</sup>Department of Physics and Astronomy, VU University, De Boelelaan 1081, 1081HV Amsterdam, The Netherlands <sup>d</sup>Nikhef, Science Park 105, NL-1098 XG Amsterdam, The Netherlands

# Sep 201 14 arXiv:1709.04922v1 [hep-ph]

Abstract

We review recent progress in the determination of the parton distribution functions (PDFs) of the proton, with emphasis on the applications for precision phenomenology at the Large Hadron Collider (LHC). First of all, we introduce the general theoretical framework underlying the global QCD analysis of the quark and gluon internal structure of protons. We then present a detailed overview of the hard-scattering measurements, and the corresponding theory predictions, that are used in state-of-the-art PDF fits. We emphasize here the role that higher-order QCD and electroweak corrections play in the description of recent high-precision collider data. We present the methodology used to extract PDFs in global analyses, including the PDF parametrization strategy and the definition and propagation of PDF uncertainties. Then we review and compare the most recent releases from the various PDF fitting collaborations, highlighting their differences and similarities. We discuss the role that QED corrections and photon-initiated contributions play in modern PDF analysis. We provide representative examples of the implications of PDF fits for high-precision LHC phenomenological applications, such as Higgs coupling measurements and searches for high-mass New Physics resonances. We conclude this report by discussing some selected topics relevant for the future of PDF determinations, including the treatment of theoretical uncertainties, the connection with lattice QCD calculations, and the role of PDFs at future high-energy colliders beyond the LHC.

*Keywords:* Parton Distributions, Quantum Chromodynamics, Large Hadron Collider, Higgs boson, Standard Model, Electroweak theory

166 pages, 82 figures, > 500 references, to appear in Physics Reports

Only time for a brief snapshot here!

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## A story of success



Global PDF fits: highly non-trivial validation of the QCD factorisation framework: i) including O(5000) data points, ii) from O(50) experiments, iii) several of them with ≈1% errors, yet still manage to achieve χ²/N<sub>dat</sub> ≈ 1!

## Higher orders matter

NNPDF3.1	NNLO	NLO
HERA	1.16	1.14
ATLAS	1.09	1.37
CMS	1.06	1.20
LHCb	1.47	1.61
TOTAL	1.148	1.168

**NNPDF3.1** is the first NNPDF analysis where we find that **NNLO** gives a **markedly superior fit quality** as compared to **NLO** 

Direct consequence of the **new high-precision measurements** included

Strongly suggests that NNLO PDFs should be the baseline in all analyses!

## The inner life of protons



## The inner life of protons





At the LHC, precise knowledge of the gluon is required **from small-x to large-x** 

## The large-x gluon from differential top quarks



Top-quark production driven by gluon-gluon scattering

NNLO calculations for stable top quarks available Czakon, Mitov et al 2015-2017

**Data from ATLAS and CMS at 8 TeV** available with breakdown of systematic uncertainties

Final formation included differential top data into NNPDF3.0: constraints on the large-*x* gluon comparable to those of inclusive jet production *Czakon et al* 2017

Improved theory uncertainties in regions crucial for BSM searches, *i.e.*, m<sub>tt</sub> > 1 TeV (while fitting only y<sub>t</sub> and y<sub>tt</sub>)



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## The medium-x gluon from NNLO Z $p_{\rm T}$

 $\frac{1}{2}$  Dominated by **quark-gluon scattering**, thus sensitive to the gluon PDF at intermediate values of *x* 

Malik and Watt 2013, Boughezal et al 2017

**NNLO corrections to the Z p**<sub>T</sub> also available: **up to 10**% **effects** for a measurement that has **sub-percent exp errors** 

Boughezal et al 2015-2017, Gehrmann et al 2015-2017

**Complementary information on the gluon** as compared to inclusive jets and differential top pair production







## The large-x gluon from NNLO jets



NNPDF3.1 NNLO: includes jet data using NNLO evolution and NLO matrix elements, with scale variations as additional TH systematic error

The **jet p**<sub>T</sub> is always used as central scale choice

Also tried variants where ATLAS and CMS 2011 7 TeV data included **using exact NNLO theory** 

Server small impact on the gluon

Moderate improvement of the chi2

Solution Only central bin of ATLAS data included - the large  $\chi^2$  once all bins are included remains there once exact NNLO theory is used

	NNPDF3.1	exact NNLO
CDF Run II $k_t$ jets	0.84	0.85
ATLAS jets $2.76 \text{ TeV}$	1.05	1.03
CMS jets 2.76 $TeV$	1.04	1.02
ATLAS jets 2010 7 TeV	0.96	0.95
ATLAS jets $2011$ 7 TeV	1.06	0.91
CMS jets 7 TeV 2011 7 TeV	0.84	0.79

See also MMHT analysis, Harland-Lang et at 16

MPP, Munich, 26/07/2017

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### One (upgraded) glue to bind them all

### NNPDF3.1 NNLO, Q = 100 GeV



## The small-x gluon from forward charm production

**D** and B meson production from LHCb allow accessing the gluon down to x=10<sup>-6</sup>, well below the HERA coverage *PROSA 2015, Gauld et al 2015* 

Gluon PDF errors **reduced by up to a factor 10**!

Allows robust estimate for the *prompt neutrino flux*, the main background for astrophysical neutrinos at IceCube

 $$\ensuremath{\$}$$  Precision calculation of the **UHE neutrino-nucleus cross-section**, with few-percent TH errors up to  $E_{\nu}=10^{12}$  GeV





**UHE** neutrino-nucleus xsecs



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## The inner life of protons



## Parton distributions with BFKL resummation

- Perturbative fixed-order QCD calculations have been extremely successful in describing a wealth of data from proton-proton and electron-proton collisions
- There are theoretical reasons that eventually we need to go beyond DGLAP: at small-x, **logarithmically enhanced terms in** *1*/*x* **become dominant** and need to be resummed to all orders
- BFKL/high-energy/small-x resummation can be matched to the DGLAP collinear framework, and thus be included into a standard PDF analysis

$$\begin{split} \mathbf{DGLAP}_{\mathbf{Evolution in } \mathbf{Q}^2} & \mu^2 \frac{\partial}{\partial \mu^2} f_i(x,\mu^2) = \int_x^1 \frac{dz}{z} P_{ij}\left(\frac{x}{z},\alpha_s(\mu^2)\right) f_j(z,\mu^2), \\ \mathbf{BFKL}_{\mathbf{Evolution in } \mathbf{x}} & -x \frac{d}{dx} f_+(x,\mu^2) = \int_0^\infty \frac{d\nu^2}{\nu^2} K\left(\frac{\mu^2}{\nu^2},\alpha_s\right) f_+(x,\nu^2) \\ \end{split}$$

$$\end{split}$$

$$\begin{split} \text{Within small-}x \text{ resummation, the N*LO fixed-order DGLAP splitting functions are complemented with the N*LLx contributions from BKFL \\ ABF, CCSS, TW + others, 94-08 & P_{ij}^{N^k \text{LO}+N^h \text{LL}x}(x) = P_{ij}^{N^k \text{LO}}(x) + \Delta_k P_{ij}^{N^h \text{LL}x}(x), \end{split}$$

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## A new world at small-x

Ultimately, the need for (or lack of) BKFL resummation in **ep and pp collider data** can only be assessed by performing a global PDF analysis based on (N)NLO+NLLx theory

Frequencies for the second sec

APFEL: Bertone, Carrazza, Rojo 13 https://apfel.hepforge.org/

HELL: Bonvini, Marzani, Peraro, Muselli 16-17 https://www.ge.infn.it/~bonvini/hell/

NNPDF3.1 (N)NLO+NLL fits stabilize the perturbative PDF expansion at small-x, in particular for the gluon, and markedly improve the fit quality to the small-x HERA data



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## Evidence for BFKL dynamics in HERA data

Using NNLO+NLL*x* theory, the NNLO instability at small-*x* of the  $\chi^2$  disappears

Excellent fit quality to **inclusive and charm HERA** data achieved in the **entire (x,Q<sup>2</sup>) region** 



## Nunca es tarde si la dicha es buena

### **Science** Life and Physics

## After 40 years of studying the strong nuclear force, a revelation

This was the year that analysis of data finally backed up a prediction, made in the mid 1970s, of a surprising emergent behaviour in the strong nuclear force



In the mid 1970s, four Soviet physicists, Batlisky, Fadin, Kuraev and Lipatov, made some predictions involving the strong nuclear force which would lead to their initials entering the lore. "BFKL" became a shorthand for a difficult-to-

### Jon Butterworth

✓ @jonmbutterworth Thu 28 Dec 2017 17.30 GMT



Jon Butterworth, the Guardian, 28/12/2018



## Quark flavour separation from LHC data

- Recent high-precision W and Z production data from ATLAS, CMS, LHCb and D0 allow a better separation of quark and antiquark flavours
- LHCb data in particular provides constrains on the **light quarks at large-***x* (*i.e.* PDF uncertainty reduction by a factor 2 in NNPDF3.1 for  $x \approx 0.1$ )
- In the **MMHT (2016) fit,** the new collider *W*,*Z* data leads to improved determination of  $x(u_V-d_V)$

MMHT2014

0.001

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40

20

0

-20

-40 -10001

23



## The inner life of protons



## How bright is the proton?

- The calculation of QED and electroweak corrections to hadron collider processes requires by consistency to introduce the PDF of the photon in the proton, γ(x,Q)
- First model-independent determination of **γ**(**x**,**Q**) from LHC *W*,*Z* data was **NNPDF2.3QED**, which however affected by **large uncertainties**, due to the limited experimental information
- Recently, y(x,Q) computed in terms of the well-known inclusive structure functions F<sub>2</sub> and F<sub>L</sub>: the resulting photon PDF, exhibits now few-percent uncertainties

$$x\gamma(x,\mu) = \frac{1}{2\pi\alpha(\mu)} \int_{x}^{1} \frac{dz}{z} \left\{ \int_{Q_{\min}^{2}}^{\mu^{2}/(1-z)} \frac{dQ^{2}}{Q^{2}} \alpha^{2}(Q^{2}) \left[ -z^{2}F_{L}(x/z,Q^{2}) + \left( zP_{\gamma q}(z) + \frac{2x^{2}m_{p}^{2}}{Q^{2}} \right) F_{2}(x/z,Q^{2}) \right] - \alpha^{2}(\mu)z^{2}F_{2}(x/z,\mu^{2}) \right\} + \mathcal{O}\left(\alpha\alpha_{s},\alpha^{2}\right)$$

pp $\rightarrow$ H W <sup>+</sup> ( $\rightarrow$ l <sup>+</sup> v) + X at 13 TeV			
non-photon induced contributions	91.2 ± 1.8 fb		
photon-induced contribs (NNPDF23)	6.0 <sup>+4.4</sup> -2.9 fb		
photon-induced contribs (LUXqed)	4.4 ± 0.1 fb		

Manohar, Nason, Salam, Zanderighi, 16-17

Crucial implications for LHC pheno: high-precision determination of photoninitiated (PI) contributions

### Juan Rojo

## Illuminating the photon content of the proton

**NNPDF3.1luxQED**: variant of the NNPDF3.1 global analysis supplemented by

*It the LUX fed theoretical constraints,* 

*Weight NLO QED corrections to DGLAP evolution,* and

Solution Provide the American Solution Solution Provident Solution Soluti Solution S

Fiterative procedure: photon PDF recomputed at each iteration until convergence is achieved



## There is light in all of us

$$\left\langle x\right\rangle _{\gamma}(Q)\equiv\int_{0}^{1}\,dx\,x\gamma(x,Q)$$

Photon momentum fraction in the proton



	$\langle x \rangle_{\gamma} \left( Q = 1.65  \text{GeV} \right)$	$\langle x \rangle_{\gamma} \left( Q = m_Z \right)$
NNPDF3.0QED	$(0.3\pm0.3)\%$	$(0.5\pm0.3)\%$
NNPDF3.1luxQED	$(0.229 \pm 0.003)\%$	$(0.420 \pm 0.003)\%$
LUXqed17	_	$(0.421 \pm 0.003)\%$

In LHC collisions, up to **0.5%** of the momentum carried by the protons corresponds to its photon component!

## Photon-intiated processes at the LHC



NNPDF3.1luxQED results are **consistent** with NNPDF3.0QED, with **much reduced uncertainties** Ş

PI effects can be up to 10% (Drell-Yan) and 30% (WW) at the LHC, with opposite sign wrt EW corrections



ZPW2018, Zurich, 17/01/2018

## The inner life of protons



## A strange conundrum

In most PDF fits, strangeness suppressed wrt **up and down quark sea** due to **neutrino dimuon data** 

Solution of the other hand, recent collider data, in particular the ATLAS W,Z 2011 rapidity distributions, prefer instead a **symmetric strange quark sea** 

$$R_s(x,Q^2) = \frac{s(x,Q^2) + \bar{s}(x,Q^2)}{\bar{u}(x,Q^2) + \bar{d}(x,Q^2)} \begin{cases} \approx 0.5 \text{ (from neutrino, CMS W+c)} \\ \approx 1.0 \text{ (from ATLAS W,Z)} \end{cases}$$

The new ATLAS data can be accommodated in the **global fits**, and *i*) indeed it **increases strangeness**, but not as much as in a collider-only fit, and *ii*) **some tension remain**s between neutrino and collider data



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## How strange is the proton?

In the global fit, the resulting strange PDF is the best compromise between the **pulls from individual** experiments. More strange-sensitive measurements are needed to shed more light in this strange mystery!



## The inner life of protons



## A charming story

In global PDF fits there are **two strategies** to treat the **charm PDF**:

The charm PDF is **generated perturbatively** via collinear splittings **off gluons and light quarks**:

$$\begin{aligned} c(x,Q) &= 0 , \quad Q < m_c \\ c(x,Q) &\simeq \alpha_s(Q^2) \ln \frac{Q^2}{m_c^2} \int_x^1 \frac{dy}{y} P_{gg}^{(0)}(y) g\left(\frac{y}{x},Q^2\right) , \quad Q \ge m_c \end{aligned}$$

The charm PDF is **fitted from data**, that is, it is treated on an **equal footing to light quarks**:

$$c(x, Q_0) = f(x, \{a_i\}), \qquad Q_0 > m_c$$

Note that the first option is necessarily an **assumption**, which can only be validated by comparing with the results of a **direct fit of the charm PDF** 

Moreover fitting the charm PDF offers several advantages:

- **W** Reduce the dependence of high-scale cross-sections with respect of the value of the charm mass
- **Mathematical Security of Figh-precision collider data** that depend on **quark flavor separation**
- **Models** of the **non-perturbative** (*``intrinsic''*) **charm content of the proton**

Fitting the charm PDFs has many advantages even with for a vanishing non-perturbative component

## Some charm really improves things

	NNLO Fitted Charm	NNLO Pert Charm	NLO Fitted Charm	NLO Pert Charm
HERA	1.16	1.21	1.14	1.15
ATLAS	1.09	1.17	1.37	1.45
CMS	1.06	1.09	1.20	1.21
LHCb	1.47	1.48	1.61	1.77
TOTAL	1.148	1.187	1.168	1.197

Fin NNPDF3.1, for collider data, **NNLO theory** leads to a markedly better fit quality that than **NLO** 

Free new precision data included has small experimental errors, with **NNLO corrections mandatory** 

From The global PDF analysis where the charm PDF is fitted leads to a **slightly superior fit quality** than assuming a perturbatively generated charm PDF

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## A charming story

The NNPDF3.1 global analysis is performed with a **fitted charm PDF**. The recent LHC *W*,*Z* data, in particular from LHCb, impose stringent constraints on the **size of the charm PDFs at Q=m**<sub>c</sub>

$$C(Q^2) \equiv \int_0^1 dx \, x \, \left(c(c, Q^2) + \bar{c}(x, Q^2)\right)$$

Charm momentum fraction in the proton



Indications of a **non-perturbative component of the charm PDF localised at large-***x*, though its statistical significance is still limited....

## PDFs and precision physics at the LHC Run II

Compare recent PDF fits with **fiducial W,Z cross-sections at 13 TeV** from ATLAS

Theory calculations with FEWZ and Horace: **NNLO QCD** + **NLO EW**. The latter improves agreement with data

Qualitative agreement between theory and data for most cases - now looking forward to **differential measurements**!





### ZPW2018, Zurich, 17/01/2018

## Summary and outlook

General Recent developments in our understanding of the quark and gluon structure of the proton have been driven by a combination of:

**Theory**: Progress in NNLO QCD and NLO EW calculations for many collider processes: differential top pairs, inclusive jets, the Z transverse momentum ... Also the calculation of the photon PDF in terms of DIS structure functions

**Data**: a wealth of high-precision measurements from HERA, Tevatron, ATLAS, CMS and LHCb, in several cases with sub-percent uncertainties.

**Methodology**: fitted charm PDF, combination/reduction methods for different PDFs, new software for PDF fits, fast NLO/NNLO interfaces, ....

✤ Improvements for many Run II analysis: Higgs couplings, M<sub>W</sub> measurements, heavy BSM particle production, precision SM studies, MC validation, …

**Theory uncertainties** now likely to be a limiting factor in PDF fits and PDF-related studies: urgent need to provide PDF analysis with robust estimates of TH errors, including MHOUs

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