

Parton Distribution Functions and their uncertainties: the polarized case

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Outline

- ① Motivation
 - Issues in Standard PDF determination
- ② The NNDPF procedure
 - A general overview
 - Monte Carlo sampling, Neural Networks and minimization
- ③ Towards NNPDFpol1.0
 - Experimental dataset, PDF parametrization, theoretical constraints
 - Preliminary results: the NNPDFpol1.0 parton set
- ④ Conclusions
 - Summary and outlook

1. Motivation

Issues in standard PDF determination

- Extraction of a set of functions with error bands from a set of data points.
- We need an error band, i.e. a probability density $\mathcal{P}[\Delta q(x)]$ in the space of PDFs:

$$\langle \mathcal{O} \rangle = \int \mathcal{D}\Delta q \mathcal{P}[\Delta q] \mathcal{O}[\Delta q]$$

$$\sigma_{\mathcal{O}}^2 = \int \mathcal{D}\Delta q \mathcal{P}[\Delta q] (\mathcal{O}[\Delta q] - \langle \mathcal{O} \rangle)^2$$

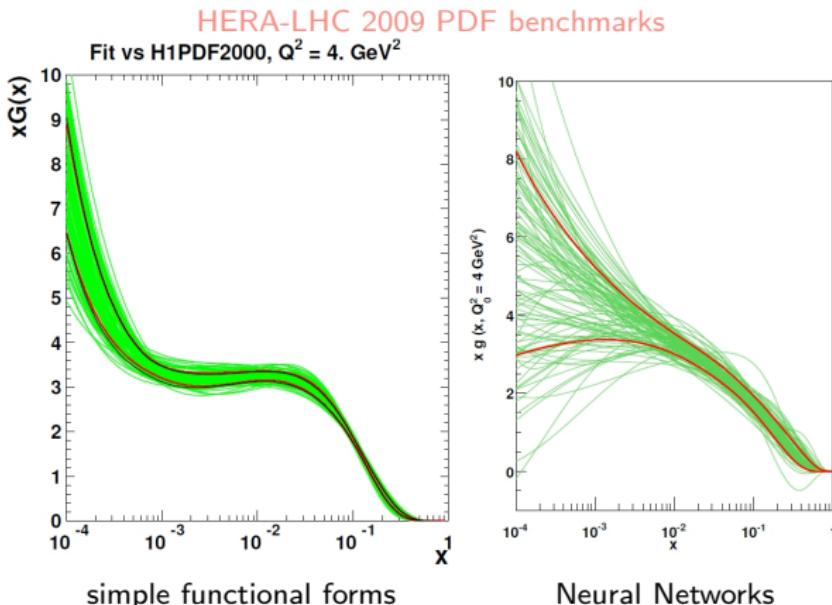
Standard approach

- ① Choose a fixed functional form like
$$\Delta q_i(x, Q_0^2) = A_i x^{b_i} (1-x)^{c_i} (1+\dots)$$
- ② Determine best-fit parameters
- ③ Errors determined via Gaussian linear error propagation

But...

- ① Is the parametrization flexible enough?
- ② What is the error associated to any particular choice?
- ③ Need to rely on linear error propagation

Simple functional forms vs Neural Networks



- Simple functional forms $\Delta q(x) = Ax^b(1-x)^cP(x)$
→ systematic underestimation of uncertainties ⇒ tolerance
- Artificial Neural Networks as universal interpolants
→ reduce theoretical bias from choice of PDF functional form

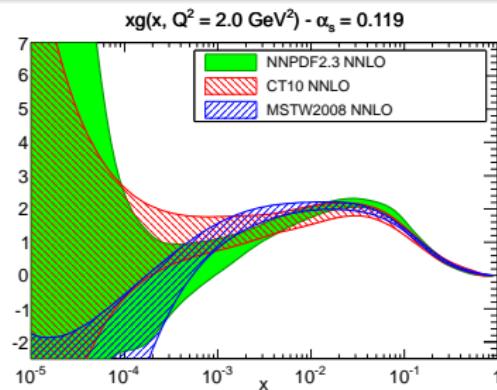
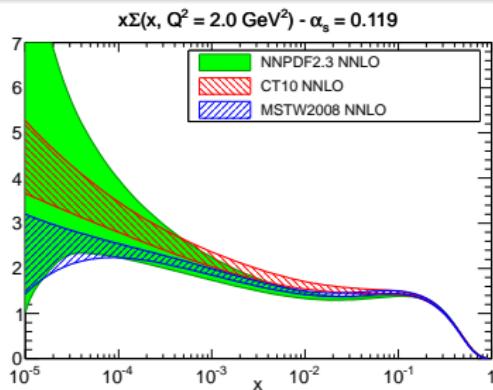
PDF fitting: a new approach

NNPDF: a new approach to PDF fitting based on
Monte Carlo sampling and Neural Networks

	NNPDF1.0	NNPDF1.2	NNPDF2.0	NNPDF2.1 NLO	NNPDF2.1 LO and NNLO	NNPDF2.2/2.3
DIS	✓	✓	✓	✓	✓	✓
Drell-Yan data	✗	✗	✓	✓	✓	✓
Jet data	✗	✗	✓	✓	✓	✓
LHC data	✗	✗	✗	✗	✗	✓
Independent param of the strange and anti-strange	✗	✓	✓	✓	✓	✓
Heavy Quark masses	✗	✗	✗	✓	✓	✓
NNLO	✗	✗	✗	✗	✓	✓

PDF fitting: a new approach

NNPDF: a new approach to PDF fitting based on
Monte Carlo sampling and Neural Networks



The NNPDF Collaboration, arXiv:1207.1303, accepted for publication in Nucl.Phys.B

Routinely used in LHC data analysis and theory prediction
Most recent global fit: NNPDF2.3 available at NLO and NNLO
The first fit including all available LHC data

NNPDF available products: technical details

- ① NNPDF are available as .LHgrid files in which we provide the value of each PDF member at gridded values (x, Q^2) suitable for spline interpolation.
- ② Different .LHgrid files are available for different values of $\alpha_s(M_Z^2)$ and datasets.
- ③ .LHgrid files can be downloaded from the LHAPDF web page

<http://lhapdf.hepforge.org/pdfsets>

and can be used through the LHAPDF interface

- ④ We also provide a Mathematica interface at the NNPDF web page

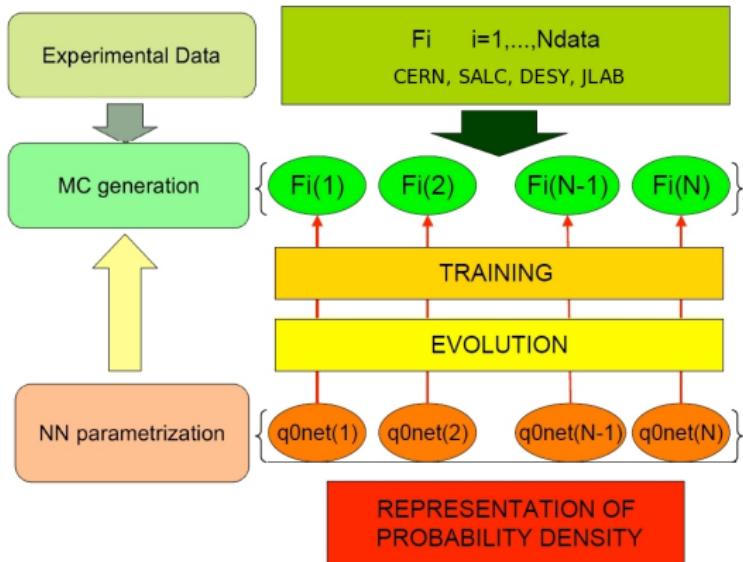
<http://nnpdf.hepforge.org/html/mathematica.html>

Further NNPDF projects

- ① Spin-dependent PDF fits: NNPDFpol1.0
- ② Nuclear PDFs with NNPDF methodology: Nuclear NNPDF1.0

2. The NNPDF approach

A general overview on the methodology



Ingredients:
Monte Carlo sampling and Neural Networks

Ingredient 1: Monte Carlo sampling of experimental data

MONTE CARLO SAMPLING

- Sample the probability density $\mathcal{P}[\Delta q]$ in the space of functions assuming multi-Gaussian data probability distribution

$$g_{1,p}^{(\text{art}), (k)}(x, Q^2) = \left[1 + r_{s,p}^{(k)} \sigma_{s,p} \right] g_{1,p}^{(\text{exp})}(x, Q^2)$$

$\sigma_{s,p}$: statistical errors (also uncorrelated systematics) $r_{s,p}^{(k)}$: Gaussian random numbers

- Generate MC ensemble of N_{rep} replicas with the data probability distribution

Ingredient 1: Monte Carlo sampling of experimental data

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- Generate MC ensemble of N_{rep} replicas with the data probability distribution

MAIN FEATURES

- Expectation values for observables are Monte Carlo integrals

$$\langle \mathcal{O}[\Delta q] \rangle = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \mathcal{O}[\Delta q_k]$$

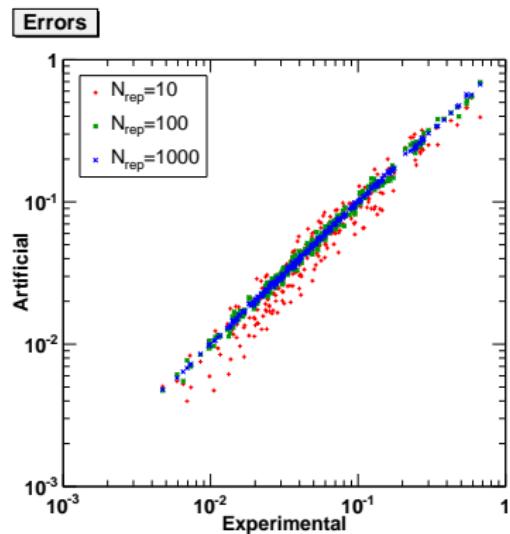
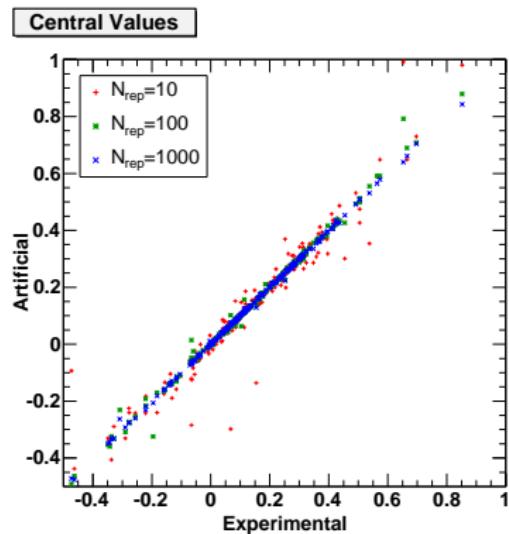
... and the same is true for errors, correlations etc.

- No need to rely on linear propagation of errors
- Possibility to test for non-Gaussian behaviour in fitted PDFs

Ingredient 1: Monte Carlo sampling of experimental data

DETERMINING SAMPLE SIZE

- Require the average over the replicas reproduces central values and errors of the original experimental data to desired accuracy

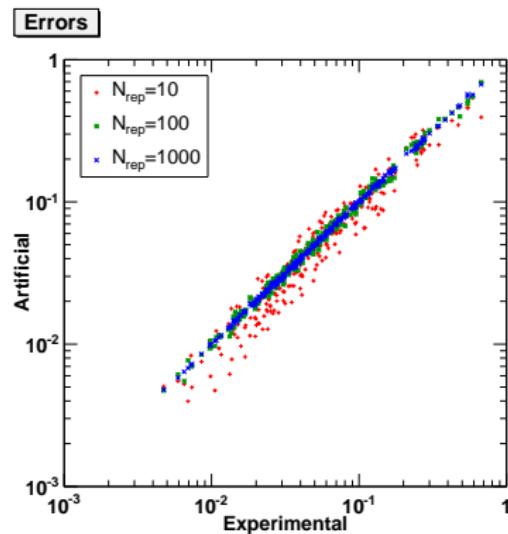
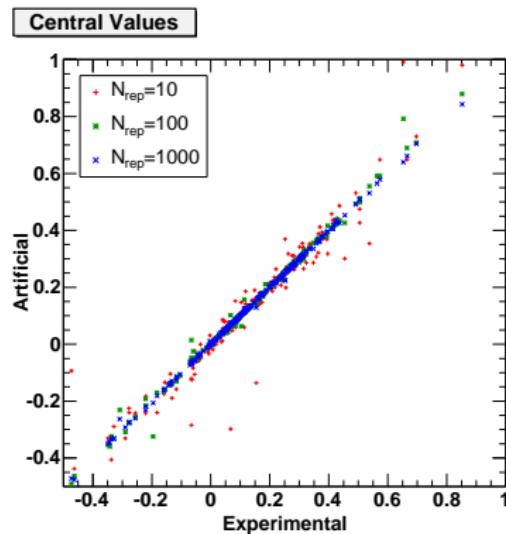


Accuracy of few % requires ~ 100 replicas

Ingredient 1: Monte Carlo sampling of experimental data

DETERMINING SAMPLE SIZE

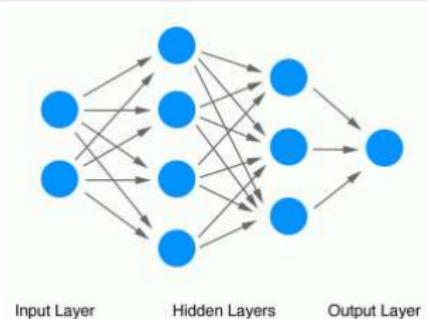
- Require the average over the replicas reproduces central values and errors of the original experimental data to desired accuracy



Accuracy of few % requires ~ 100 replicas

Ingredient 2: Neural Networks

A convenient **functional form**
providing **redundant** and **flexible** parametrization
used as a generator of random functions in the PDF space



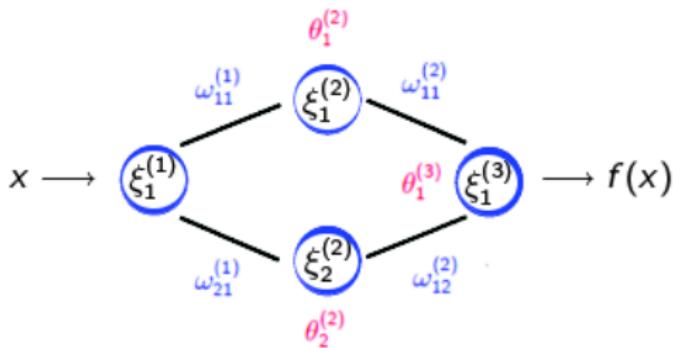
$$\xi_i^{(l)} = g \left(\sum_j^{n_l-1} \omega_{ij}^{(l-1)} \xi_j^{(l-1)} - \theta_i^{(l)} \right)$$

$$g(x) = \frac{1}{1 + e^{-x}}$$

- made of neurons grouped into layers (define the architecture)
- each neuron receives input from neurons in preceding layer (feed-forward NN)
- activation determined by parameters (**weights** and **thresholds**)
- activation determined according to a **non-linear function** (except the last layer)

Ingredient 2: Neural Networks

EXAMPLE: THE SIMPLEST 1-2-1 NN



$$f(x) \equiv \xi_1^{(3)} = \left\{ 1 + \exp \left[\theta_1^{(3)} - \frac{\omega_{11}^{(2)}}{1 + e^{\theta_1^{(2)} - x \omega_{11}^{(1)}}} - \frac{\omega_{12}^{(2)}}{1 + e^{\theta_2^{(2)} - x \omega_{21}^{(1)}}} \right] \right\}^{-1}$$

Recall: $\xi_i^{(l)} = g \left(\sum_j^{n_l-1} \omega_{ij}^{(l-1)} \xi_j^{(l-1)} - \theta_i^{(l)} \right) ; \quad g(x) = \frac{1}{1 + e^{-x}}$

Ingredient 2: Neural Networks

NEURAL NETWORKS

- **Parametrize** each polarized PDF replica with flexible Neural Network

DSSV, AAC, LSS, BB

$\mathcal{O}(10 - 20)$ parameters

NNPDFpol

$\mathcal{O}(200)$ parameters

- **Train** NN to determine the best fit for each replica
- Compute an ensemble of observables and compare to experimental data

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DSSV, AAC, LSS, BB

$\mathcal{O}(10 - 20)$ parameters

NNPDFpol

$\mathcal{O}(200)$ parameters

- Train NN to determine the best fit for each replica
- Compute an ensemble of observables and compare to experimental data

MAIN FEATURES

- Only require **smoothness** of the fitted function
- Do not require any other prejudice on *a priori* functional form
- Reduce the **bias** associated to the choice of some functional form

One more ingredient: minimization and stopping

GENETIC ALGORITHM

Standard minimization unefficient owing to the large parameter space
and non-local x -dependence of the observables

Genetic algorithm provides better exploration of the whole parameter space

- Set Neural Network parameters randomly
- Make clones of the parameter vector and mutate them
- Define a **figure of merit** or error function for the k -th replica

$$E^{(k)} = \frac{1}{N_{\text{rep}}} \sum_{i,j=1}^{N_{\text{rep}}} \left(g_{1,i}^{(\text{art})(k)} - g_{1,i}^{(\text{net})(k)} \right) \left((\text{cov})^{-1} \right)_{ij} \left(g_{1,j}^{(\text{art})(k)} - g_{1,j}^{(\text{net})(k)} \right)$$

$g_{1,i}^{(\text{art})(k)}$: generated from Monte Carlo sampling

$g_{1,i}^{(\text{net})(k)}$: computed from Neural Network PDFs

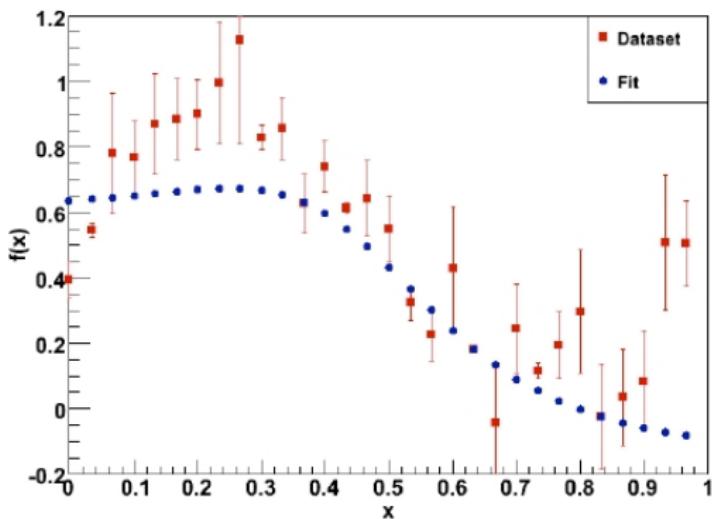
- Select the best set of parameters and perform other manipulations (crossing, mutating, ...) until stability is reached.

One more ingredient: minimization and stopping

DRAWBACK

- NN can learn fluctuations owing to their flexibility

UNDERLEARNING

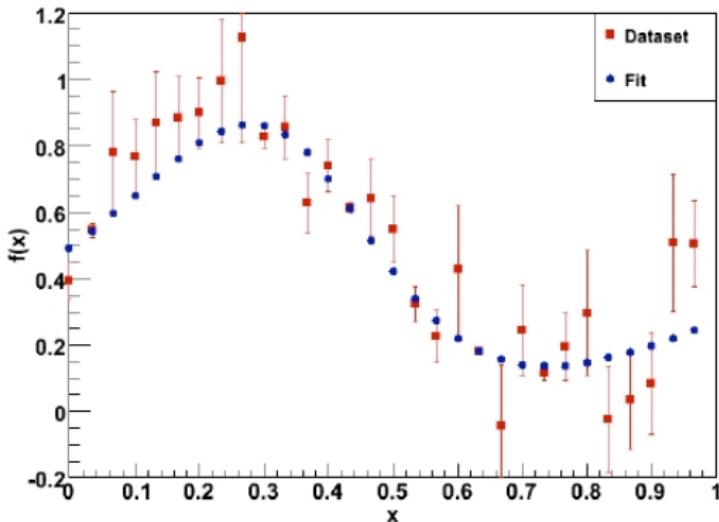


One more ingredient: minimization and stopping

DRAWBACK

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PROPER LEARNING

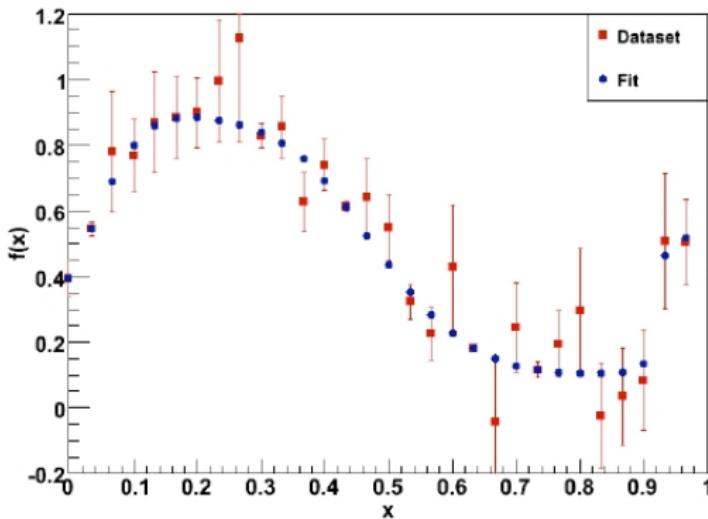


One more ingredient: minimization and stopping

DRAWBACK

- NN can learn fluctuations owing to their flexibility

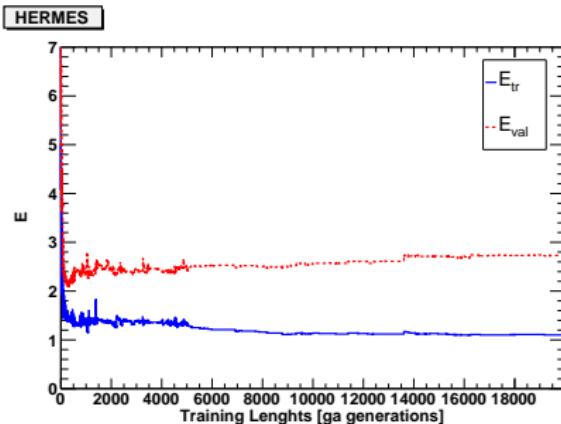
OVERLEARNING



One more ingredient: minimization and stopping

CROSS-VALIDATION METHOD

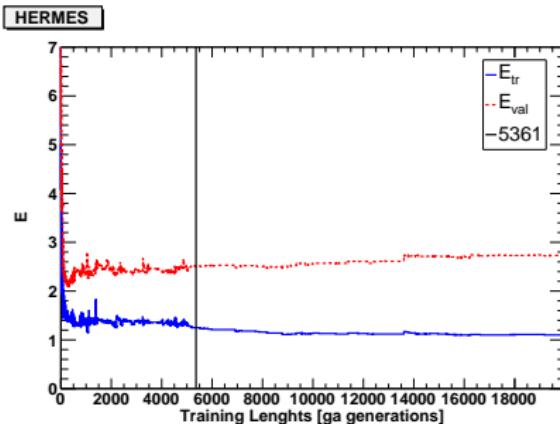
- divide data into two subsets (**training & validation**)
- train the NN on training subset and compute χ^2 for each subset
- stop when χ^2 of validation subset no longer decreases
(NN are learning fluctuations!)



One more ingredient: minimization and stopping

CROSS-VALIDATION METHOD

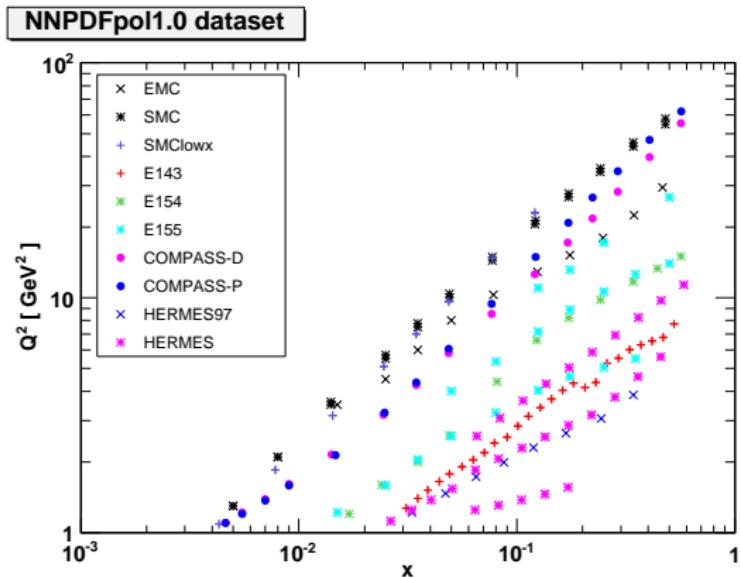
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The best fit does not coincide with the χ^2 absolute minimum

3. Towards NNPDFpol1.0

Experimental dataset



- ① All inclusive polarized DIS data on proton, neutron and deuteron targets
 - ② Kinematical cuts to remove the sensitivity to dynamical higher-twist
 - $Q^2 > 1 \text{ GeV}^2$
 - $W^2 = Q^2(1-x)/x \geq 6.25 \text{ GeV}^2$ (C. Simolo, Ph.D. Thesis. [arXiv:0807.1501](https://arxiv.org/abs/0807.1501))
- (higher twist terms added to observables and fitted to data become compatible with zero)

Input polarized PDF basis

Four polarized PDFs (gluon + linear combinations of light quarks)

① singlet $\Delta\Sigma(x) \equiv \sum_{i=1}^{n_f} \Delta q_i(x)$

② gluon $\Delta g(x)$

③ triplet $\Delta T_3(x) \equiv \Delta u(x) - \Delta d(x)$

④ octet $\Delta T_8(x) \equiv \Delta u(x) + \Delta d(x) - 2\Delta s(x)$

$$\Delta q_i(x, Q^2) = q_i^{\uparrow\uparrow}(x, Q^2) + \bar{q}_i^{\uparrow\uparrow}(x, Q^2) - q_i^{\uparrow\downarrow}(x, Q^2) + \bar{q}_i^{\uparrow\downarrow}(x, Q^2)$$

$$\Delta g(x, Q^2) \equiv g^{\uparrow\uparrow}(x, Q^2) - g^{\uparrow\downarrow}(x, Q^2)$$

Inclusive neutral-current DIS data do not allow disentangling the contributions from q and \bar{q} .
In our notation, Δq takes into account flavor plus anti-flavor contributions.

- At **initial scale** $Q_0^2 = 1 \text{ GeV}^2$
- Assume all heavy quarks are generated radiatively
- Adopt $\alpha_s(M_Z^2) = 0.119$, $m_c = 1.4 \text{ GeV}$, $m_b = 4.75 \text{ GeV}$

PDF Parametrization

$$\Delta\Sigma(x, Q_0^2) = (1-x)^{m_{\Delta\Sigma}} x^{-n_{\Delta\Sigma}} \textcolor{red}{NN}_{\Delta\Sigma}(x)$$

$$\Delta g(x, Q_0^2) = (1-x)^{m_{\Delta g}} x^{-n_{\Delta g}} \textcolor{red}{NN}_{\Delta g}(x)$$

$$\Delta T_3(x, Q_0^2) = A_{\Delta T_3} (1-x)^{m_{\Delta T_3}} x^{-n_{\Delta T_3}} \textcolor{red}{NN}_{\Delta T_3}(x)$$

$$\Delta T_8(X, Q_0^2) = A_{\Delta T_8} (1-x)^{m_{\Delta T_8}} x^{-n_{\Delta T_8}} \textcolor{red}{NN}_{\Delta T_8}(x)$$

- ➊ Each polarized PDF parametrized with a multi-layer feed-forward NN (2-5-3-1)
- ➋ Parametrization supplemented with a preprocessing polynomial:
 - exponents m and n randomly chosen in fixed intervals;
 - intervals must be sufficient large not to introduce a bias on the fit
 - check *a posteriori* by studying asymptotic exponents
- ➌ Overall normalization constant factored out for triplet and octet.

$$A_{\Delta T_3} = \frac{a_3}{\int_0^1 dx [(1-x)^{m_{\Delta T_3}} x^{-n_{\Delta T_3}} \textcolor{red}{NN}_{\Delta T_3}(x)]}$$

$$A_{\Delta T_8} = \frac{a_8}{\int_0^1 dx [(1-x)^{m_{\Delta T_8}} x^{-n_{\Delta T_8}} \textcolor{red}{NN}_{\Delta T_8}(x)]}$$

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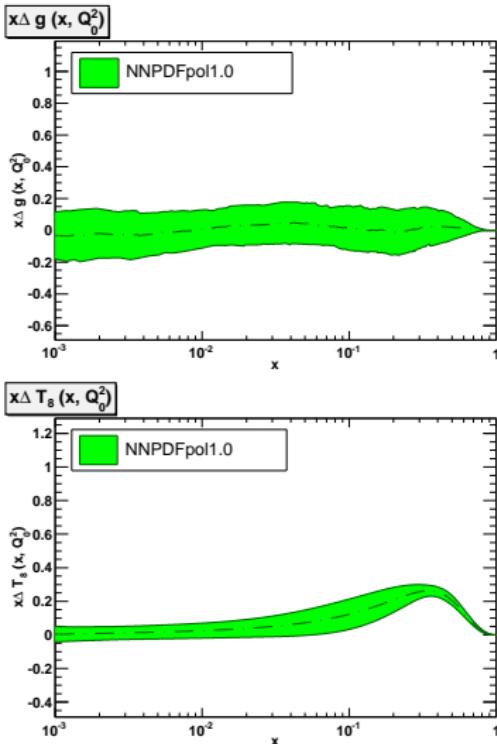
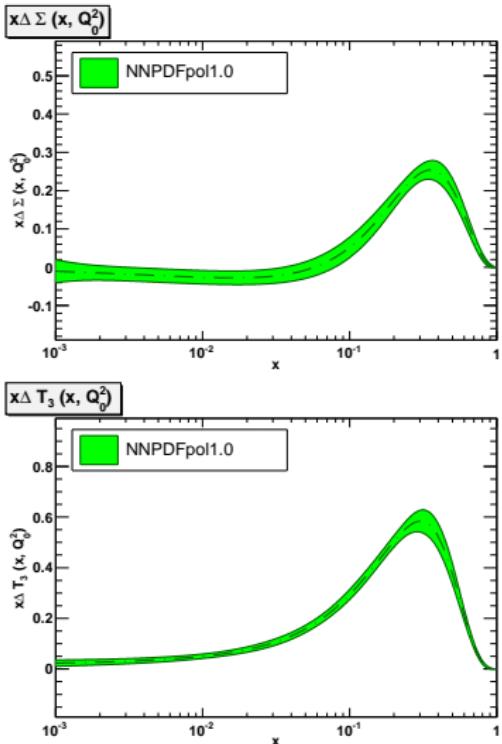
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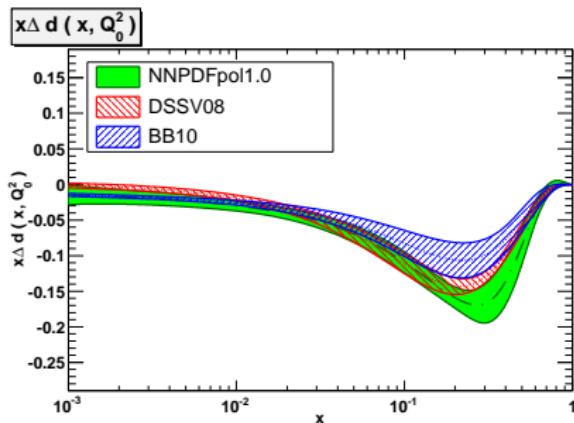
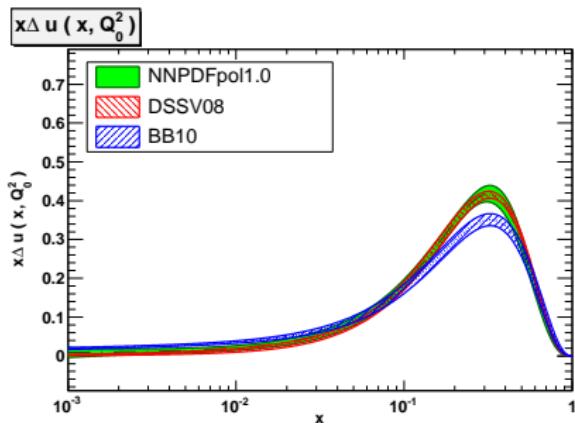
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The NNPDFpol1.0 parton set: parametrization basis



Comparison with other parton sets: flavor basis

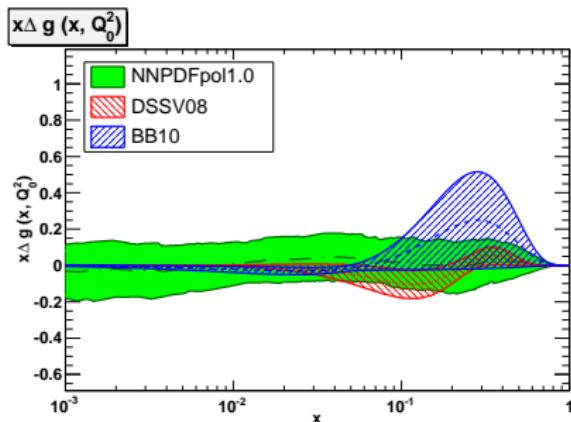
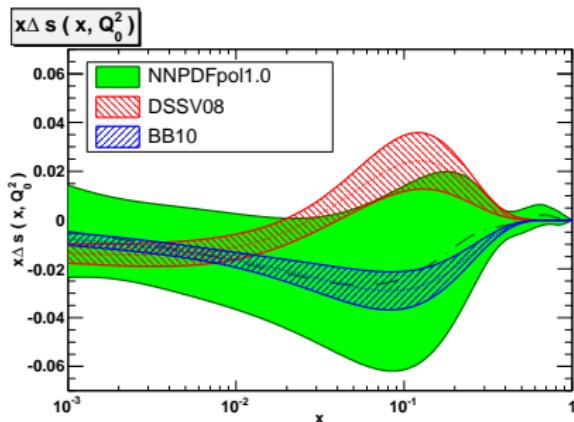
DSSV09: D. de Florian et al., Phys.Rev. D80 (2009) 034030
BB10: J. Blumlein and H. Bottcher, Nucl.Phys. B841 (2010) 205



- NNPDF Δu and Δd uncertainties comparable with other PDF sets
- Much larger error band for strangeness and gluon PDFs

Comparison with other parton sets: flavor basis

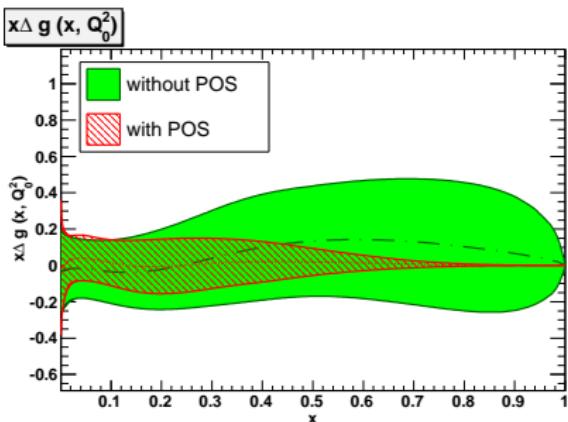
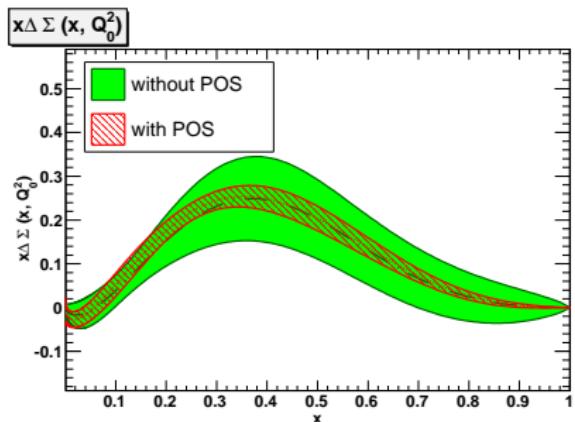
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Theoretical constraints: positivity

Polarized PDFs are only loosely constrained by the data.
At high- x positivity constraints are as important as data!



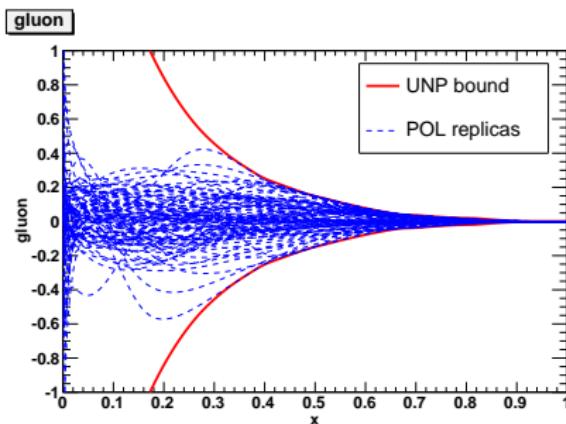
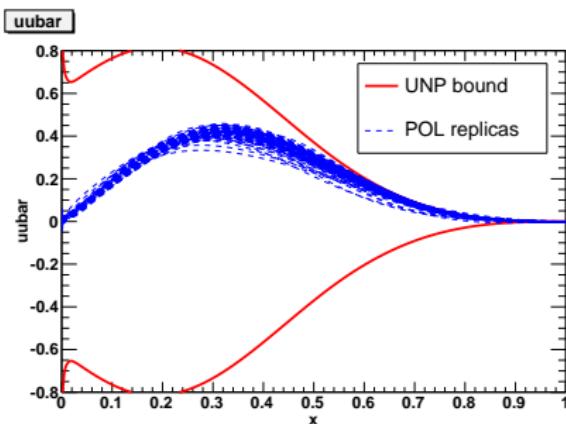
$$\chi_{\text{tot}}^2 = 0.72 \text{ (no positivity constraints)}$$
$$\chi_{\text{tot}}^2 = 0.75 \text{ (with positivity constraints)}$$

Theoretical constraints: positivity

We impose LO positivity, i.e.

$$|\delta f(x, Q_0^2)| \leq f(x, Q_0^2)$$

at the parametrization scale $Q_0 = 1 \text{ GeV}^2$ separately for each flavor (u,d,s) and gluon (g).



The spin content of the proton

Singlet and Gluon first moments in $\overline{\text{MS}}$ scheme at $Q_0^2 = 4 \text{ GeV}^2$

	NNPDFpol1.0	DSSV08 (tr)	AAC08 (pos Δg)	BB10	LSS10 (pos Δg)
$[\Delta\Sigma]$	0.26 ± 0.16	0.37 ± 0.02	0.24 ± 0.07	0.19 ± 0.08	0.21 ± 0.03
$[\Delta g]$	0.1 ± 1.7	-0.03 ± 0.16	0.63 ± 0.81	0.46 ± 0.43	0.32 ± 0.19

Notice the large uncertainty on the first moments:

Singlet between two and four times

Gluon almost one order of magnitude

DSSV09: D. de Florian et al., Phys.Rev. D80 (2009) 034030

AAC08: M. Hirai and S. Kumano, Nucl.Phys. B813 (2009) 106

BB10: J. Blumlein and H. Bottcher, Nucl.Phys. B841 (2010) 205

LSS10: E. Leader, A.V. Sidorov and D.B. Stamenov, Phys.Rev. D82 (2010) 114018

4. Conclusions

Final remarks

Summary

- ① The NNPDF methodology provides a statistically sound procedure for PDF fitting based on:
 - Monte Carlo sampling of experimental data;
 - Neural Networks used as unbiased interpolants.
- ② NNPDFpol1.0 is the first polarized parton determination using the NNPDF methodology:
 - Δu and Δd quark combinations have uncertainties similar to other parton sets;
 - Δs and Δg have larger uncertainties (both can be compatible with zero)
⇒ what is their (true) contribution to the proton spin content?

Outlook

- ① Include data sets from other processes (open charm and jet production with fixed target, inclusive jet production, W boson production at RHIC, ...) in order to test how uncertainty bands can be improved, in particular Δs and Δg .

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Thank you for your attention!

5. Backup

PDF fitting: state of the art

- ➊ First stage: first **moments** of polarized PDFs and polarized **sum rules** (last 25 years)
→ “historical” experimental collaborations (at CERN, SLAC, DESY, JLAB): EMC, SMC, E142, E143, E154, E155, COMPASS, HERMES, CLAS, ...
- ➋ Second stage: polarized **PDF fits** from **global NLO QCD analysis** (last \sim 15 years)
→ different choice of datasets, parton parametrization, treatment of higher twists, ...
ABFR ([arXiv:hep-ph/9803237](https://arxiv.org/abs/hep-ph/9803237), 1998), BB ([arXiv:1005.3113](https://arxiv.org/abs/1005.3113), 2010) (DIS only); AAC ([arXiv:0808.0413](https://arxiv.org/abs/0808.0413), 2008), LSS ([arXiv:1010.0574](https://arxiv.org/abs/1010.0574), 2010) (DIS+SIDIS); DSSV ([arXiv:0904.3821](https://arxiv.org/abs/0904.3821), 2009) (DIS+SIDIS+pp)
- ➌ Third stage: provide **uncertainties** on polarized PDFs (last \sim 10 years)
→ Gaussian error propagation, Lagrange multiplier + Hessian method; fit with orthogonal polynomials ([arXiv:1011.4873](https://arxiv.org/abs/1011.4873), 2010)

Polarized PDF evolution

In Mellin space the DGLAP equations

$$\begin{aligned}\mu^2 \frac{\partial}{\partial \mu^2} \Delta q_{NS}^{\pm,\nu}(N, \mu^2) &= \Delta \gamma_{NS}^{\pm,\nu} q_{NS}^{\pm,\nu}(N, \mu^2) \\ \mu^2 \frac{\partial}{\partial \mu^2} \begin{pmatrix} \Delta \Sigma \\ \Delta g \end{pmatrix}(N, \mu^2) &= \begin{pmatrix} \Delta \gamma_{qq}(N, \alpha_s(Q^2)) & \Delta \gamma_{qg}(N, \alpha_s(Q^2)) \\ \Delta \gamma_{gq}(N, \alpha_s(Q^2)) & \Delta \gamma_{gg}(N, \alpha_s(Q^2)) \end{pmatrix} \begin{pmatrix} \Delta \Sigma \\ \Delta g \end{pmatrix}\end{aligned}$$

can be solved analitically

$$\Delta q_{NS}^{\pm,\nu}(N, Q^2) = \Gamma_{NS}^{\pm,\nu}(N, a_s, a_0) \Delta q_{NS}^{\pm,\nu}(N, Q_0^2), a_s \equiv \alpha_s/2\pi$$

where, at NLO,

$$\Gamma_{NS,NLO}^{\pm,\nu}(N, a_s, a_0) = \exp \left\{ \frac{U_1^{\pm,\nu}}{b_1} \ln \left(\frac{1 + b_1 a_s}{1 + b_1 a_0} \right) \right\} \left(\frac{a_s}{a_0} \right)^{-R_0^{NS}}$$

Polarized PDF evolution

NNPDF NLO polarized PDF evolution (**Fast Kernel method**) benchmarked with the Les Houches PDF benchmarks ([G. Salam and a. Vogt, hep-ph/0511119](#))

x	$\epsilon_{\text{rel}}(\Delta u_V)$	$\epsilon_{\text{rel}}(\Delta d_V)$	$\epsilon_{\text{rel}}(\Delta \Sigma)$	$\epsilon_{\text{rel}}(\Delta g)$
10^{-3}	$1.1 \cdot 10^{-4}$	$9.2 \cdot 10^{-5}$	$9.9 \cdot 10^{-5}$	$1.1 \cdot 10^{-4}$
10^{-2}	$1.4 \cdot 10^{-4}$	$1.9 \cdot 10^{-4}$	$3.5 \cdot 10^{-4}$	$9.3 \cdot 10^{-5}$
0.1	$1.2 \cdot 10^{-4}$	$1.6 \cdot 10^{-4}$	$5.4 \cdot 10^{-6}$	$1.7 \cdot 10^{-4}$
0.3	$2.3 \cdot 10^{-6}$	$1.1 \cdot 10^{-5}$	$7.5 \cdot 10^{-6}$	$1.7 \cdot 10^{-5}$
0.5	$5.6 \cdot 10^{-6}$	$9.6 \cdot 10^{-6}$	$1.6 \cdot 10^{-5}$	$2.5 \cdot 10^{-5}$
0.7	$1.2 \cdot 10^{-4}$	$9.2 \cdot 10^{-7}$	$1.6 \cdot 10^{-4}$	$7.8 \cdot 10^{-5}$
0.9	$3.5 \cdot 10^{-3}$	$1.1 \cdot 10^{-2}$	$4.1 \cdot 10^{-3}$	$7.8 \cdot 10^{-3}$

Very accurate evolution!

Target mass corrections

- Extracting both structure functions $g_1(x, Q^2)$ and $g_2(x, Q^2)$ from data requires measuring longitudinal and transverse spin asymmetries A_{\parallel} and A_{\perp}
- Experimental information on A_{\perp} is rather poor (in most cases only A_{\parallel} is measured), thus $g_1(x, Q^2)$ and $g_2(x, Q^2)$ are related

$$g_1(x, Q^2) = \frac{F_1(x, Q^2)}{1 + \gamma\eta} \frac{A_{\parallel}}{D} + \frac{\gamma(\gamma - \eta)}{\gamma\eta + 1} g_2(x, Q^2)$$

$$\gamma = \frac{2m_N x}{Q} ; \quad \eta = \frac{\epsilon\gamma y}{1 - \epsilon(1 - y)} ; \quad D = \frac{1 - (1 - y)\epsilon}{1 + \epsilon R(x, Q^2)} ; \quad \epsilon = \frac{4(1 - y) - \gamma^2 y^2}{2y^2 + 4(1 - y) + \gamma^2 y^2} ; \quad y = 1 - \frac{E'}{E}$$

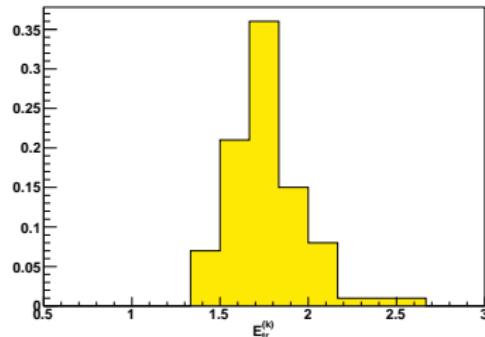
MUST MAKE SOME ASSUMPTION ON $g_2(x, Q^2)$

$g_2 = 0$ OR $g_2 = g_2^{WW}$ (relate g_2 to g_1 by means of the Wandzura-Wilczek relation)

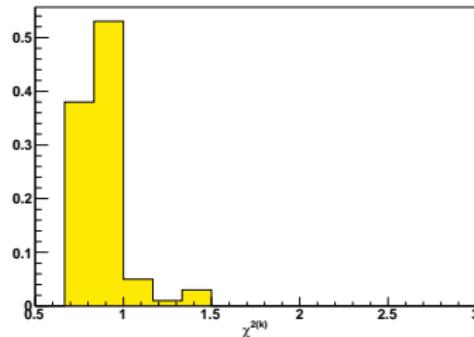
- Target mass corrections implemented iteratively during the minimization procedure

NNPDFpol1.0: global χ^2

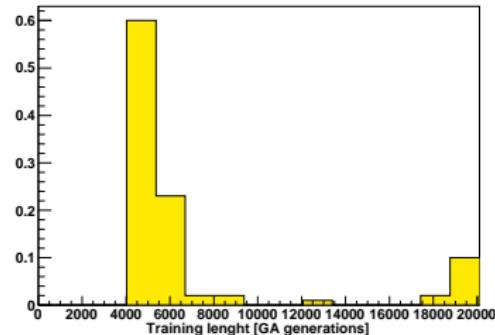
E_{tr} distribution for MC replicas



$\chi^{2(k)}$ distribution for MC replicas

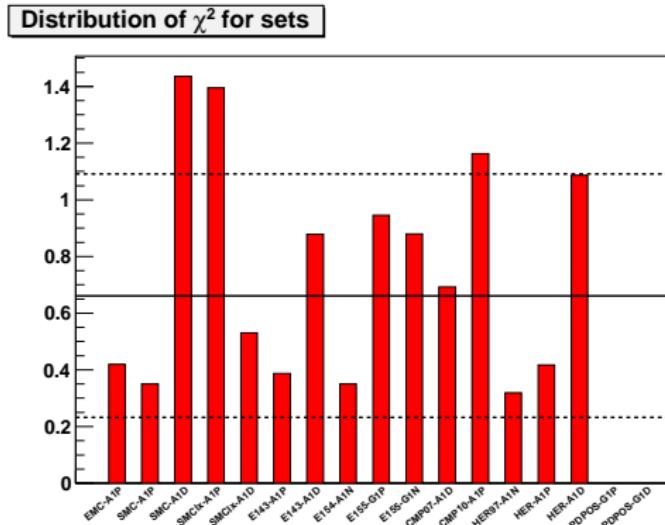


Distribution of training lengths



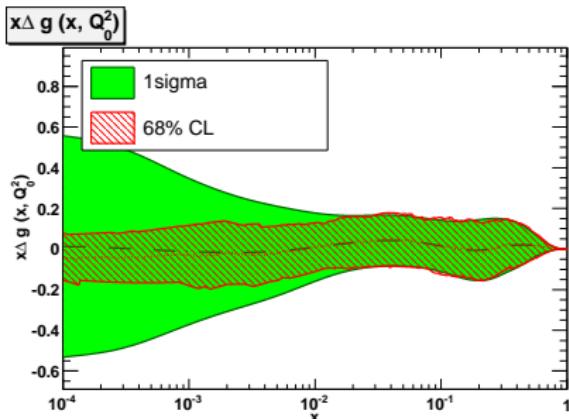
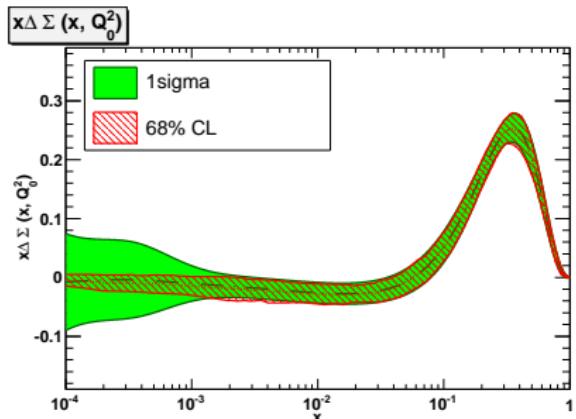
χ^2_{tot}	0.76
$\langle E \rangle \pm \sigma_E$	1.81 ± 0.19
$\langle E_{tr} \rangle \pm \sigma_{E_{tr}}$	1.59 ± 0.56
$\langle E_{val} \rangle \pm \sigma_{E_{val}}$	1.76 ± 0.72
$\langle \chi^{2(k)} \rangle \pm \sigma_{\chi^2}$	0.88 ± 0.13
$\langle TL \rangle$	7247

NNPDFpol1.0: individual experiments χ^2



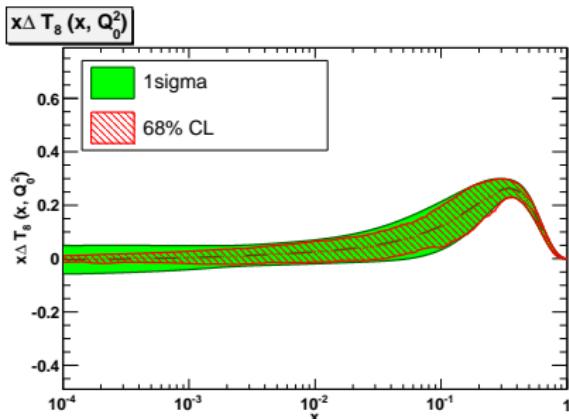
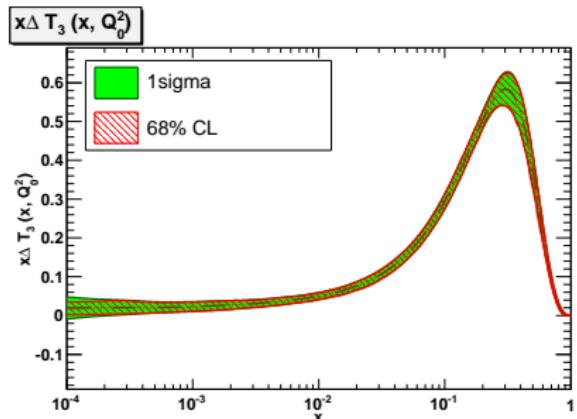
- No evidence of any specific dataset being inconsistent with each other
- Distribution of individual χ^2 values broadly consistent with statistical expectations

NNPDFpol1.0: 68% confidence levels



comparison between 1σ error bands and 68% confidence level
test for non-Gaussian behaviour
sizable deviations from Gaussian behaviour in the extrapolation region

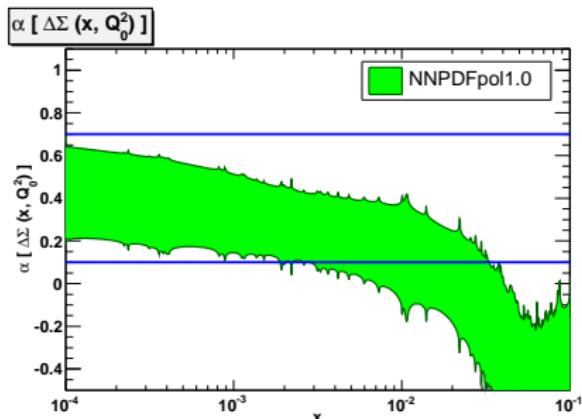
NNPDFpol1.0: 68% confidence levels



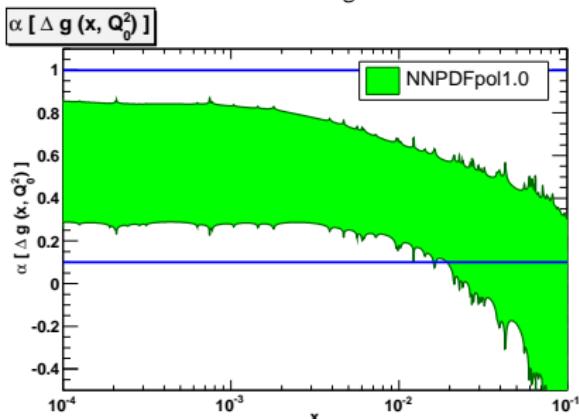
comparison between 1σ error bands and 68% confidence level
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sizable deviations from Gaussian behaviour in the extrapolation region

Preprocessing: effective asymptotic exponents

$$0.0 \leq n_{\Delta\Sigma} \leq 0.6$$



$$0.0 \leq n_{\Delta g} \leq 0.6$$

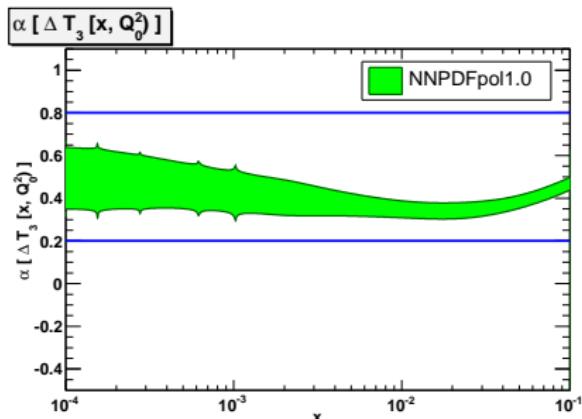


$$\alpha_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1/x) \text{ at } Q^2 = Q_0^2 = 1 \text{ GeV}^2$$

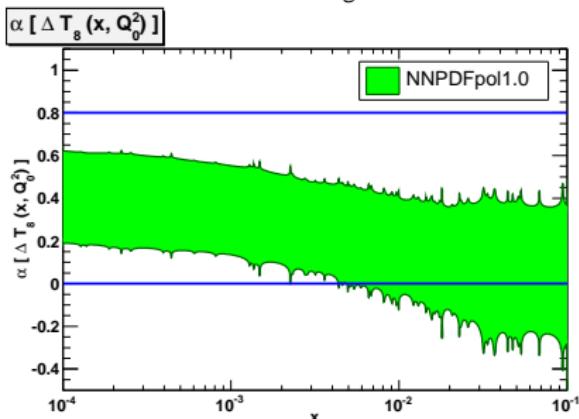
Effective exponents always contained in the preprocessing exponents range
The polarized PDF is driven only by experimental data

Preprocessing: effective asymptotic exponents

$$0.0 \leq n_{\Delta\Sigma} \leq 0.6$$



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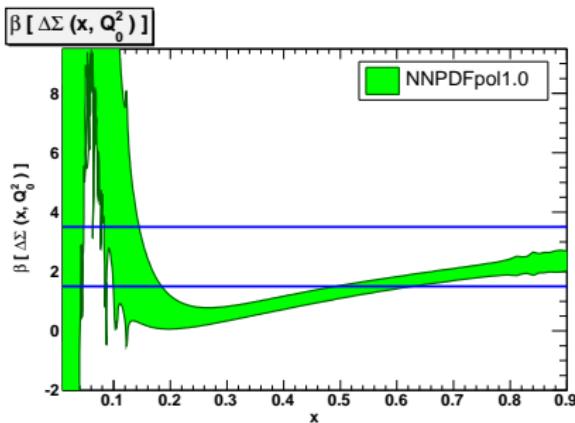


$$\alpha_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1/x) \text{ at } Q^2 = Q_0^2 = 1 \text{ GeV}^2$$

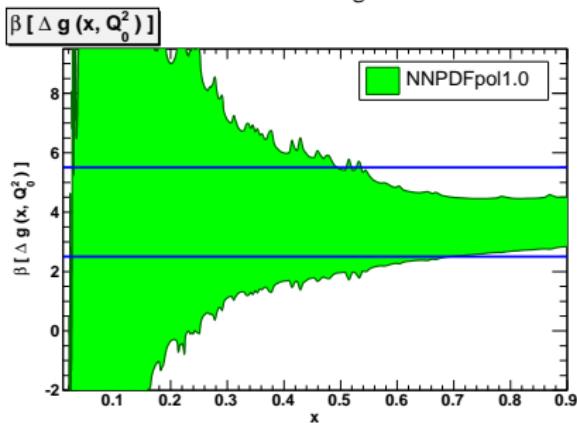
Effective exponents always contained in the preprocessing exponents range
The polarized PDF is driven only by experimental data

Preprocessing: effective asymptotic exponents

$$0.5 \leq m_{\Delta\Sigma} \leq 5.0$$



$$0.5 \leq m_{\Delta g} \leq 4.0$$

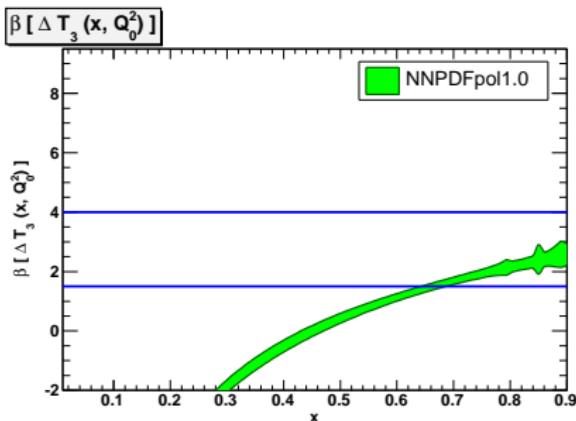


$$\beta_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1-x) \text{ at } Q^2 = Q_0^2 = 1 \text{ GeV}^2$$

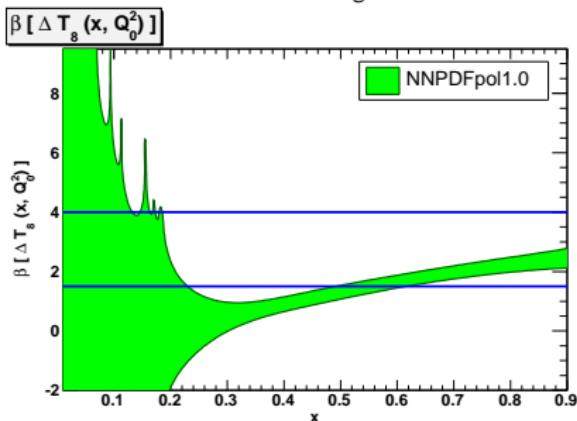
Effective exponents always contained in the preprocessing exponents range
The polarized PDF is driven only by experimental data

Preprocessing: effective asymptotic exponents

$$0.5 \leq m_{\Delta\Sigma} \leq 4.0$$



$$0.5 \leq m_{\Delta g} \leq 7.0$$



$$\beta_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1-x) \text{ at } Q^2 = Q_0^2 = 1 \text{ GeV}^2$$

Effective exponents always contained in the preprocessing exponents range
The polarized PDF is driven only by experimental data

Distances

Compare two sets of $N_{\text{rep}}^{(1)}$ and $N_{\text{rep}}^{(2)}$ replicas coming from different fits
Do they have belong to the same underlying probability distribution?

MEAN VALUE

$$d^2 \left(\langle q^{(k)} \rangle_{(1)}, \langle q^{(k)} \rangle_{(2)} \right) = \frac{\left(\langle q^{(k)} \rangle_{(1)} - \langle q^{(k)} \rangle_{(2)} \right)^2}{\sigma^2 \left[\langle q^{(k)} \rangle_{(1)} \right] + \sigma^2 \left[\langle q^{(k)} \rangle_{(2)} \right]}$$

$$\langle q^{(k)} \rangle_{(i)} = \frac{1}{N_{\text{rep}(i)}} \sum_{l=1}^{N_{\text{rep}(i)}} q_l^{(k)} \quad \sigma^2 \left[\langle q^{(k)} \rangle_{(i)} \right] = \frac{1}{N_{\text{rep}(i)}} \sigma^2 \left[q_{(i)}^{(k)} \right] = \frac{1}{N_{\text{rep}(i)} - 1} \sum_{l=1}^{N_{\text{rep}(i)}} \left(q_{l,(i)} - \langle q \rangle_{(i)} \right)^2$$

UNCERTAINTY

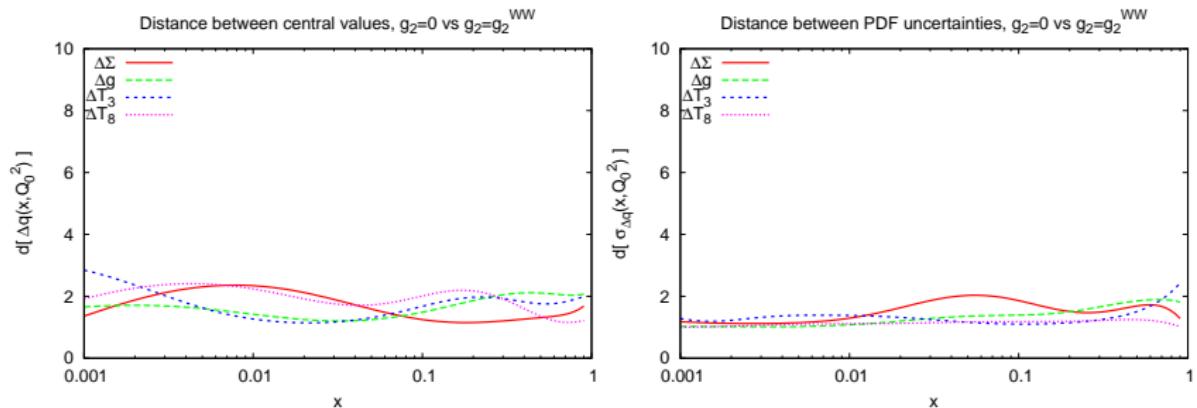
$$d^2 \left(\sigma_{(1)}^2, \sigma_{(2)}^2 \right) = \frac{\left(\bar{\sigma}_{(1)}^2 - \bar{\sigma}_{(2)}^2 \right)}{\sigma^2 \left[\sigma_{(1)}^2 \right] + \sigma^2 \left[\sigma_{(2)}^2 \right]}$$

$$\bar{\sigma}_{(i)}^2 \equiv \sigma^2 \left[q_{(i)}^{(k)} \right] \quad \sigma^2 \left[\sigma_{(i)}^2 \right] = \frac{1}{N_{\text{rep}(i)}} \left[\frac{1}{N_{\text{rep}(i)}} \sum_{l=1}^{N_{\text{rep}(i)}} \left(q_{l,(i)} - \langle q \rangle_{(i)} \right)^4 - \frac{N_{\text{rep}(i)} - 3}{N_{\text{rep}(i)} - 1} \left(\bar{\sigma}_{(i)}^2 \right)^2 \right]$$

By definition, the distances have a χ^2 probability distribution with one degree of freedom
mean $\langle d \rangle = 1$ and $d \lesssim 2.3$ at 90% confidence level

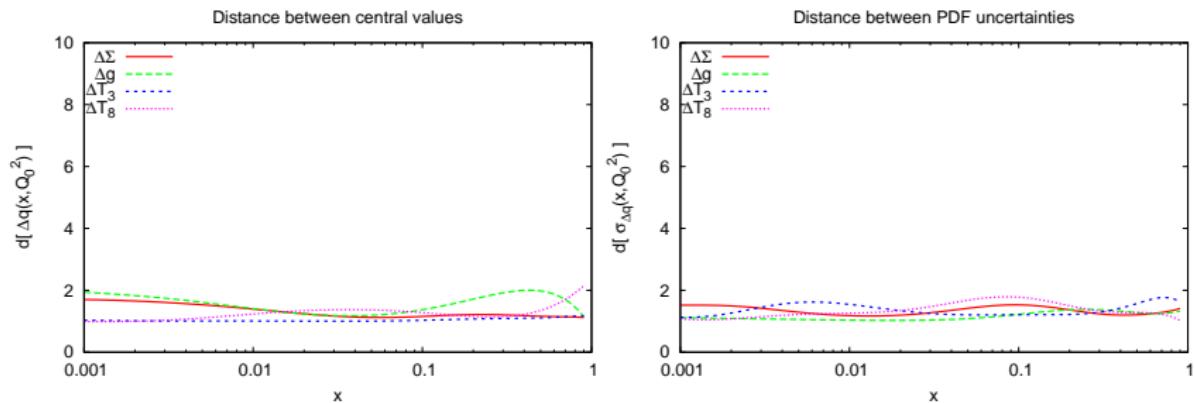
Distances

Comparison between different inclusion of Target Mass Corrections



Distances

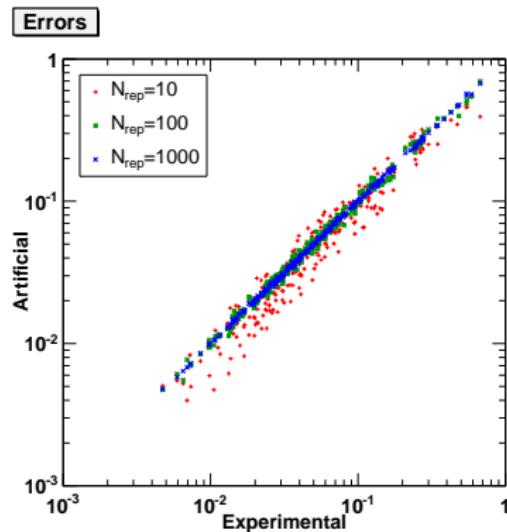
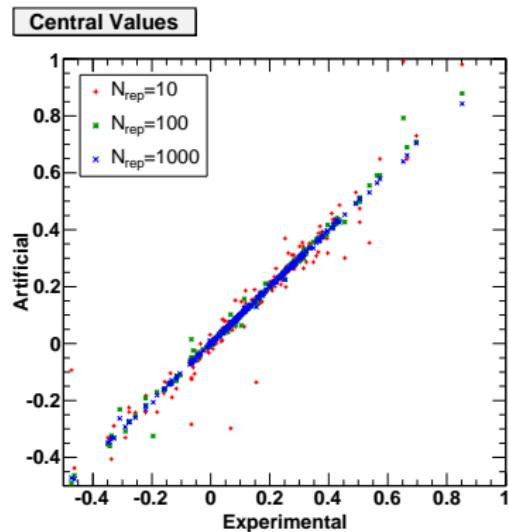
Impact of uncertainty on a8 coupling



Ingredient 1: Monte Carlo sampling of experimental data

DETERMINING SAMPLE SIZE

- Require the average over the replicas reproduces central values and errors of the original experimental data to desired accuracy



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Quantitative approach: devise proper statistical estimators

	$\left\langle PE \left[\langle g_1^{(\text{art})} \rangle \right] \right\rangle [\%]$				$r \left[g_1^{(\text{art})} \right]$		
N_{rep}	10	100	1000		10	100	1000
EMC	23.7	3.5	2.9		.76037	.99547	.99712
SMC	19.4	5.6	1.2		.94789	.99908	.99993
...

$$\left\langle PE \left[\langle F^{(\text{art})} \rangle_{\text{rep}} \right] \right\rangle_{\text{dat}} = \frac{1}{N_{\text{dat}}} \sum_{i=1}^{N_{\text{dat}}} \left| \frac{\langle F_i^{(\text{art})} \rangle_{\text{rep}} - F_i^{(\text{exp})}}{F_i^{(\text{exp})}} \right| \quad \text{Percentage Error}$$

$$r \left[F^{(\text{art})} \right] = \frac{\left\langle F^{(\text{exp})} \langle F^{(\text{art})} \rangle_{\text{rep}} \right\rangle_{\text{dat}} - \langle F^{(\text{exp})} \rangle_{\text{dat}} \left\langle \langle F^{(\text{art})} \rangle_{\text{rep}} \right\rangle_{\text{dat}}}{\sigma_s^{(\text{exp})} \sigma_s^{(\text{art})}} \quad \text{Scatter Correlation}$$

Ingredient 1: Monte Carlo sampling of experimental data

DETERMINING SAMPLE SIZE

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Quantitative approach: devise proper statistical estimators

	$\left\langle PE \left[\langle \delta g_1^{(\text{art})} \rangle \right] \right\rangle [\%]$				$r \left[\delta g_1^{(\text{art})} \right]$	
N_{rep}	10	100	1000	10	100	1000
EMC	12.8	4.9	2.0	.97397	.99521	.99876
SMC	22.4	5.4	1.7	.96585	.99489	.99980
...

$$\left\langle PE \left[\langle F^{(\text{art})} \rangle_{\text{rep}} \right] \right\rangle_{\text{dat}} = \frac{1}{N_{\text{dat}}} \sum_{i=1}^{N_{\text{dat}}} \left| \frac{\langle F_i^{(\text{art})} \rangle_{\text{rep}} - F_i^{(\text{exp})}}{F_i^{(\text{exp})}} \right| \quad \text{Percentage Error}$$

$$r \left[F^{(\text{art})} \right] = \frac{\left\langle F^{(\text{exp})} \langle F^{(\text{art})} \rangle_{\text{rep}} \right\rangle_{\text{dat}} - \langle F^{(\text{exp})} \rangle_{\text{dat}} \left\langle \langle F^{(\text{art})} \rangle_{\text{rep}} \right\rangle_{\text{dat}}}{\sigma_s^{(\text{exp})} \sigma_s^{(\text{art})}} \quad \text{Scatter Correlation}$$

Ingredient 1: Monte Carlo sampling of experimental data

DETERMINING SAMPLE SIZE

- Require the average over the replicas reproduces central values and errors of the original experimental data to desired accuracy

Accuracy of few % requires ~ 100 replicas