







News from NNPDF: QED, small-x, and α_S fits

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on behalf of the NNPDF Collaboration



Outline

The NNPDF3.1 global analysis was released in June 2017 (*arxiv:1706.00482, EPJC in press*).

Since then, most of our efforts have been focused on **three spin-offs of the NNPDF3.1 fits**:

MNPDF3.1QED and the photon PDF using the **LUXqed formalism**

Minimum NNPDF3.1 fits with **small-x (BFKL) resummation**

 \mathbf{M} A **determination of** $\alpha_{s}(\mathbf{m}_{z})$ based on the NNPDF3.1 analysis

In this talk we report on the progress in these three topics

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NNPDF3.1QED



Motivation

The NNPDF2.3/3.0QED fits were data-driven determinations of the photon PDF $\gamma(x,Q)$, freely parametrised in terms of an ANN, and the constrained by LHC Drell-Yan measurements

NNPDF 13, Bertone and Carrazza 15

Data-driven QED fits are not competitive anymore with the semi-analytical calculation of the photon PDF using the LUXqed formalism in terms of the inclusive structure functions



Where the structure functions are decomposed in 3 parts:

$$F_2(x,Q) = F_2^{\text{elastic}}(x,Q) + F_2^{\text{inelastic}}(x,Q) + F_2^{\overline{\text{MS}}}(x,Q)$$

It is clearly more advantageous to **perform a QED fit imposing the LUXqed theory constraints** on the photon PDF $\gamma(x,Q)$, rather than extracting it from experimental measurements

Manohar, Nason, Salam, Zanderighi 16, 17

Motivation



Few-percent PDF uncertainties on $\gamma(x,Q)$

Agreement within errors with NNPDF3.0QED 10 NNPDF3.0 3 = 100 GeV 0.3 10⁻² 10⁻³ 10⁻⁵ 10-4 .3 .5 0.1 .7 .9 Х

Even using one of the most sensitive processes to photon-initiated contributions, high-mass DY at 8 TeV, uncertainties in $\gamma(x,Q)$ still at the 30% level



xFitter Developer's Team 17

NNPDF3.1QED: strategy

The **NNPDF3.1QED fits** will impose the **LUXqed formalism** as an external theoretical constrain:



Another important update in the NNPDF3.1QED fits is the use of **NLO QED theory** both in **splitting functions** and in the **DIS coefficient functions**, implemented in the **APFEL code**

Bertone, Carrazza, Rojo 13

Preliminary results



Good agreement between NNPDF3.1QED and the LUXqed photon PDF

The LUXqed systematic uncertainties on **y**(**x**,**Q**) are included as **extra Gaussian fluctuations**

$$\gamma_{\mathrm{unc}}^{k}(x, Q) = \gamma^{k}(x, Q) + \sum_{j=1}^{N_{\mathrm{sys}}} \delta \sigma_{j}^{\mathrm{LUX17}} \mathcal{N}(0, 1)$$

PDF4LHC WG meeting, CERN, 06/08/2017

Preliminary results





Juan Rojo

The iterative procedure converges very fast (within 2 iterations at most)

At the level of quarks and gluons, the differences between the NNPDF3.1 QCD and QCD+QED fits are minimal

PDF4LHC WG meeting, CERN, 06/08/2017

Fits with photon-initiated contributions

Free previous results are based on fits where the **PI contributions are added only to the DIS SFs**

In principle one needs to add them to **all hadronic processes**, but this is very cumbersome

We have checked that NNPDF3.1QED results are stable once **PI contributions added to the LHCb Z production data**, which are directly sensitive to the photon PDF at large *x*



The fits are mostly **insensitive to the inclusion of PI effects in the LHCb cross-sections** Even smaller effects on $\gamma(x,Q)$ would then arise for the rest of the datasets in NNPDF3.1

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Parton Distributions with Small-x Resummation



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Motivation I: beyond fixed-order DGLAP

Perturbative fixed-order QCD calculations have been extremely successful in describing a wealth of data from proton-proton and electron-proton collisions

However, there are theoretical indications that eventually we might need to go beyond DGLAP:

- At very small-x, **logarithmically enhanced terms in 1/x become dominant** and need to be resummed to all orders: **small-x/high-energy/BFKL resummation** formalism.
- The steep rise in the small-x gluon will eventually trigger non-linear recombinations: gluon saturation, BK/JIMWLK equations
- **BFKL resummation** can be matched to the **DGLAP collinear framework**, and thus can in principle be included into a standard PDF analysis

DGLAP
Evolution in Q2
$$\mu^2 \frac{\partial}{\partial \mu^2} f_i(x,\mu^2) = \int_x^1 \frac{dz}{z} P_{ij}\left(\frac{x}{z},\alpha_s(\mu^2)\right) f_j(z,\mu^2),$$
BFKL
Evolution in x $-x \frac{d}{dx} f_+(x,\mu^2) = \int_0^\infty \frac{d\nu^2}{\nu^2} K\left(\frac{\mu^2}{\nu^2},\alpha_s\right) f_+(x,\nu^2)$

Altarelli, Ball, Forte 08 Ciafalini, Colferai, Salam, Stasto 07 Thorne and White 07 + others Within small-*x* resummation, the N^kLO fixed-order DGLAP splitting functions are complemented with the N^hLL*x* contributions from BKFL

$$P_{ij}^{\mathcal{N}^k\mathcal{LO}+\mathcal{N}^h\mathcal{LL}x}(x) = P_{ij}^{\mathcal{N}^k\mathcal{LO}}(x) + \Delta_k P_{ij}^{\mathcal{N}^h\mathcal{LL}x}(x),$$

Motivation II: tensions in inclusive HERA data?

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- Several groups have reported that the **fit quality to the legacy HERA inclusive data** gets worse in the **small**-*x* **and small**-*Q* **region**
- Fypically this trend is **more marked at NNLO**
- Several explanations have been advocated, from higher twists (*i.e.* saturation), issues with the heavy quark schemes, experimental systematics, ...
- What happens if the **PDF fit includes NLL resummation**?







PDFs with BFKL resumation

- Ultimately, the need for (or lack of) BKFL resummation can only be assessed by performing a global PDF analysis with (N)NLO+NLLx matched theory



$$\alpha_{\rm s} = 0.20, \ n_{\rm f} = 4, \ Q_0 \overline{\rm MS}$$

Bonvini, Marzani, Peraro 16 Bonvini, Marzani, Muselli 17

 $P_{ij}^{\mathbf{N}^{k}\mathbf{LO}+\mathbf{N}^{h}\mathbf{LL}x}(x) = P_{ij}^{\mathbf{N}^{k}\mathbf{LO}}(x) + \Delta_{k}P_{ij}^{\mathbf{N}^{h}\mathbf{LL}x}(x),$

https://www.ge.infn.it/~bonvini/hell/

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PDFs with BFKL resumation

Ultimately, the need for (or lack of) BKFL resummation can only be assessed by performing a global PDF analysis with (N)NLO+NLLx matched theory

At small-x, **NNLO+NNL***x* **coefficient functions for DIS structure functions** soften the steep rise of the N3LO result



Bonvini, Marzani, Peraro 16 Bonvini, Marzani, Muselli 17

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Global PDFs with BFKL resummation

- Performed a variant of the NNPDF3.1 global fits using NLO+NLLx and NNLO+NLLx theory
- Small-x resummation of partonic cross-sections included **only for DIS structure functions**: remove all collider data with potential sensitivity to small-x effects

Using NNLO+NLL*x* theory **stabilises small-x gluon**, leading to improved agreement with NLO result



Fit quality at small-x and Q^{2}

In order to assess the impact of small-x resummation for the description of the small-*x* and Q^2 HERA data, compute the χ^2 removing data points in the region where resummation effects are expected



Fit quality at small-x and $Q^{\mbox{\scriptsize 2}}$

Using NNLO+NLL*x* theory, the NNLO instability of the χ^2 disappears

Excellent fit quality to inclusive HERA data achieved in the **entire (x,Q²) region** accessible experimentally



Comparison with HERA data



Using **NNLO+NLL***x* **theory**, improved description of the **small-x NC cross-sections**, in particular of the change of slope

Also **improved description of F**_L, which moreover remains markedly **positive** down to the smallest values of *x* and Q probed

From the LHC to Neutrino Telescopes



Constraints from BFKL resummation



The strong coupling constant α_S(m_Z) from the NNPDF3.1 analysis



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Recap: $\alpha_{s}(m_{z})$ determination from NNPDF2.1

The last determination of $\alpha_S(m_Z)$ within the NNPDF framework was based on the NNPDF2.1 fit NNPDF 10,11

A large number of MC replicas (between 1000 and 100) generated for a range in $\alpha_S(m_Z)$, and then a **parabola was fitted to the** χ^2 **results of the best fit** for each $\alpha_S(m_Z)$ value

Theoretical uncertainties from MHOU estimated using the **Cacciari-Houdeau method**

Result in good agreement with the PDG average and with the MSTW08 determination

 $\alpha_{\rm S}(m_Z) = 0.1173 \pm 0.0007^{\rm PDF} \pm 0.0009^{\rm MHO}$ NNPDF 11







Motivation for an updated $\alpha_{\rm S}(m_{\rm Z})$ fit

Several of the new datasets included in NNPDF3.1 provide a direct handle on $\alpha_s(m_z)$: inclusive jets, top quark pair differential distributions, the Z p_T distribution, ...

Exact NNLO theory is used for the calculations of all collider cross-sections, also for inclusive jet production (removed the LHC 2.76 TeV jet data from the baseline fit since NNLO calculations not yet available)

Finproved methodology for the $\alpha_s(m_z)$ extraction, based on performing parabolic fits replica by replica, keeping the correlations within the generated pseudo-data: this way, $\alpha_s(m_z)$ is determined on exactly the same footing as the PDFs themselves, keeping all PDF- $\alpha_s(m_z)$ correlations

Finite This bypasses the need of using the $\Delta \chi^2$ =1 criterion (or any other value for the tolerance), which might not be adequate in the presence of **inconsistent experiments**

Fitting strategy

- Fully correlated pseudo-data replicas are generated, and then fits with a range of $\alpha_S(m_Z)$ values are performed for each MC replica
- A parabolic fit determines the **preferred value of** $\alpha_s(m_z)$ **for each replica** (points that fail to satisfy quality criteria are removed, *i.e.*, fits that have not converged)
- From The mean and variance of the preferred $\alpha_s(m_z)$ over the replica sample determines the central value of $\alpha_s(m_z)$ and the associated 68% CL PDF uncertainty



Fitting strategy

Some processes (jets, top, FT NC DIS, ...) lead to **marked parabolas**, illustrating **sensitivity to** $\alpha_s(m_z)$ Other processes, such as LHC Drell-Yan, are **less sensitive**, except for very small values of $\alpha_s(m_z)$







Preliminary results

Using the **correlated replica method**, we find that the NNPDF3.1 global analysis leads to:

 $\alpha_{\rm S}(m_{Z}, \text{NNLO}) = 0.1190 \pm 0.0005^{\rm PDF}$ $\alpha_{\rm S}(m_{Z}, \text{NLO}) = 0.1219 \pm 0.0007^{\rm PDF}$

 $\stackrel{\scriptstyle \odot}{}$ This is to be compared with the results obtained using the central fits and the $\Delta \chi^2$ =1 criterion:

 $\alpha_{\rm S}(m_Z, \text{NNLO}) = 0.1183 \pm 0.0004^{\rm PDF}$ $\alpha_{\rm S}(m_Z, \text{NLO}) = 0.1214 \pm 0.0004^{\rm PDF}$

For the best-fit α_s(m_z) value is driven by the jet, top, and FT NC DIS data. The higher value of α_s(m_z) as compared to the NNPDF2.1 fit is partly due to the LHC DY and Z p_T data.



Correlated replica method vs " $\Delta \chi^2 = 1$ " method



In general there is reasonable consistency between the two methods, highlighting the robustness of the results

The " $\Delta \chi^2 = 1$ " method can underestimate the real PDF uncertainties, by a small amount in some cases (total dataset, jets, top) and by a large amount for those processes with little sensitivity to $\alpha_S(m_Z)$, such as HERA and FT Drell-Yan

Preliminary results



Good agreement with the PDG17 average, and consistent with the MMHT2014 determination

Also repeated the analysis based on a **collider-only dataset:** not competitive with the global fit, presumably because the constraints from the **fixed-target DIS data still important**

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\alpha_{s}(m_{Z}, NNLO, global) = 0.1190 \pm 0.0005^{PDF}
\alpha_{s}(m_{Z}, NNLO, collider-only) = 0.1229 \pm 0.0026^{PDF}
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Summary and outlook

MNPDF3.1QED and the photon PDF using the LUXqed formalism

Preliminary results presented, **final NNPDF3.1QED NLO and NNLO close to completion**. Good agreement with LUXqed, no changes for quarks and gluons

Release of the **public library** *FiatLUX* that implements the LUXqed formalism

Study of phenomenological implications at the LHC for PI processes

☑ NNPDF3.1 fits with **small-x (BFKL) resummation**

Using NNLO+NLL*x* theory improves the perturbative expansion at small-x, cures the chi2 instability, and the allows a better description of the **inclusive HERA data**

Figure For **UHE astrophysics**, as well as for f**uture colliders (LHeC, FCC-hh/eh)**

 \mathbf{M} A **determination of** $\alpha_{s}(\mathbf{m}_{z})$ based on the NNPDF3.1 analysis

New approach developed, based on correlated MC replicas

[§] The new fit result, $\alpha_s(m_z, NNLO) = 0.1190 \pm 0.0005^{PDF}$, is consistent with the PDG average but **higher than other PDF-based determinations**, mostly due to the pull of the recent high-precision LHC data. Collider-only determinations still not competitive

Solution of the second second