

Recent progress in PDFs

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Anatomy of hadronic collisions

In high-energy **hadron colliders** the collisions involve **composite particles** (protons) with internal substructure (quarks and gluons): the LHC is actually a quark/gluon collider!



Calculations of **cross-sections** in hadron collisions require the combination of **perturbative cross-sections** with **non-perturbative parton distribution functions (PDFs)**

Why we need better PDFs?

Dominant TH unc for M _W measurements at LHC							LHC		ATLA	.S 2017
Channel	$m_{W^+} - m_{W^-}$ [MeV]	Stat. Unc.	Muon Unc.	Elec. Unc.	Recoil Unc.	Bckg. Unc.	QCD Unc.	EW Unc.	PDF Unc.	Total Unc.
$\begin{array}{l} W \to e\nu \\ W \to \mu\nu \end{array}$	-29.7 -28.6	17.5 16.3	0.0 11.7	4.9 0.0	0.9 1.1	5.4 5.0	0.5 0.4	0.0 0.0	24.1 26.0	30.7 33.2
Combined	-29.2	12.8	3.3	4.1	1.0	4.5	0.4	0.0	23.9	28.0





Recent progress in proton PDF fits



The PDF fitting landscape

April 2017	NNPDF3.0	MMHT2014	CT14	HERAPDF2.0	CJ15	ABMP16
Fixed Target DIS	 	 	v	×	 	v
JLAB	×	×	×	×	V	×
HERA I+II	 	 	v	 	 ✓ 	~
HERA jets	×	 	×	×	×	×
Fixed Target DY	 	 Image: A second s	 ✓ 	×	 	~
Tevatron W,Z	 	 Image: A set of the set of the	 	×	 	~
Tevatron jets	 	 Image: A set of the set of the	 	×	 	×
LHC jets	 	 Image: A set of the set of the	 	×	×	×
LHC vector boson	 ✓ 	V	~	×	×	~
LHC top	 ✓ 	×	×	×	×	 ✓
Stat. treatment	Monte Carlo	Hessian Δχ² dynamical	Hessian Δχ² dynamical	Hessian Δχ²=1	Hessian Δχ²=1.645	Hessian Δχ²=1
Parametrization	Neural Networks (259 pars)	Chebyshev (37 pars)	Bernstein (30-35 pars)	Polynomial (14 pars)	Polynomial (24 pars)	Polynomial (15 pars)
HQ scheme	FONLL	TR'	ΑСΟΤ-χ	TR'	ΑСΟΤ-χ	FFN (+BMST)
Order	NLO/NNLO	NLO/NNLO	NLO/NNLO	NLO/NNLO	NLO	NLO/NNLO

Ubiali, DIS2017

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The inner life of protons



The inner life of protons





At the LHC, precise knowledge of the gluon is required **from small-x to large-x**

The large-x gluon from differential top quarks



Fop-quark pair production driven by gluon-gluon lumi

NNLO calculations for stable top quarks available Czakon, Mitov et al 2015-2017

Data from ATLAS and CMS at 8 TeV available with breakdown of systematic uncertainties

✤ Included differential top data into NNPDF3.0: constraints on the large-x gluon comparable to those of inclusive jet production Czakon et al 2017

Improved theory uncertainties in regions crucial for BSM
 searches, i.e., m_{tt} > 1 TeV (while fitting only y_t and y_{tt})



The medium-x gluon from NNLO Z $p_{\rm T}$

 $\frac{1}{2}$ Dominated by **quark-gluon scattering**, thus sensitive to the gluon PDF at intermediate values of *x*

Malik and Watt 2013, Boughezal et al 2017

NNLO corrections to the Z p_T also available: **up to 10**% **effects** for a measurement that has **sub-percent exp errors**

Boughezal et al 2015-2017, Gerhmann et al 2015-2017

Complementary information on the gluon as compared to inclusive jets and differential top pair production









The small-x gluon from forward charm production

D and **B** meson production from LHCb allow accessing the **gluon down to** x=10⁻⁶, well below the HERA coverage

PROSA 2015, Gauld et al 2015
Gluon PDF errors reduced by up to a factor 10!

Allows robust estimate for the *prompt* **neutrino flux**, the main background for astrophysical neutrinos at IceCube

 $$\ensuremath{\$}$ Precision calculation of the **UHE neutrino-nucleus cross-section**, with few-percent TH errors up to $E_{\nu}=10^{12}$ GeV





One (upgraded) glue to bind them all

NNPDF3.1 NNLO, Q = 100 GeV



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Quark flavour separation from LHC data

- Recent high-precision W and Z production data from ATLAS, CMS, LHCb and D0 allow a better separation of quark and antiquark flavours
- **LHCb** data in particular provides constrains on the **light quarks at large-**x (*i.e.* error reduction by a factor 2 in NNPDF3.1 for $x \approx 0.1$)
- In the **MMHT (2016)** fit, the new collider $W_{,Z}$ data leads to improved determination of $x(u_V-d_V)$





The inner life of protons



A strange conundrum

In most PDF fits, strangeness suppressed wrt up and down quark sea due to neutrino dimuon data

Solution of the other hand, recent collider data, in particular the ATLAS W,Z 2011 rapidity distributions, prefer instead a **symmetric strange quark sea**

$$R_s(x,Q^2) = \frac{s(x,Q^2) + \bar{s}(x,Q^2)}{\bar{u}(x,Q^2) + \bar{d}(x,Q^2)} \begin{cases} \approx 0.5 \text{ (from neutrino, CMS W+c)} \\ \approx 1.0 \text{ (from ATLAS W,Z)} \end{cases}$$

The new ATLAS data can be accommodated in the **global fits**, and *i*) indeed it **increases strangeness**, but not as much as in a collider-only fit, and *ii*) **some tension remain**s between neutrino and collider data



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How strange is the proton

In the global fit, the resulting strange PDF is the best compromise between the **pulls from individual experiments.** More strange-sensitive measurements are needed to shed more light in this **strange mystery!**



The inner life of protons



How bright is the proton?

- For the calculation of **QED** and electroweak corrections to hadron collider processes requires by consistency to introduce the PDF of the photon in the proton, $\gamma(x,Q)$
- For the first model-independent determination of γ(x,Q) from LHC W,Z data was NNPDF2.3QED, which however affected by large uncertainties, O(100%), due to limited experimental information
- Last year it was shown how to compute y(x,Q) in terms of the well-known inclusive structure functions F₂ and F_L: the resulting photon PDF, LUXqed, exhibits now few-percent uncertainties



 $\mathbf{y}(x,Q)$, ratio to LUXqed

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How bright is the proton?

- The precision determination of $\gamma(x,Q)$ in LUXqed reduces theory uncertainties in several high-mass processes relevant for BSM searches, i.e., **Dilepton** (Drell-Yan) production at the TeV
- Also for **Higgs cross-sections** with EW corrections
- The upcoming NNPDF3.1QED set will be based on the LUX qed master formula for $\gamma(x,Q)$

G. Salam, Oxford colloquium 02/17

pp \rightarrow H W ⁺ (\rightarrow l ⁺ v) + X at 13 TeV						
non-photon induced contributions	91.2 ± 1.8 fb					
photon-induced contribs (NNPDF23)	6.0 ^{+4.4} -2.9 fb					
photon-induced contribs (LUXqed)	4.4 ± 0.1 fb					



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The inner life of protons



A charming story

In global PDF fits there are **two options** to treat the **charm PDF**:

• The charm PDF is **generated perturbatively** via collinear splittings **off gluons and light quarks**:

$$\begin{split} c(x,Q) &= 0 \,, \quad Q < m_c \\ c(x,Q) &\simeq \alpha_s(Q^2) \ln \frac{Q^2}{m_c^2} \int_x^1 \frac{dy}{y} P_{gg}^{(0)}(y) g\left(\frac{y}{x},Q^2\right) \,, \quad Q \geq m_c \end{split}$$

The charm PDF is **fitted from data**, that is, it is treated on an **equal footing to light quarks**:

$$c(x, Q_0) = f(x, \{a_i\}), \qquad Q_0 > m_c$$

Note that the first option is necessarily an **assumption**, which can only be validated by comparing with the results of a **direct fit of the charm PDF**

Moreover fitting the charm PDF offers several advantages:

- **Markov Reduce the dependence of high-scale cross-sections** with respect of the value of the **charm mass**
- **Mathematical Security of Figh-precision collider data** that depend on **quark flavor separation**
- **Compare with models of the non-perturbative (``***intrinsic''***) charm content of the proton**

Fitting the charm PDFs has many important advantages even with for a vanishing non-perturbative component

A charming story

The NNPDF3.1 global analysis is performed with a **fitted charm PDF**. The recent LHC $W_{,Z}$ data, in particular from LHCb, impose stringent constraints on the **size of the charm PDFs at Q=m**_c

$$C(Q^2) \equiv \int_0^1 dx \, \left(xc(x,Q^2) + x\bar{c}(x,Q^2)\right)$$

Charm momentum fraction in the proton



We also find some indications of a **non-perturbative component of the charm PDF localised at large-***x*, though evidence is still limited at the 1.5-sigma level

PDFs and precision physics at the LHC Run II

Compare recent PDF fits with **fiducial W,Z cross-sections at 13 TeV** from ATLAS

Theory calculations with FEWZ and Horace: **NNLO QCD** + **NLO EW**. The latter improves agreement with data

Qualitative agreement between theory and data for most cases - now looking forward to **differential measurements**!





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PDFs and precision physics at the LHC Run II

Compare recent PDF fits for the theory predictions for **Higgs production cross-sections at 13 TeV**

The dominant production mode, gluon-fusion, is **stable wrt the addition of the recent LHC data**

Note improved agreement between **ABMP16 and the global fits** (common value of $\alpha_{S}(m_{Z})=0.118$)

Comparison to be updated with the upcoming new releases from CT and MMHT









Recent progress in nuclear PDFs



Why nuclear PDFs?

Quark and gluons inside bound nucleons exhibit different behaviour than free-nucleon PDFs



So far constraints available mostly from fixed-target DIS, but now LHC is the game-changer!

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Why nuclear PDFs?

- A robust understanding of **cold nuclear matter effects** are of paramount importance as baseline for the **interpretation of heavy-ion collisions**
- Deviations from linear DGLAP evolution (*i.e.* saturation) should be enhanced in nuclei
- The **QCD factorisation theorem** in collisions involving **heavy ions** has been validated only in a restricted kinematical range

Experiment	Process	Observables	Motivation	
ALICE	Heavy quark production	$p_T^D, p_T^{D \to e}$	large-x gluon	The availability of LHC p+Pb data makes
ALICE	Charged jets	$p_T^{\rm ch-jet}$	medium and large-x gluon	possible realising global nuclear PDF fits!
ALICE	Charged hadron production	p_T^h	medium and large-x gluon	
ALICE	Dijet correlations	$k_{Ty} = p_{T,\text{jet}} \sin(\Delta \phi_{\text{dijet}})$	medium and large-x gluon	
ALICE	Inclusive W production	$\mathcal{Y}W \rightarrow l$	quark flavor separation large-x quarks	
ATLAS	Inclusive charged particles	рт	large-x gluon	$\underbrace{9}_{2}$ 10 ⁴
ATLAS	Inclusive W production	$\mathcal{Y}W \rightarrow l$	quark flavor separation medium-x quarks	$\begin{array}{c c} & & & Paukkunen, DIS201\\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & & \\ 10^3 & & & \\ \hline & & & \\ 10^3 & & \\ 10^3 & & \\ 10^3 & & \\ 10^3 & & \\ 10^3 & & \\ 10$
ATLAS	Inclusive Z production	y_Z and p_T^Z	quark flavor separation medium-x quarks	LHC dijets LHC W & Z
CMS	Heavy quark jets	$p_T^{\mathrm{b-jet}}$	medium- <i>x</i> gluon in-medium fragmentation	$10^2 - PHENIX \pi^0$
CMS	<i>B</i> meson production	p_T^B	medium- <i>x</i> gluon in-medium fragmentation	10
CMS	Dijet production	$\eta^{ ext{dijet}}$ and $E_T^{ ext{dijet}}$	medium and large-x gluon	
CMS	Inclusive W,Z production	$y_{W \to l}$ y_Z and p_T^Z	quark flavour separation medium-x quarks	1
LHCb	D meson production	p_T^{D}	small-x gluon	10^{-4} 10^{-3} 10^{-2} x 10^{-1}
LHCb	Inclusive Z production	$\sigma(Z ightarrow l^+ l^-)$	quark flavor separation small-x quarks	I HCP2017, Shanghai, 16/05/2017

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The nuclear PDF landscape

	EPS09	NCTEQ15	EPPS16
Order in α_s	NLO	NLO	NLO
Neutral current DIS <i>ℓ</i> +A/ <i>ℓ</i> +d	\checkmark	\checkmark	\checkmark
Drell-Yan dilepton p+A/p+d	\checkmark	\checkmark	\checkmark
RHIC pions d+Au/p+p	\checkmark	\checkmark	\checkmark
Neutrino-nucleus DIS			\checkmark
LHC p+Pb jet data			\checkmark
LHC p+Pb W, Z data			\checkmark
Drell-Yan dilepton $\pi + A$			\checkmark
Q^2 cut in DIS	$1.3\mathrm{GeV}$	$2{ m GeV}$	$1.3{ m GeV}$
datapoints	929	708	1811
free parameters	15	17	20
error analysis	Hessian	Hessian	Hessian
error tolerance $\Delta\chi^2$	50	35	52
Free proton baseline PDFs	CTEQ6.1	CTEQ6M-LIKE	CT14NLO
Heavy quark treatment	ZM-VFNS	GM-VFNS	GM-VFNS
Flavour separation	none	some	\checkmark
Weight data in χ^2	yes	no	no
Reference	[JHEP 0904 065]	[PR D93 085037]	[EPJ C77, 163]

Paukkunen, DIS2017

Quark nuclear PDFs from W,Z in p+Pb

- W and Z production in p+Pb collisions at the LHC allow disentangling quark and antiquark nuclear PDFs
- Adding LHC data into nuclear PDF fits improves its description, but constraints on nPDFs still moderate
- Looking forward to the 8 TeV p+Pb data!





Gluon nuclear PDF from dijets in p+Pb



The *description of the CMS dijet* data improves sizeably once it is included in the global nPDF fit

Eskola et al 2016

- Dijet production in p+Pb provides information on poorly known gluon PDF
- Available CMS data already allows to discriminate between nPDF sets
- In the EPPS16 analysis this data is central to establish the gluon EMC effects



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nPDFs: the big picture



Despite being based on the widest dataset, EPPS16 exhibits largest uncertainties, showcasing that methodological uncertainties are still a dominant component of nPDFs fits

PDF fits getting closer to proton global fits, but need more LHC p+Pb data

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Summary and outlook

♣ Recent developments in our understanding of the quark/gluon structure of the proton have been driven by a combination of:

Matheory: Progress in NNLO QCD and NLO EW calculations for many collider processes: differential top pairs, inclusive jets, the Z transverse momentum

Data: a wealth of high-precision measurements from ATLAS, CMS and LHCb, in several cases with sub-percent uncertainties.

Methodology: fitted charm PDF, combination / reduction methods for different PDFs,...

Find the second second

The availability of **LHC data on hard probes in p+Pb collision** is a game-changer for nuclear PDF fits: getting closer to global proton PDF fits. Now looking forward to the p+Pb data from Run II.

Summary and outlook

✤ Recent developments in our understanding of the quark/gluon structure of the proton have been driven by a combination of:

