#### The PDF-Lattice Study:

#### Connecting PDFs from phenomenology and lattice QCD

Smowmass EF06 meeting

Emanuele R. Nocera – Nikhef

July 1, 2020



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# Foreword: (collinear) leading twist PDF map $f_1 =$ parton $g_1 =$ hadron p s $h_1 =$ $\phi_{ij}(k;p,s) = 2\pi \sum_{\mathbf{x}} \int \frac{d^3 \mathbf{P}_X}{2E_{\mathbf{x}}} \delta^4(p-k-P_X) \langle p, s | \bar{\psi}_j(0) | X \rangle \langle X | \psi_i(0) | p, s \rangle$ $\phi(x,s) = \frac{1}{2} \left[ \mathbf{f_1}(x) \mathbf{/}_+ + s_L g_1(x) \gamma^5 \mathbf{/}_+ + \frac{h_1 i \sigma_{\mu\nu} \gamma^5 n_+^{\mu} s_T^{\nu}}{h_+^{\nu} \sigma_+^{\nu} n_+^{\nu} n$ In this talk $\mathbf{f_1} \to f$ , $q_1 \to \Delta f$ and $h_1 \to \delta f$ $f(x) = \frac{1}{4\pi} \int dy^- e^{-ixp^+y^-} \langle p, s | \bar{\psi}_f(0,0,\mathbf{0}_\perp) \gamma^+ \mathcal{P}\psi_f(0,y^-,\mathbf{0}_\perp) | p, s \rangle$ $\Delta f(x) = \frac{1}{4\pi} \int dy^- e^{-ixp^+y^-} \langle p, s | \bar{\psi}_f(0,0,\mathbf{0}_\perp) \gamma^+ \gamma^5 \mathcal{P} \psi_f(0,y^-,\mathbf{0}_\perp) | p, s \rangle$ $\delta f(x) = \frac{1}{4\pi} \int dy^- e^{-ixp^+y^-} \langle p, s | \bar{\psi}_f(0,0,\mathbf{0}_\perp) i \sigma^{1+} \gamma^5 \mathcal{P} \psi_f(0,y^-,\mathbf{0}_\perp) | p, s \rangle$

#### Strategy to fit PDFs from data

Collinear, leading-twist factorisation of physical observables

$$\mathcal{O}_I = \sum_{f=q,\bar{q},g} C_{If}(y,\alpha_s(\mu^2)) \otimes f(y,\mu^2) + \text{p.s. corrections} \qquad f \otimes g = \int_x^1 \frac{dy}{y} f\left(\frac{x}{y}\right) g(y)$$

2 Parametrisation: general, smooth, flexible at an initial scale  $Q_0^2$ 

$$\begin{split} xf_i(x,Q_0^2) &= A_{f_i} \, x^{a_{f_i}} \, (1-x)^{b_{f_i}} \, \mathscr{F}(x,\{c_{f_i}\}) \\ xf_i(x,Q^2) \xrightarrow{x \to 0} x^{a_{f_i}} & \xrightarrow{\mathscr{F}(x,\{c_{f_i}\}) \xrightarrow{x \to 0} \text{finite}} \\ \xrightarrow{\text{smooth interpolation in between}} & xf_i(x,Q^2) \xrightarrow{x \to 1} (1-x)^{b_{f_i}} \end{split}$$

A prescription to determine/compute expectation values and uncertainties

$$\begin{split} \chi^2 &= \sum_{i,j}^{N_{\rm dat}} [T_i[\{\vec{a}\}] - D_i](\mathrm{cov}^{-1})_{ij}[T_j[\{\vec{a}\}] - D_j] \\ E[\mathcal{O}] &= \int \mathcal{D}f\mathcal{P}(f|data)\mathcal{O}(f) \qquad V[\mathcal{O}] = \int \mathcal{D}f\mathcal{P}(f|data)[\mathcal{O}(f) - E[\mathcal{O}]]^2 \\ \text{Monte Carlo: } \mathcal{P}(f|data) &\longrightarrow \{f_k\} \qquad \qquad \text{Maximum likelihood: } \mathcal{P}(f|data) &\longrightarrow f_0 \\ E[\mathcal{O}] &\approx \frac{1}{N} \sum_k \mathcal{O}(f_k) \qquad \qquad E[\mathcal{O}] \approx \mathcal{O}(f_0) \\ V[\mathcal{O}] &\approx \frac{1}{N} \sum_k [\mathcal{O}(f_k) - E[\mathcal{O}]]^2 \qquad \qquad V[\mathcal{O}] \approx \text{Hessian, } \Delta\chi^2 \text{envelope, } \dots \end{split}$$

Experimental, theoretical and procedural uncertainties

More details in Ann.Rev.Nucl.Part.Sci. (2020) 70

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Connecting PDFs from pheno and lattice

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### Strategies to reconstruct PDFs from lattice QCD

Hadronic tensor [PRL 72 (1994) 1790] Auxiliary scalar quarks [PLB 441 (1998) 371] Ficticious heavy quark [PRD 73 (2006) 014501] Higher moments [PRD 86 (2012) 054505] Quasi-PDFs (LaMET) [PRL 110 (2013) 262002] Good Cross Sections [PRL 120 (2018) 022003] Compton Amplitudes [PRL 118 (2017) 242001] Pseudo-PDFs [PRD 96 (2017) 034025]



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#### Strategies to reconstruct PDFs from lattice QCD

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$$\begin{split} h(z,p_z) &= \frac{1}{4p_\alpha} \sum_{s=1}^2 \langle p,s | \, \bar{\psi}(z) \gamma_\alpha e^{ig \int_0^z A_z(z')dz'} \psi(0) \, | p,s \rangle \\ & \underline{\text{Quasi-PDFs}} \\ \tilde{q}(x,\Lambda,p_z) &= \int \frac{dz}{2\pi} e^{-ixzp_z} p_z h(z,p_z) \\ \tilde{q}(x,\Lambda,p_z) &= \int_{-1}^1 \frac{dy}{|y|} Z\left(\frac{x}{y},\frac{\mu}{p_z},\frac{\Lambda}{p_z}\right)_{\mu^2 = Q^2} q(y,Q^2) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{p_z^2},\frac{M^2}{p_z^2}\right) \\ & \underline{\text{Pseudo-PDFs}} \\ \mathcal{P}(x,z^2) &= \int \frac{d\nu}{2\pi} e^{-ix\nu} \overline{h}(\nu,z^2) \quad \overline{h}(\nu,z^2) \equiv h(z,p_z) \quad \mathcal{M}(\nu,z^2) = \frac{\overline{h}(\nu,z^2)}{\overline{h}(0,z^2)} \\ & q(x,\mu^2) = \int \frac{d\nu}{2\pi} e^{-ix\nu} \mathcal{M}(\nu,z^2) + \mathcal{O}(z^2) \end{split}$$

Each step is associated to systematic uncertainties and theoretical challenges

More details in Adv. High Energy Phys. 2019 (2019) 3036904

### Kinematic coverage



Fits of *f* from **thousands** of data CT, MMHT, NNPDF, .... Fits of  $\Delta f$ from hundreds of data DSSV, JAM, NNPDF, ... Fits of  $\delta f$ from tens of data Kang; Anselmino; Bacchetta

Obvious spread in the distribution of measurements across PDF species

#### Connecting two faces of the same world



Define a mutually agreed conventional notation for relevant PDF-related quantities, such as PDF moments.

Assess the sources of systematic uncertainties in lattice-QCD calculations.

Identify a best-set of quantities to benchmark lattice-QCD calculations against global-fit determinations.

Set precision targets for lattice-QCD calculations with respect to global-fit determinations.

Assess the impact of lattice-QCD calculations on global-fit determinations within their current/projected precision.

#### PDFLattice2017, Balliol College, Oxford, 22-24 March 2017

#### PDFLattice2019, Kellogg Biological Station, Hickory Corners, 25-27 September 2019

Prog.Part.Nucl.Phys. 100 (2018) 107; arXiv:2006.08636

# 1. Appraising lattice QCD

#### Define a quantitative benchmark for PDF moments Benchmark quantities

$$\begin{split} \mathbf{f} & \langle x \rangle_{u^{+}-d^{+}} = \int_{0}^{1} dx \, x \begin{bmatrix} u^{+}(x,Q^{2}) - d^{+}(x,Q^{2}) \end{bmatrix} \\ \langle x \rangle_{q^{+}} = \int_{0}^{1} dx \, xq^{+}(x,Q^{2}), \ q = u, d, s \quad \langle x \rangle_{g} = \int_{0}^{1} dx \, xg(x,Q^{2}) \\ \Delta f & g_{A} = \langle 1 \rangle_{\Delta u^{+}-\Delta d^{+}} = \int_{0}^{1} dx \left[ \Delta u^{+}(x,Q^{2}) - \Delta d^{+}(x,Q^{2}) \right] \\ \langle 1 \rangle_{\Delta q^{+}} = \int_{0}^{1} dx \, \Delta q^{+}(x,Q^{2}), \ q = u, d, s \quad \langle x \rangle_{\Delta u^{-}-\Delta d^{-}} = \int_{0}^{1} x dx \left[ \Delta u^{-}(x,Q^{2}) - \Delta d^{-}(x,Q^{2}) \right] \\ \delta f & g_{T} = \langle 1 \rangle_{\delta u^{-}-\delta d^{-}} = \int_{0}^{1} dx \left[ h_{1}^{u^{-}}(x,Q^{2}) - h_{1}^{d^{-}}(x,Q^{2}) \right] \\ g_{T}^{q} = \langle 1 \rangle_{\delta q^{-}} = \int_{0}^{1} dx \, h_{1}^{q^{-}}(x,Q^{2}), \ q = u, d, s \end{split}$$

#### Benchmark criteria

	*	0	
discretisation	$ \begin{cases} a_1, \dots, a_i, \dots \} & i \ge 3 \\ a_l, a_m < 0.1 \text{ fm } \left( \frac{a_{\max}}{a_{\min}} \right)^2 \ge 2 \end{cases} $	$ \begin{array}{l} \{a_1, \ldots, a_i, \ldots\}  i \geq 2 \\ a_l < 0.1 \; \mathrm{fm}  \left(\frac{a_{\max}}{a_{\min}}\right)^2 \geq 1.4 \end{array} $	otherwise
chiral extrapolation	$m_{\pi,i},i\geq 3$ $m_{\pi,1,2}<250\;{ m MeV}\;m_{\pi,3}<200\;{ m MeV}$	$m_{\pi,i}, i \geq 3$ $m_{\pi,1,2} < 300 \; { m MeV}$	otherwise
finite volume	$\begin{split} m_{\pi,\min}L &\geq 4\\ L_1 \neq L_2 \neq L_3 > 2.5 \text{ fm} \end{split}$	$\begin{array}{l} m_{\pi,\min}L\geq 3.4\\ L_1\neq L_2>2.5 \ \mathrm{fm} \end{array}$	otherwise
renormalisation	non-perturbative (RI-MOM)	perturbative (one-loop or ohigher)	otherwise
excited states	$(source-sink)_i$ $i \ge 3 \ \forall \ m_{\pi}, L$	$( ext{source-sink})_i \ i \geq 2 \ orall \ m_\pi, L$	otherwise
Emanuelo P. Noco	Connecting RDEc from phone	a and lattice	1. 1 2020

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### Moments of f



Moment		Lattice QCD		Global Fit
$\left\langle x \right\rangle_{u^+ - d^+}$	0.153 — 0.194	0.111 — 0.209	0.166 — 0.212	0.161(18)
$\langle x \rangle_{u^+}$	0.359(30)†	0.307(35) <sup>†</sup>	_	0.353(12)
$\langle x \rangle_{d^+}$	0.188(19)†	0.160(48) <sup>†</sup>	_	0.192(6)
$\langle x \rangle_{s+}$	0.052(12)†	0.051(26) <sup>†</sup>	_	0.037(3)
$\langle x \rangle_g$	0.427(92)†	0.353 — 0.587	_	0.411(8)
+				

<sup>†</sup> Single lattice result

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#### Moments of $\Delta f$



Moment	Lattice QCD Global Fit				
$g_A$	1.179 — 1.309	1.202 — 1.314	$1.268(36)^{\dagger}$	1.258(28)	
$\left< 1 \right>_{\Delta u} +$	0.738 — 0.875	0.810 — 1.001	_	0.813(25)	
$\left< 1 \right>_{\Delta d} +$	-0.473 — -0.403	-0.431 — -0.278	—	-0.462(29)	
$\left< 1 \right>_{\Delta s} +$	-0.0538 — -0.0379	-0.0035(9)†	_	-0.114(43)	
$\langle x \rangle_{\Delta u^ \Delta d^-}$	0.174 — 0.239	$0.221(^{+27}_{-25})^{\dagger}$	_	0.199(16)	
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<sup>†</sup> Single lattice result

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#### Moments of $\delta f$



Moment		Lattice QCD		Global Fit
$g_T$	0.894 — 1.023	0.909 — 1.175	0.941 — 1.039	0.10 — 1.1
$\left< 1 \right>_{\delta u} -$	0.688 — 0.814	0.85(8)	0.782(21)	-0.14 — 0.91
$\left< 1 \right>_{\delta d} -$	-0.221 — -0.189	-0.24(3)†	-0.219(17)†	-0.97 — 0.47
$\left<1\right>_{\delta s}-$	-0.0085 — 0.0031	-0.012(18)†	-0.00319(72)†	_
A				

<sup>†</sup> Single lattice result

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#### Qualitative comparison of lattice QCD and global PDF fits Unpolarised PDFs



ETMC and LP3 determinations are both at the physical pion mass Lattice determinations are qualitatively similar (among them) and similar to global fits Nucleon momentum is limited by lattice spacing Different procedures lead to slightly different behaviour in x

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# Qualitative comparison of lattice QCD and global PDF fits Helicity PDFs



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# 2. Exploiting lattice QCD

### Impact of lattice QCD moments on f

Generate lattice QCD pseudodata assuming NNPDF3.1 central values for  $\langle x \rangle_{u^+}$ ,  $\langle x \rangle_{d^+}$ ,  $\langle x \rangle_{s^+}$ ,  $\langle x \rangle_g$ ,  $\langle x \rangle_{u^+-d^+}$ 

scenario	$\left\langle x\right\rangle _{u}+$	$\left\langle x\right\rangle _{d}+$	$\left\langle x\right\rangle _{s}+$	$\langle x\rangle_g$	$\left\langle x\right\rangle _{u^{+}-d^{+}}$
A	3%	3%	5%	3%	5%
C	2% 1%	2% 1%	4% 3%	2% 1%	4% 3%
current	17%	30%	45%	13%	60%

Assume percentage uncertainties according to three scenarios

Reweight NNPDF3.1 with lattice pseudodata and look at the impact



#### Impact of lattice QCD moments on $\Delta f$

Generate lattice QCD pseudodata assuming NNPDFpol1.1 central values for

$$g_A \equiv \langle 1 \rangle_{\Delta u^+ - \Delta d^+}$$
,  $\langle 1 \rangle_{\Delta u^+}$ ,  $\langle 1 \rangle_{\Delta d^+}$ ,  $\langle 1 \rangle_{\Delta s^+}$ ,  $\langle x \rangle_{\Delta u^- - \Delta d^-}$ 

scenario	$g_A$	$\left< 1 \right>_{\Delta u} +$	$\left<1\right>_{\Delta d}+$	$\left<1\right>_{\Delta s}+$	$\left\langle x\right\rangle _{\Delta u^{-}-\Delta d^{-}}$
A B C	5% 3% 1%	5% 3% 1%	10% 5% 2%	100% 50% 20%	70% 30% 15%
current	3%	3%	5%	70%	65%

Assume percentage uncertainties according to three scenarios

Reweight NNPDFpol1.1 with lattice pseudodata and look at the impact



### Impact of lattice QCD moments on $\delta f$

Simultaneous fit to the Collins asymmetry data from HERMES and COMPASS of

$$f_1^q(x,k_{\perp}^2) = h_1^q(x,k_{\perp}^2) = D_1^{h/q}(z,p_{\perp}^2) = H_1^{\perp h/q}(z,p_{\perp})$$

and to three lattice *data sets* with an estimate of systematic uncertainties PDNME [Bhattacharya et al. (2016)] RQCD [Bali et al. (2015)] LHPC [Green et al. (2012)] using Monte Carlo techniques for the representation of uncertainties



PRL 120 (2018) 152502

Excellent description of the data with and without lattice results ( $\chi^2/N_{dat} = 0.65$ ) Lattice results seem compatible with measured asymmetries Lattice results are able to reduce the uncertainty on  $h_1$  and  $H_1^{\perp}$  significantly

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Connecting PDFs from pheno and lattice

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#### Impact of lattice QCD moments on $\delta f$ [Courtesy of M. Radici]



[1] ETMC 19; [2] Mainz 19; [3] LHPC 19; [4] JLQCD 18; [5] PNDME 18; [6] ETMC 17; [7] RQCD 14; [8] LHPC 12 global fit [Radici, in progress]; JAM [PRL 120 (2018) 152502]; TMD [PRD 93 (2016) 014009]; Torino [PRD 92 (2015) 114023]



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#### Reconstructing PDFs from lattice moments



Detmold et al. [EPJ direct 3 (2001) 13]

 $u\,-\,d$  from the lowest few lattice moments, ensure the correct behavior in the chiral and heavy quark limits

Haegler et al. [PRD 77 (2008) 094502]

non-perturbative renormalization factor for the axial vector current, only connected diagrams are included

Bacchetta et al. [PRD 95 (2017) 014036]

supplement lattice moments with quasi-PDFs (using results of a diquark spectator model) matched at a fixed point  $x_0$ 



0.2

0.4

х

0.6

0.8

0.6

0.4

0.2

0

0

#### Impact of lattice calculations of x-space PDFs

Apply Bayesian reweighting to the isotriplet PDF combinations

 $\begin{array}{ll} f & {\sf NNPDF3.1} & u(x_i,Q^2) - d(x_i,Q^2) & \bar{u}(x_i,Q^2) - \bar{d}(x_i,Q^2) \\ \Delta f & {\sf NNPDFpol1.1} & \Delta u(x_i,Q^2) - \Delta d(x_i,Q^2) & \Delta \bar{u}(x_i,Q^2) - \Delta \bar{d}(x_i,Q^2) \end{array} i = 1, \ldots, N_x$ 

Consider uncorrelated lattice pseudodata  $Q^2 = 4 \text{ GeV}^2$  and  $x_i = 0.70, 0.75, 0.80, 0.85, 0.90$  for three scenarios: (D)  $\delta_L^{(i)} = 12\%$ ; (E)  $\delta_L^{(i)} = 6\%$ ; (F)  $\delta_L^{(i)} = 3\%$ 



No large differences among the three scenarios (PDF variations are correlated) Moderate precision required for lattice QCD to make an impact on antiquarks at large xCaveat: rough assumptions ( $\{x_i\}, \delta_L^{(i)}$ ); can we do something better?

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Quasi-PDFs defined as momentum-dependent nonlocal static matrix elements for nucleon states at finite momentum, with an ultraviolet cut-off scale  $\Lambda\sim 1/a$ 

$$\widetilde{q}(x,\Lambda,p_z) = \int \frac{dz}{4\pi} e^{-ixzp_z} \frac{1}{2} \sum_{s=1}^2 \langle p,s | \, \bar{\psi}(z) \gamma_\alpha e^{ig \int_0^z A_z(z')dz'} \psi(0) \, | p,s \rangle$$

Must be related to the corresponding light-front PDF, usually within LaMET

$$\widetilde{q}(x,\Lambda,p_z) = \int_{-1}^1 \frac{dy}{|y|} Z\left(\frac{x}{y},\frac{\mu}{p_z},\frac{\Lambda}{p_z}\right)_{\mu^2 = Q^2} q(y,Q^2) + \mathcal{O}\left(\frac{\Lambda_{\mathsf{QCD}}^2}{p_z^2},\frac{m^2}{p_z^2}\right)$$

Restrict to the isotriplet distributions and consider ETMC lattice data  $V_3 = u - \bar{u} - [d - \bar{d}]$  $T_3 = u + \bar{u} - [d + d]$  $\mathcal{O}^{\mathrm{Re}}_{\gamma^0}(z,\mu) \equiv \mathrm{Re}[h_{\gamma^0,3}(zp_z,z^2,\mu^2)] = \mathcal{C}^{\mathrm{Re}}_3 \circledast V_3 \quad \mathcal{O}^{\mathrm{Im}}_{\gamma^0}(z,\mu) \equiv \mathrm{Im}[h_{\gamma^0,3}(zp_z,z^2,\mu^2)] = \mathcal{C}^{\mathrm{Im}}_3 \circledast T_3$ 0.75 0.50 0.5 0.25 0.00 0.0 -0.25-0.50-0.5 •  $\gamma_0$  Dirac structure non-perturbative renormalization -0.75-1.0 • simulation at the physical pion mass -1.00•  $P_z = 10\pi/L (1.38 \text{GeV})$ -1.258 10 12 14 Ó 2 à 10 12 14 2 8 z 7

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Scenario	Cut-off	FVE	Excited states	Truncation
S1	10%	2.5%	5%	10%
S2	20%	5%	10%	20%
S3	30%	$e^{-3+0.062z/a}\%$	15%	30%
S4	0.1	0.025	0.05	0.1
S5	0.2	0.05	0.1	0.2
S6	0.3	$e^{-3+0.062z/a}$	0.15	0.3

Consider various scenarios for systematic uncertainties

Percentage values for scenarios S1-S3 should be understood as a given fraction of the central value of the matrix element Absolute values for scenarios S4-S6 are shifts independent from the matrix element



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# 3.Outlook and conclusions

#### Towards 3D structure



Spectacular theoretical effort on both the lattice and the phenomenological sides Case for the physics of new facilities and/or upgrades (JLab-12, RHIC, after@LHC, EIC)

#### Towards 3D structure - Form Factors



Neutron Electric and Magnetic



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### Towards 3D structure - GPDs





Determining GPDs is generally more challenging than determining PDFs More variables: Wilson line z, hadron momentum  $P_z$ , momentum transfer t, skewness  $\xi$ Perturbative matching depends on skewness, but not on momentum transfer Phenomenological fits complicate because of a general lack of data (but significant progress ongoing, see *e.g.* PARTONS framework [EPJC78 (2018)478])

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#### Towards 3D structure - TMDs

	quark polarization							
auon		U	L	Т				
חמווב	U	f1		$h_1^{\perp}$				
	L		<b>g</b> 1∟	$h_{1L}^{\perp}$				
ווחרופ	Т	$f_{1T}{}^{\!\!\perp}$	<b>g</b> 1T	$h_1 h_{1T}^{\perp}$				

	Framework	HERMES	COMPASS	DY	Z production	N of points	χ²/N <sub>points</sub>
Pavia 2017 arXiv:1703.10157	NLL	~	~	~	~	8059	1.55
SV 2017 arXiv:1706.01473	NNLL'	×	×	r	r	309	1.23
BSV 2019 arXiv:1902.08474	NNLL'	×	×	~	~	457	1.17
SV 2019 arXiv:1912.06532	NNLL'	r	r	r	r	1039	1.06
Pavia 2019 arXiv:1912.07550	NºLL	×	×	r	r	353	1.02



Phenomenology: various challenges (data, theoretical framework)

Lattice: Sivers and Boer-Mulders shifts; Collins-Soper kernel

### Summary

There has been an undeniable progress in the determination of PDFs from both the global fit to data and the lattice QCD sides

Such a progress cannot be ignored

Opportunity to gain further knowledge by improving cross-talk between the two sides

Attempt to realise such an opportunity within the PDFLattice joint effort benchmark + impact studies

Some substantial effort is ongoing

#### the definition and renormalisation of the non-local operators involved in the lattice simulation

the proof of the factorization theorem between PDFs and quasi-PDFs

the computation of the matching coefficients

relating lattice-computable quantities to PDFs in different renormalization schemes

the implementation of efficient methods

to incorporate lattice QCD information into global PDF fit determinations (and viceversa)

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## Thank you

Emanuele R. Nocera