





Parton Distributions and the search for New Physics at the LHC

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The inner life of the protons

- The Large Hadron Collider collides proton, but these are not fundamental particles: really what the LHC is doing is colliding quarks and gluons
- ✓ The distribution of momentum that the quarks and gluons carry is quantified by the Parton Distribution Functions (PDFs), determined by non-perturbative dynamics: cannot be computed from first principles and need to be extracted from experimental data
- An accurate determination of PDFs is of paramount importance to be able to do precision physics at hadronic colliders as the LHC



LHC collisions in a nutshell



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LHC collisions in a nutshell



Drawing by K. Hamilton

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1) PDFs fundamental limit for Higgs boson characterization in terms of couplings



Solid: no TH unc Hatched: with TH unc

ATLAS Simulation Preliminary $\sqrt{s} = 14 \text{ TeV}: \int \text{Ldt} = 300 \text{ fb}^{-1}; \int \text{Ldt} = 3000 \text{ fb}^{-1}$



Massive development of NNLO higher-order calculations ...

Markov ... now we even have the **Higgs gluon fusion xsec at N3LO! Scale uncertainties down to 2**%!



Finally, the computation of the hadronic cross-section relies crucially on the knowledge of the strong coupling constant and the parton densities. After our calculation, the uncertainty coming from these quantities has become dominant. Further progress in the determination of parton densities must be anticipated in the next few years due to the inclusion of LHC data in the global fits and the impressive advances in NNLO computations, improving the theoretical accuracy of many standard candle processes.

Anastasiou et al, arxiv:1503.06056

PDF uncertainties are now **dominant** for a number of crucial LHC processes, and thus it is crucial to match the **accuracy of hard-cross section calculations** with that of the PDFs

2) Very large PDF uncertainties (>100%) for BSM heavy particle production

 $K_{NLO+NLL} = (NLO+NLL)/NLO$

Gluino Pair Production

Squark-Antisquark Pair Production



Beenakker, Borschensky, Kramer, Kulesza, Laenen, Marzani, JR, arXiv:1509.aaaaa

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3) PDFs dominant systematic for precision measurements, like W boson mass, that provide consistency stress-tests of the Standard Model



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The inner life of protons : Parton Distribution Functions



Deep-Inelastic scattering and the discovery of quarks

- Deep-inelastic **lepton-proton scattering**: First evidence for **proton structure**
- Generation Sealering Cross-Section Constant as resolution scale 1/Q decreases.
- Evidence for **point-like constituents in the proton: the quarks**



$$x_{\mathrm{Bj}} = \frac{Q^2}{2 p \cdot q}, \quad Q^2 = -q^2 \quad y = \frac{q \cdot p}{k \cdot p}$$

- If the proton had a different structure, a form factor F(Q) would be expected
- Analogous to **Rutherford's discovery of the point-like atomic nucleus**, while expecting Thomson's Plum model



Deep-Inelastic scattering and the discovery of quarks

- Deep-inelastic **lepton-proton scattering**: First evidence for **proton structure**
- Measured scattering cross-section constant as **resolution scale 1/Q decreases**.
- Evidence for **point-like constituents in the proton: the quarks**



QCD Factorization and PDFs

QCD Factorization Theorem: separate the hadronic cross section into a perturbative, process dependent partonic cross section and non-perturbative, process independent Parton Distributions. In DIS we have:

$$F_i(x,Q^2) = x \sum_i \int_x^1 \frac{dz}{z} C_i\left(\frac{x}{z}, \alpha_s(Q^2)\right) f_i(z,Q^2).$$

Parton-level cross-section



The same Factorization Theorem allows to use the same universal PDFs to provide predict ions for proton-proton collisions at the LHC:

$$\sigma_X(s, M_X^2) = \sum_{a,b} \int_{x_{\min}}^1 dx_1 \, dx_2 \, f_{a/h_1}(x_1, M_X^2) f_{b/h_2}(x_2, M_X^2) \hat{\sigma}_{ab \to X} \left(x_1 x_2 s, M_X^2 \right)$$

Hadron-level cross section

Hadron-level cross section

(2) Parton Distributions

Parton Distribution

Parton-level cross-section

To make sense of LHC collisions, we need first of all to determine the parton distributions of the proton with good precision!



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Parton Distributions

- There is one independent PDF for each parton in the proton: u(x,Q²), d(x,Q²), g(x,Q²), …
- A total of **13 PDFs**, but **heavy quark PDFs generated radiatively** from gluon and light quarks
- At Leading Order, PDFs understood as the **probability of finding a parton of a given flavor that carries a fraction** x of the total proton's momentum
- Once QCD corrections included, PDFs become schemedependent and have no probabilistic interpretation
- Shape and normalization of PDFs are very different for each flavor, reflecting the different underlying **dynamics** that determine each PDF flavor
- QCD imposes valence and momentum sum rules valid to all orders in perturbation theory

Momentum Sum Rule

$$\int_0^1 dx \ x \left[\Sigma(x) + g(x) \right] = 1$$

Valence Sum Rules

$$\int_0^1 dx \ (u(x) - \bar{u}(x)) = 2 \ , \quad \int_0^1 dx \ (d(x) - \bar{d}(x)) = 1$$



PDG Review 2014

Perturbative evolution equations

The dependence of PDFs on Bjorken-x (momentum fraction) is determined by non-perturbative QCD dynamics, but that on the scale Q² (resolution) is instead known from perturbative QCD: the DGLAP evolution equations

$$\frac{\partial q_i(x,Q^2)}{\partial \ln Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dz}{z} P_{ij}\left(z,\alpha_s\left(Q^2\right)\right) q_j\left(\frac{x}{z},Q^2\right)$$

Once *x*-dependence $q(x,Q^2_0)$ extracted from data, pQCD determines PDFs at other scales $q(x,Q^2)$



Experimental data in global PDF fits

A global dataset covering a wide set of hard-scattering observables is required to constrain all possible PDF combinations in the whole range of Bjorken-x

For example, **inclusive jets** are sensitive to the **large-x gluon**, while **HERA neutral current** data pins down the **small-x quarks**

Process	Subprocess	Partons	x range				
$\ell^{\pm} \{p, n\} \to \ell^{\pm} X$	$\gamma^* q \rightarrow q$	q, \overline{q}, g	$x \gtrsim 0.01$				
$\ell^{\pm} n/p \rightarrow \ell^{\pm} X$	$\gamma^* d/u \to d/u$	d/u	$x\gtrsim 0.01$				
$ ho p ightarrow \mu^+ \mu^- X$	$uar{u}, dar{d} ightarrow \gamma^*$	\overline{q}	$0.015 \lesssim x \lesssim 0.35$				
pn/pp $ ightarrow \mu^+\mu^-$ X	$(u\bar{d})/(u\bar{u}) \rightarrow \gamma^*$	\bar{d}/\bar{u}	$0.015 \lesssim x \lesssim 0.35$				
$ u(ar{ u}) N ightarrow \mu^-(\mu^+) X$	$W^* q ightarrow q'$	q, \overline{q}	$0.01 \lesssim x \lesssim 0.5$				
$\nu N ightarrow \mu^- \mu^+ X$	$W^*s \to c$ s		$0.01 \lesssim x \lesssim 0.2$				
$\bar{\nu} N \rightarrow \mu^+ \mu^- X$	$W^* \bar{s} ightarrow ar{c}$	5	$0.01 \lesssim x \lesssim 0.2$				
$e^{\pm} p ightarrow e^{\pm} X$	$\gamma^* q \rightarrow q$	$g,q,ar{q}$	$0.0001 \lesssim x \lesssim 0.1$				
$e^+ p \rightarrow \bar{\nu} X$	$W^+ \left\{ d, s ight\} ightarrow \left\{ u, c ight\}$	d , s	$x \gtrsim 0.01$				
$e^{\pm}p \rightarrow e^{\pm}c\bar{c}X$	$\gamma^* c ightarrow c$, $\gamma^* g ightarrow c ar c$	с, g	$0.0001 \lesssim x \lesssim 0.01$				
$e^{\pm}p ightarrow$ jet $+ X$	$\gamma^*g \rightarrow q\bar{q}$	g	$0.01 \lesssim x \lesssim 0.1$				
$p\bar{p} ightarrow$ jet $+ X$	$gg, qg, qq \rightarrow 2j$	g, q	$0.01 \lesssim x \lesssim 0.5$				
$p\bar{p} \rightarrow (W^{\pm} \rightarrow \ell^{\pm} \nu) X$	$ud \to W, \bar{u}\bar{d} \to W$	$u, d, \overline{u}, \overline{d}$	$x \gtrsim 0.05$				
$p\bar{p} ightarrow (Z ightarrow \ell^+ \ell^-) X$	$uu, dd \rightarrow Z$	d	$x \gtrsim 0.05$				
MSTW08, arXiv:0901.0002							

Global PDF analysis

Parton Distributions, and their associated **uncertainties**, need to be determined from a **global analysis of hard-scattering data** that requires as input the most updated **theory calculations**, precise and varied **experimental data** and robust **statistical methodology**

Theory

- **Martonic cross-sections**
- **M** DGLAP evolution
- Meavy quark treatment
- Fast interfaces to NLO calculations
- MNLO QCD effects
- ☑ QED and electroweak corrections

Data

- **M** Deep-inelastic scattering structure functions
- **M** DIS charm production
- **Mathebasis** DIS Neutrino production
- **I**ets at pp colliders
- **M** Drell-Yan (dilepton pair) production at pp colliders
- 🗹 Top quark data

Methodology

- **M** PDF parametrisation
- PDF uncertainty estimation and error propagation to LHC crosssections
- **Matheoretical uncertainties**
- ☑ Minimisation strategy

Parton Distributions for the LHC

- Higgs cross-sections
- SSM New Physics searches
- Precision Standard Model measurements

Contraction of the local division of the



The Neural Network Approach to Parton Distributions

The NNPDF approach

- The limitations of available PDF sets circa 2005, and the requirements of precision physics at the upcoming LHC, prompted us to develop a completely novel approach to PDF determination
- PDF sets typically based on **restrictive functional forms** leading to strong theoretical bias

$$g(x, Q_0^2) = A_g(1-x)^{m_{\Sigma}} x^{-n_{\Sigma}} \left(1 + a_g \sqrt{x} + b_g x + \ldots\right)$$

Minimum NNPDF solution: use **artificial neural networks** as universal unbiased interpolants

$$g(x, Q_0^2) = A_g(1-x)^{m_{\Sigma}} x^{-n_{\Sigma}} \operatorname{NN}_g(x)$$

PDF sets often rely on the the Gaussian/linear approximation for error estimation and propagation

$$F_0 = F(S_0), \quad \sigma_F = \sqrt{\sum_{i=1}^{N_{\text{par}}} [F(S_i) - F(S_0)]^2}.$$

MODE Solution: Use the **Monte Carlo method** to create a probability distribution in the space of PDFs

$$F_{I,p}^{(\text{art})(k)} = S_{p,N}^{(k)} F_{I,p}^{(\text{exp})} \left(1 + \sum_{l=1}^{N_c} r_{p,l}^{(k)} \sigma_{p,l} + r_p^{(k)} \sigma_{p,s} \right) , \ k = 1, \dots, N_{\text{rep}}$$
Consistent error propagation to LHC xsecs no Gaussian assumptions

 \Im Traditional PDF analyses based on **deterministic minimisation** of the χ^2 to reach convergence in the fit

MACHAE NNPDF solution: Use **Genetic Algorithms** to be able to explore efficiently the vast parameter space

$$\chi^2 = \sum_{i=1}^{N_{\text{dat}}} \sum_{i'=1}^{N_{\text{dat}}} (D_i - T_i) (V^{-1})_{ii'} (D_{i'} - T_{i'}).$$

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Subscription Inspired by biological brain models, Artificial Neural Networks are mathematical algorithms widely used in a wide range of applications, from high energy physics to targeted marketing and finance forecasting



Artificial neural networks aimed to excel in the same domains as their biological counterparts: pattern recognition, forecasting, classification, where our evolution-driven biology outperforms traditional algorithms

Example 1: Marketing. A bank wants to offer a new credit card to their clients. Two possible strategies:

- **Contact all customers**: slow and costly
- Contact 5% of the customers, **train a ANN with their input** (sex, income, loans) and **their output** (yes/no) and use the information to contact only clients likely to accept the offer

Cost-effective method to improve marketing performance



Example 2: **Classification.** Discriminate between signal and background events in complicated final states

- Given S → Find a sequence of the set of the
- Identify automatically the kinematical variables with most discrimination power

Redundancy of **NN-based Multivariate Analysis** guarantees the optimisation of signal/background separation Boosted analysis MVA weights



2

Behr, Bortoletto, Frost, Hartland, Issever, JR, in prep

22

16

14

12

Total associated weight 9 & 0 ¹7

2

m2fj

Pthh

Example 2: Classification. Discriminate between signal and background events in complicated final states

- Improve S/VB as compared to cut-based analyses
- Identify automatically the kinematical variables with most discrimination power

Redundancy of NN-based Multivariate Analysis guarantees the optimisation of signal/background separation

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HH->4b feasibility study



Behr, Bortoletto, Frost, Hartland, Issever, JR, in prep

YZE MHiggsi

Artificial Neural Networks (ANNs) provide universal unbiased interpolants to parametrize PDFs at

low input scales

$$\Sigma(x, Q_0^2) = (1 - x)^{m_{\Sigma}} x^{-n_{\Sigma}} \mathrm{NN}_{\Sigma}(x)$$

$$g(x, Q_0^2) = A_g (1 - x)^{m_{\Sigma}} x^{-n_{\Sigma}} \mathrm{NN}_g(x)$$

Solution Field The ANN class that we adopt are **feed-forward multilayer neural networks** (perceptrons)



Solutional PDF determinations, the input *ansatz* is a simple **polynomial**

$$\Sigma(x, Q_0^2) = (1-x)^{m_{\Sigma}} x^{-n_{\Sigma}} \left(1 + a_{\Sigma} \sqrt{x} + b_{\Sigma} x + \dots \right) ,$$

$$g(x, Q_0^2) = A_g (1-x)^{m_{\Sigma}} x^{-n_{\Sigma}} \left(1 + a_g \sqrt{x} + b_g x + \dots \right)$$

^C The use of Artificial Neural Networks allows:

So theory bias introduced in the PDF determination by the choice of *ad-hoc* functional forms

- The use of very flexible parametrizations for all PDFs regardless of the dataset used. The NNPDF analysis allow for **O(400) free parameters**, to be compared with **O(10-20) in traditional PDFs**
- **Faithful extrapolation**: PDF uncertainties **blow up** in regions with scarce experimental data

PDF Replica Neural Network Learning

The minimisation of the **data vs theory** χ^2 is performed using **Genetic Algorithms** Each **green curve** corresponds to a **gluon PDF Monte Carlo** replica



Artificial Neural Networks vs. Polynomials

• Compare a **benchmark PDF analysis** where **the same dataset** is fitted with **Artificial Neural Networks** and with **standard polynomials** (everything else identical)

ANN avoid biasing the PDFs, **faithful extrapolation at small-x** (very few data, thus error blow up)



Precision tests of the Factorisation Theorem

Perturbative QCD requires that the **momentum integral** should be unity to all orders

$$[M]\left(Q^{2}\right) \equiv \int_{0}^{1} dx \left(xg\left(x, Q^{2}\right) + x\Sigma\left(x, Q^{2}\right)\right)$$

♀ Is it possible to **determine** the value of the momentum integral from the global PDF analysis, rather than **imposing it?** Check in LO*, NLO* and NNLO* fits **without setting M=1**



$$\begin{split} [M]_{\rm LO} &= 1.161 \pm 0.032 \,, \\ [M]_{\rm NLO} &= 1.011 \pm 0.018 \,, \\ [M]_{\rm NNLO} &= 1.002 \pm 0.014 \,. \end{split}$$

Experimental data beautifullyconfirms the pQCD expectation

Extremely non trivial test of the global analysis framework and the **factorization hypotheses**

♀ Very good convergence of the QCD perturbative expansion

PDFs and New Physics searches at the LHC

Higgs Boson Characterisation



The study of the Higgs boson properties is a cornerstone of the LHC program. **All production cross sections** require accurate knowledge of PDFs

- **gg fusion, ttH**: gluon luminosity
- vector-boson fusion: quark-quark luminosity
- **associated production with W/Z**: quark-antiquark luminosity

For many crucial channels, **QCD uncertainties, including PDFs, hatched areas,** limit the accuracy of **Higgs coupling extraction** form LHC data

		σ (8 TeV)	une	certainty
NNLL QCD +NLO EW	gg→H	19.5 pb	14.7%	
	VBF	1.56 pb	2.9%	
NNLO QCD +NLO EW	WH	0.70 pb	3.9%	scale PDF+αs
	ZH	0.39 pb	5.1%	
NLO QCD	ttH	0.13 pb	14.4%	

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Top quarks as gluon luminometers

Free recent NNLO top quark cross section make top data the only LHC observable that is both directly sensitive to the gluon PDF and can be included consistently in a NNLO global analysis

Figure The precise 7 and 8 TeV LHC data can be used to **discriminate between PDF sets** and **to reduce the PDF uncertainties on the poorly known large-x gluon**

The **improved large-x gluon** leads to more accurate theory predictions for **BSM searches**



Czakon, Mangano, Mitov, J.R. 13

Top quark data now included in many global PDF fits: ABM12, MMHT14, HERAfitter

PDFs with QED corrections

QED and **electroweak corrections** are essential for **precision LHC phenomenology**: **W** and **Z production**, **W mass determination**, **WW boson pair** production, **TeV scale** jet and top quark pair production, searches for new **W'**, **Z' bosons**

- Generation of electroweak effects require PDFs with **QED corrections** and a **photon PDF**
- NNPDF2.3 QED: first-ever determination of the photon PDF from LHC data
- Solution Neglecting photon-initiated contributions leads to systematically underestimating theory errors in crucial BSM search channels



Updated squark and gluino cross-sections with threshold-improved PDFs

Bonvini, Marzani, JR, Rottoli, Ubiali, Ball, Bertone, Carrazza, Hartland, arXiv:1507.01006 Beenakker, Borschensky, Kramer, Kulesza, Laenen, Marzani, **JR**, arXiv:1509.aaaaa

Many LHC calculations supplement **NLO and NNLO fixed-order results** with the **resummation of soft threshold logarithms to all orders**: Higgs, top pair production, high-mass supersymmetry



Threshold-improved PDFs can differ substantially wrt fixed-order PDFs: **up to -20% for gg luminosity and -40% for quark-antiquark luminosity,** in the high-mass region relevant for new BSM heavy particles

Marceleter PDF uncertainties in high-mass SUSY cross-sections are very large!

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This is because large-x PDFs are being probed, and these are affected by large errors due to the **lack of direct experimental constraints**



Markon Support How does this affect **existing limits on SUSY production** from LHC searches?

Markov First update NLO+NLL cross-sections with NNPDF3.0NLO



Now include the effect of **NLO+NLL** threshold-improved PDF

Substantial shift, changes qualitatively and quantitatively the behaviour of NLO+NLL SUSY xsecs

Shift within total theory band, so **current exclusion limits unaffected**

W But will become crucial if we ever need to **characterise SUSY particles from LHC data**, much in the same way as we do for the Higgs boson



PDFs at Future Hadron Colliders

Going Beyond: PDFs at a 100 TeV collider



Growing consensus that the next big machine more suitable to explore the energy frontier should be a 100 TeV hadron collider, possibly with also **e+e-** and **ep** operation modes

The **phenomenology of PDFs** at such extreme energies is very rich: top quark PDFs, electroweak effects on PDFs and W/Z boson PDFs, ultra-low-x physics, BFKL dynamics, BSM physics with polarized PDFs,, lots of fun!

First studies being now performed in the context of the CERN FCC working group



Kinematical coverage



The top quark PDF t(x,Q)

✦ At 10 TeV, the top quark PDF **t**(**x**,**Q**) is only a **factor 2 smaller that all other quark PDFs**, with charm and bottom essentially massless flavours there

◆ Shall we include top quarks in any FCC calculation? Yes, but it does not seem particularly useful to introduce a top PDF



The top quark PDF t(x,Q)

- **Meavy-quark PDFs** are constructed **perturbatively** from the resumption of multiple collinear emissions from initial-state light quarks and gluons
- They are only a convenient way to express a **theoretical calculation**, whose usefulness depends on the specific process. No intrinsic top sin the proton!

At the FCC, **resummation of collinear logs** into a top PDFs is not necessary, since these are always small



Electroweak Parton Distributions at 100 TeV

Free analogous of **DGLAP** evolution equations in QCD can be derived in the **electroweak sector** of the Standard Model, but the resulting equations are very different (**Ciafaloni and Comelli 2002**)

Evolution equation for the structure function of W bosons

$$-\frac{\partial}{\partial t} \mathop{\mathcal{F}}_{g}_{AB} = \frac{\alpha_{W}}{2\pi} \left\{ C_{g} \mathop{\mathcal{F}}_{AB} \otimes P_{gg}^{V} + (T_{V}^{C} \mathop{\mathcal{F}}_{g} T_{V}^{C})_{AB} \otimes P_{gg}^{R} + \left(\sum_{L} \operatorname{Tr} \left[t^{B} \mathop{\mathcal{F}}_{L} {}^{t} t^{A} \right] + \sum_{\bar{L}} \operatorname{Tr} \left[t^{A} \mathop{\mathcal{F}}_{L} {}^{t} t^{B} \right] \right) \otimes P_{fg}^{R} + \operatorname{Tr} \left[\mathcal{T}_{L}^{B} \mathop{\mathcal{F}}_{\phi} {}^{t} \mathcal{T}_{L}^{A} \right] \otimes P_{\phi g}^{R} \right\}$$

No numerical implementation of EW evolution equations exist. Very different flavor/coupling structure as compared to QCD evolution equations

Finaddition, **EW must be combined with pure QED evolution**, and then combined with **QCD** into a **complete set of Standard Model PDF evolution equations**

Striking new phenomena: **W and Z PDFs**, internal polarisation of quarks due to chiral couplings, resumption of multiple (massless) weak boson emission

What can we do with polarised protons?

$$\mathcal{L}_{ij} = \frac{1}{s} \int_{\tau}^{1} \frac{dx}{x} \frac{1}{1+\delta_{ij}} \left[q_i\left(x, m_X\right) q_j\left(\frac{\tau}{x}, m_X\right) + q_i\left(\frac{\tau}{x}, m_X\right) q_j\left(x, m_X\right) \right]$$
$$\mathcal{L}_{ij}^L = \frac{1}{s} \int_{\tau}^{1} \frac{dx}{x} \frac{1}{1+\delta_{ij}} \left[\Delta q_i\left(x, m_X\right) q_j\left(\frac{\tau}{x}, m_X\right) + q_i\left(\frac{\tau}{x}, m_X\right) \Delta q_j\left(x, m_X\right) \right]$$

$$A_L = \frac{\sigma^{\uparrow\uparrow} - \sigma^{\downarrow\downarrow} - \sigma^{\uparrow\downarrow} + \sigma^{\downarrow\uparrow}}{\sigma^{\uparrow\uparrow} + \sigma^{\downarrow\downarrow} + \sigma^{\uparrow\downarrow} + \sigma^{\downarrow\uparrow}} \quad A_L =$$

Fuks, Proudom, JR, Schienbein, arXiv:1403.2383



 $rac{\mathcal{L}_{ij}^L}{\mathcal{L}_{ij}}$

Measurement of spin-asymmetries discriminate between **different BSM scenarios** (different production modes) and would allow to **pin down couplings** and branching fractions of

Summary

- Parton Distributions are an essential ingredient for LHC phenomenology
- Accurate PDFs are required for precision SM measurements, Higgs characterisation and New Physics searches
- The determination of **fundamental SM parameters** like the **W mass** or α_s **from LHC data** also greatly benefits from improved PDFs
- PDFs are also a basic component for LO and NLO **Monte Carlo event generators**
- The NNPDF approach provides parton distributions based on a robust, unbiased methodology, the most updated theoretical information and all the relevant hard scattering data including LHC data



Extra Material

PDF Uncertainties: The Monte Carlo Method

Generate a large number of Monte Carlo replicas of the experimental data with the same underlying probability distribution

$$F_{I,p}^{(\operatorname{art})(k)} = S_{p,N}^{(k)} F_{I,p}^{(\exp)} \left(1 + \sum_{l=1}^{N_c} r_{p,l}^{(k)} \sigma_{p,l} + r_p^{(k)} \sigma_{p,s} \right) , \ k = 1, \dots, N_{\operatorname{rep}} >> \mathsf{I}$$

lumi error random numbers

Perform a **PDF determination** on each of these MC replicas

The set of PDF replicas form a representation of the probability density in the space of parton distribution functions

PDF uncertainties can be propagated to physical cross sections using textbook statistics, no need of linear/gaussian assumptions

Central PDF prediction = Expectation
Value of MC sample
$$\langle \mathcal{O} \rangle = \int \mathcal{O}[f] \mathcal{P}(f) Df = \frac{1}{N} \sum_{k=1}^{N} \mathcal{O}[f_k]$$
PDF Uncertainty = Standard
Deviation of MC sample
$$\Delta f = \sqrt{\frac{1}{N} \sum_{k=1}^{N} f_k^2 - \left(\frac{1}{N} \sum_{k=1}^{N} f_k\right)^2}$$

PDF comparisons made easy: APFEL-Web

Comparing different PDF sets is now really easy thanks to the new APFEL-Web online PDF plotter
 Just log in, select the PDF sets that you want to compare, the plotting settings, and have fun!

Welcome to **APFEL** online cluster!

This web-application is a tool designed for High Energy Physics by providing a simple and intuitive interface to plot and compute the most common observables with Parton Distribution Functions (PDFs).

To begin to produce on-line plots, please register and login!

The APFEL library

APFEL, a PDF evolution library, is a computer library specialized in the solution of DGLAP evolution equations up to NNLO in QCD and to LO in QED, both with Pole and $\overline{\rm MS}$ masses. With APFEL you can replace the evolution of LHAPDF sets and check the impact on the choice of evolution parameters. APFEL also computes deep-inelastic scattering processes using multiple schemes.



Web developers: D. Palazzo, S. Carrazza, A. Ferrara APFEL developers: V. Bertone, S. Carrazza, J. Rojo. (Contact)

Bertone, Carrazza, J.R. 13

Parton Distributions with LHC data

A major breakthrough in the recent years has been the inclusion of LHC data into global PDF fits

PDF constraints from a wide variety of LHC processes have been studied, many of which **for first time**



Isolated photon LHC data constraints gluons at medium-x: relevant for Higgs production in gluon fusion

> **W production in association with charm quarks** provides direct access to the proton strangeness





Charm production and the small-x gluon

From The production of **charm and bottom mesons in the forward region** is directly sensitive to the **small-x gluon**, where PDF uncertainties are huge from lack of direct constraints

Using the FONLL calculation and **normalised LHCb 7 TeV D meson data**, we have included these measurements in **NNPDF3.0 NLO** and found a substantial reduction of PDF errors

Semi-analytical FONLL results validated with POWHEG and aMC@NLO calculations

Important implications for the calculations of charm-induced prompt neutrino fluxes at IceCube



Gauld, J.R., Rottoli, Talbert in prep

Jet data in NNLO PDF fits

Jet data at NNLO included using **threshold calculation**, validated by **bin-by-bin comparison with the exact NNLO calculation** in the gg channel (Currie, Gehrmann, Glover, Pires 12-15)

- A substantial amount of data from excluded from NNLO fit at low p_{T,jet} and forward rapidities
- $\stackrel{_{\tiny e}}{_{\scriptstyle =}}$ Perturbative convergence improved if $p_{T,jet}$ used as central scale, as opposed to $p_{T,leading}$
- Validity of threshold calculation extended if **larger values of R used** (0.6 vs 0.4 for ATLAS data)



PDFs and Monte Carlo generators

PDFs are an essential ingredient for the **tuning of soft and semi-hard physics** in LO Monte Carlo event generators like **Pythia8**, **Herwig++** or **Sherpa**

- Most updated tune of **Pythia8**, the **Monash 2013 Tune**, is based on the **NNPDF2.3LO** set
- The harder small-x gluon in NNPDF2.3LO is essential to improve the description of the LHC forward data
- Wext step: specific tunes for NLO Monte Carlo event generators



 NNPDF2.3 is the baseline PDF set in MadGraph5_aMC@NLO and in Pythia8
 NNPDF3.0 is the baseline PDF in Sherpa v2.2.0





Unraveling the gluon polarisation

- Contribution of **gluon polarisation to the proton spin** has been of the **big unknowns** in the last 30 years
- The analysis of recent RHIC polarised jet data in the NNPDFpol and DSSV frameworks has provided first ever evidence for positive (non-zero) polarisation of the gluon in the proton





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AME	RICA	\mathbf{N}^{M}		Search ScientificAmeric	Q		
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More Science » N	lews			- 19	:: 🖂 Email	:: 🚔 Print	

Proton Spin Mystery Gains a New Clue

Physicists long assumed a proton's spin came from its three constituent quarks. New measurements suggest particles called gluons make a significant contribution July 21, 2014 | By Clara Moskowitz

Protons have a constant spin that is an intrinsic particle property like mass or charge. Yet where this spin comes from is such a mystery it's dubbed the "proton spin crisis." Initially physicists thought a proton's spin was the sum of the spins of its three constituent quarks. But a 1987 experiment showed that quarks can account



Total contribution of gluons to proton spin still unknown since **large uncertainties at small-x** from lack of data: need an **Electron-Ion Collider**

Unbiased determination of nuclear PDF

- PDFs of nucleons inside nuclei are modified as compared to free proton PDFs
- Knowledge of nuclear PDFs in lead is important to understand the initial state of pPb and PbPb collisions at the LHC
- Using the NNPDF technology, there is work in progress towards a unbiased nuclear PDF fit, at NNLO and including heavy quark mass effects for the first time
- Allows an essential cross-check for the nuclear models currently used for the LHC heavy ion program



Juan Rojo

Cross section Ratios between 7, 8 and 14 TeV

The staged increase of the LHC beam energy provides a new class of interesting observables: cross section ratios for different beam energies

$$R_{E_2/E_1}(X) \equiv \frac{\sigma(X, E_2)}{\sigma(X, E_1)} \quad R_{E_2/E_1}(X, Y) \equiv \frac{\sigma(X, E_2)/\sigma(Y, E_2)}{\sigma(X, E_1)/\sigma(Y, E_1)}$$

- These ratios can be computed with very high precision due to the large degree of correlation of theoretical uncertainties at different energies
- **Experimentally** these ratios can also be measured accurately since many systematics, like luminosity or jet energy scale, **cancel partially in the ratio**s
- These ratios allow stringent precision tests of the SM, like PDF discrimination



Mangano, J. R., 12

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Particle Physics in the headlines

- **The Higgs Boson** is the most important discovery in particle physics in 25 years
- The Higgs completes the **extremely successful** Standard Model of particle physics, but at the same time opens a number of crucial questions for the field that we need to address
- The LHC will play a crucial role in exploring the energy frontier in the next 20 years

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Queen's Park

El CERN anuncia el descubrimiento de una partícula que podría ser el bosón de Higgs

El CERN anuncia el descubrimiento de una partícula que podría ser el bosón de Higgs, cuya existencia está predicha por el modelo estándar de la física de partículas

Ciencia | 04/07/2012 - 09:46h | Actualizado el 04/07/2012 - 11:27h



Juan Rojo



Example 2: **Marketing.** A bank wants to offer a new credit card to their clients. Two possible strategies:

- **Contact all customers**: slow and costly
- Contact 5% of the customers, train a ANN with their input (sex, income, loans) and their ourput (yes/no) and use the information to contact only clients likely to accepy the offer

Cost-effective method to improve marketing performance

Example 1: **Pattern recognition.** During the Yugoslavian wars, the NATO used ANNs to recognise hidden military vehicles

A military aircraft is identified, despite being hidden below a commercial plane.

Many other applications of ANN in **pattern recognition**: OCR software, hand writing recognition, automated anti-plagiarism software,



PDF fits at NLO+PS accuracy

NLO+PS is current standard for LHC event simulation, and **improves in many directions over fixed-order** NLO results: improved pert. behaviour, direct relation with measured quantities, less need for kin cuts ...

Using NLO+PS calculations in global PDF fits should have many important applications, like for the W mass among others, and is now technically possible thanks to aMCfast, the fast interface to MadGraph5_aMC@NLO based on the applgrid library



aMCfast: Bertone, Frixione, Frederix, J.R., Sutton, arXiv:1406.7693 (for NLO), NLO+PS in preparation

Given the crucial aspect to explore is the role of the PDF used by the MC shower, since this is fixed even in the fast NLO+PS grid

Quite small effect in most
 observables, except extreme
 kinematics like forward
 rapidities

Future NNPDF releasescould be performed atNLO+PS accuracy

aMCfast makes possible to include easily **hadronlevel measurements** directly into PDF fits



Juan Rojo

List from Ian Shipsey

