

PineAPPL: towards NLO EW in NNPDF

NLO EW corrections (not only) for PDF fits
arXiv:2008.12789

Christopher Schwan

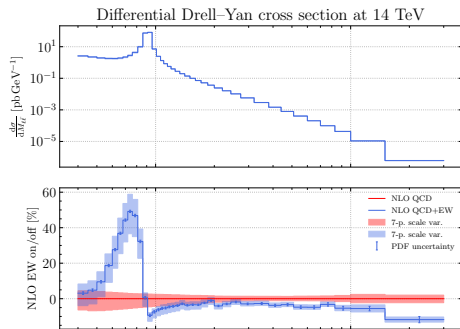
DIS 2021, 13 March 2021



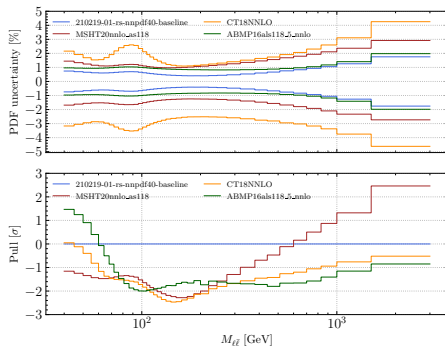
This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 740006

NLO EW PDF fit with NNPDF: motivation

- NNPDF4.0—see Emanuele's talk yesterday—aims for $\lesssim 1\%$ accuracy/precision
- Further improvement need previously neglected contributions: theory uncertainties, **NLO EW**, ...
- Has never been done in a global (all processes) and consistent way (correct data)



- Can include extreme phase-space regions that show large EW corrections: large $M_{\ell\ell}$ in Drell–Yan, large p_T of Z boson, ... \rightarrow constrain high x
- Those regions are going to be measured more precisely in the future!



NLO EW in a PDF fit: What has to be done?

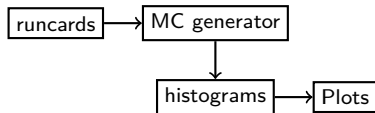
- ✓ Theory: Need corrections in the form of **interpolation grids**: PineAPPL
 - interfaced with MG5_aMC@NLO, see [S. Carrazza, E.R. Nocera, C.S., M. Zaro]; now released in v3.1.0:
<https://launchpad.net/mg5amcnlo>
 - WIP: SHERPA/MCgrid [E. Bothmann et al.]
 - API available in C, C++, Fortran, Python, Rust
- ✓ Data: Needs careful selection
 - no subtraction of FSR
 - no photon-initiated subtraction
 - proper observable definition
- ✓ Write/test runcards for all PDF processes and ...
- ✗ Run them (WIP)
- ✗ Implement changed data (WIP)
- Run fit

What is PineAPPL?

PineAPPL

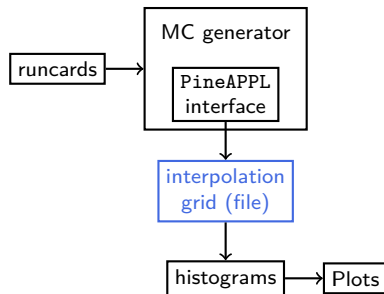
A tool/library for storing **PDF-independent theoretical predictions** in interpolation grids [S. Carazza, E.R. Nocera, C.S., M. Zaro]

Typical setup:



- 1 write runcards for your process/distributions
- 2 MC generates histograms **with PDFs baked-in**
- 3 changing PDFs is **slow**—need to rerun everything

With PineAPPL:



- grids can be convolved with **arbitrary PDF sets** in a **matter of seconds**

Interpolation techniques and applications

Not a new idea:

- APPLgrid [T. Carli et al.] or
- fastNLO [T. Kluge, K. Rabbertz, M. Wobisch]

Interpolate PDFs $f_a(x)$ with kernels $L_i(x)$, $x \mapsto (0, 1)$:

$$f_a(x) = \sum_{i=1}^{\infty} f_a^i L_i(x)$$

Convolution turns into a sum:

$$\frac{d\sigma}{d\mathcal{O}} = \sum_a \int dx f_a(x) \frac{d\sigma_a}{d\mathcal{O}}(x) = \sum_{a,i} f_a^i G_i$$

Interpolation grid $\{G_i\}_{i=1}^{\infty}$ indep. of PDFs:

$$G_i = \int dx L_i(x) \sigma_a(x)$$

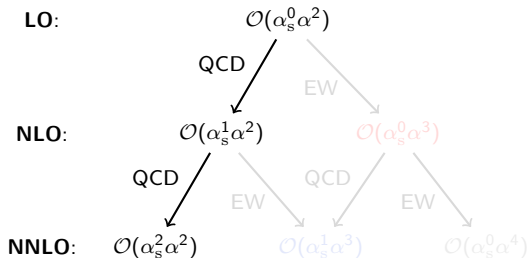
Applications of interpolation grids

- **input for PDF fits** or
- study impact of PDFs uncertainties of observables

Why PineAPPL?

- need support **EW corrections** (any powers of α and/or α_s)
- **performance** becomes very important
- powerful tooling (see later slides)

Which EW corrections will be included for $pp \rightarrow \ell\bar{\ell}$?



- NNLO QCD corrections included in PDF fits

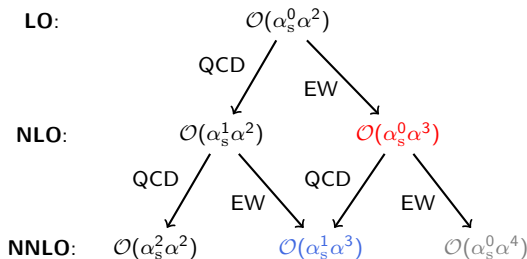
→ include also higher-order α contributions in NNPDF 4.x

- for all processes
- check impact of the corrections,
- be more inclusive, etc.

→ PineAPPL supports all higher-orders

- $\mathcal{O}(\alpha_s \alpha^3)$ might become available at some point ...

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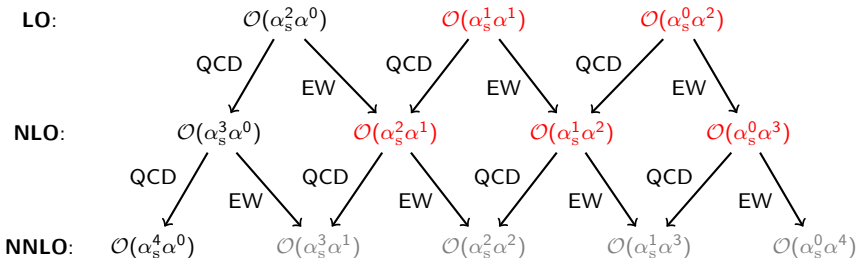
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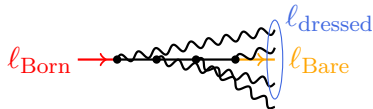
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Which datasets do we fit? Final-state radiation (FSR) subtraction



- **pre-FSR data/Born leptons**: leptons “before they radiate”, calculated using shower inversion (PHOTOS), from
- **post-FSR data/dressed leptons**: leptons with photons recombined around $\Delta R_{f\gamma}$, typically $\Delta R_{f\gamma} = 0.1$
- **bare leptons**: non-collinear safe
- dressing factors

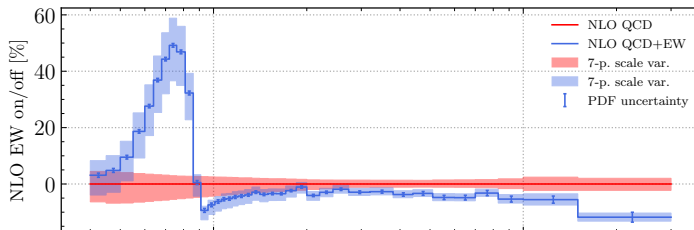
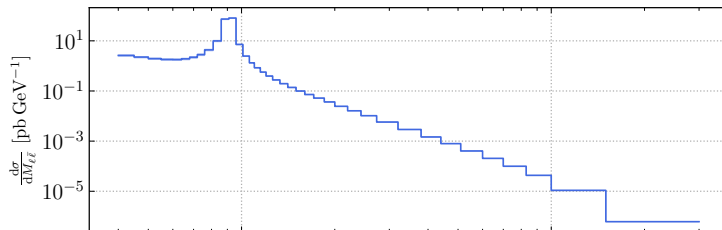
$$C_{\text{dress}} = \frac{d\sigma_{\text{post-FSR}}/d\mathcal{O}}{d\sigma_{\text{pre-FSR}}/d\mathcal{O}}$$

can be very large, **up to 50 %** in invariant mass distributions

- **pre-FSR data** for comparisons with **QCD**-only theory predictions
- **post-FSR data** for comparisons with **EW** corrections (up to one photon at NLO)
 - For some analyses post-FSR data (preferred choice) not published: double counting issue with pre-FSR data!
 - Often C_{dress} (+uncertainty) and pre-FSR dataset given \Rightarrow need to change systematic uncertainties!
- NLO EW PDF dataset largely determined by whether dressed-lepton observables/post-FSR dataset is available

Effects of FSR in Drell–Yan

Differential Drell–Yan cross section at 14 TeV



- distortions of the NLO EW around the Z peak: QED FSR (one-photon emission)
- not described by FSR: weak corrections for large $M_{\ell\bar{\ell}}$ → need full EW corrections
- Experiments: **please** publish dressed-lepton observables/post-FSR datasets!

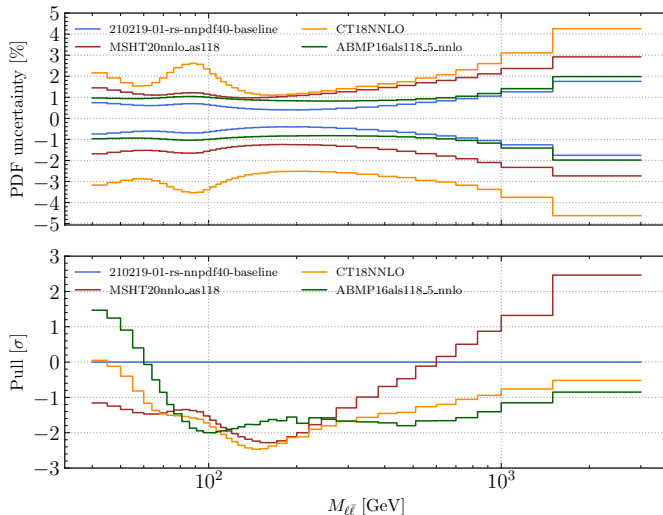
How can I use Madgraph5_aMC@NLO+PineAPPL?

- 1 Install Madgraph5_aMC@NLO: <https://launchpad.net/mg5amcnlo>
- 2 Install PineAPPL: <https://github.com/N3PDF/pineappl>
- 3 Example for generating the DY plot (next page) available at <https://github.com/N3PDF/pineappl/tree/master/examples/mg5amcnlo>

- only one line needs to be added in the mg5amc runcard (rhs)
- No two-phase generation of the grids needed
- Replaces the aMCfast [V. Bertone et al.] interface in Madgraph5_aMC@NLO v2.x
- PineAPPL's CLI allows to easily produce convolutions and plots

```
launch processname
[...]  
set pt1 = 25.0  
set etal = 2.5  
set mll_sf = 116  
set req_acc_F0 0.001  
set pineappl True  
done  
quit
```

Example: DY plot



```
$ pineappl --silent-lhapdf plot \
  DY_14.pineappl.lz4 \
  210219-01-rs-nnpdf40-baseline \
  MSHT20nnlo_as118 \
  CT18NNLO \
  ABMP16als118.5_nnlo > plot.py
$ python3 plot.py
```

- PDF uncertainties for
 - NNPfD 4.0 candidate fit
 - MSHT20
 - CT18 (only main set)
 - ABMP16
- pull: weighted difference w.r.t. NNPfD 4.0 in units of σ

Summary

- NLO EW corrections for PDF fits: upgrading the precision of PDF processes
- PineAPPL: interpolation tool/library for storing PDF-independent theoretical predictions
<https://github.com/N3PDF/pineappl>
- Built-in support in Madgraph5_aMC@NLO v3.1.0
- data issues: FSR subtraction, photon-initiated subtraction, ...

Outlook:

- We will publish grids for ATLAS/CMS/LHCb analyses (PDF processes) soon
- PineAPPL already public and Open Source: <https://n3pdf.github.io/pineappl>
- NLO EW PDF fit not too far away

Interpolation grids

For PDF fitting we need **PDF independent** predictions. Use Lagrange interpolation,

$$f_a(x_1, Q^2)f_b(x_2, Q^2) \approx \sum_{i,j,k} f_a(x_i, Q_k^2)f_b(x_j, Q_k^2)L_i(x_1)L_j(x_2)L_k(Q^2),$$

with Lagrange polynomials L_i over the 3D grid $\{(x_i, x_j, Q_k^2)\}_{i,j,k}$. Insert into master formula:

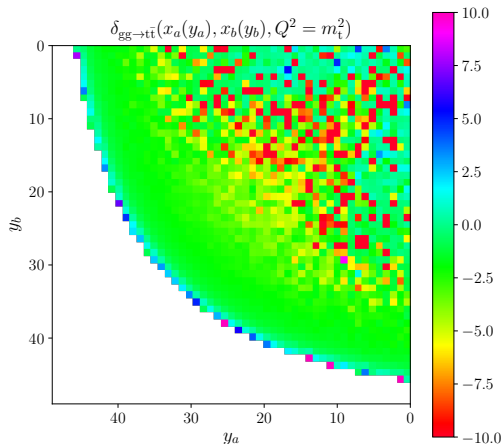
$$\begin{aligned} \frac{d\sigma}{d\mathcal{O}} &= \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 \int_{Q_{\min}^2}^{Q_{\max}^2} dQ^2 f_a(x_1, Q^2)f_b(x_2, Q^2) \frac{d\sigma_{ab}}{d\mathcal{O}}(x_1, x_2, Q^2, \mathcal{O}) \\ &= \sum_{a,b} \sum_{i,j,k} \sum_{m,n} f_a(x_i, Q_k^2)f_b(x_j, Q_k^2)\alpha_s^m(Q^2)\alpha^n \frac{d\Sigma_{abijkmn}}{d\mathcal{O}} \end{aligned}$$

where

$$\frac{d\Sigma_{abijkmn}}{d\mathcal{O}} = \int_0^1 dx_1 \int_0^1 dx_2 \int_{Q_{\min}^2}^{Q_{\max}^2} dQ^2 L_i(x_1)L_j(x_2)L_k(Q^2) \frac{d\sigma_{ab}^{(i,k)}}{d\mathcal{O}}(x_1, x_2, Q^2, \mathcal{O})$$

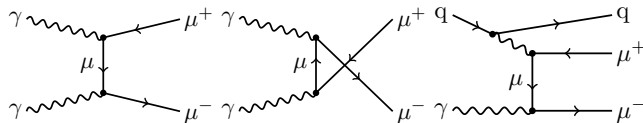
→ generate $\frac{d\Sigma_{abijkmn}}{d\mathcal{O}}$ **once**, perform PDF convolutions very **quickly off-line**

Example: $\Sigma_{ggij021}/\Sigma_{ggij020}$, $\mathcal{O}(\alpha_s^2\alpha)/\mathcal{O}(\alpha_s^2)$ for $gg \rightarrow t\bar{t}$ @ 8 TeV



- no interpolation in y_a , y_b , or Q^2
- correction for ixs roughly -0.5%
- $y_{a/b}(x) = -\ln x_{a/b} + 5(1 - x_{a/b})$, $y(1) = 0$
- lower left corner \rightarrow production threshold
- at threshold: Coulomb singularity
- $y_a \leftrightarrow y_b$ symmetry: initial-state symmetry of $gg \rightarrow t\bar{t}$
- negative correction for larger x_a , x_b

Subtraction of photon–photon contribution

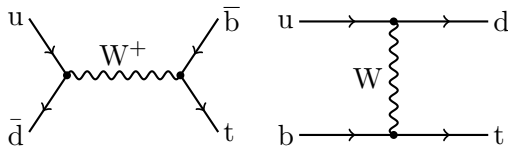


- For ATLAS and CMS it seems to be standard procedure to subtract double-photon induced contributions:
The photon-induced process, $\gamma\gamma \rightarrow \ell\bar{\ell}$, is simulated at LO using Pythia 8 and the MRST2004qed PDF set.
- I am not sure why this is done
- This is a problem: proton contains photons, should be counted towards signal!
- Size of the LO contribution can become significant in large-invariant-mass bins (3 %) depending on the used PDF—up to twice as large for pre-LUXQED photon PDFs

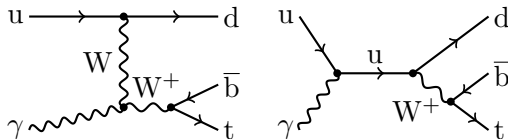
t -channel single-top production

Not properly definable (!?) at NLO EW:

- Analyses, e.g. [ATLAS collaboration], treat s -channels as irreducible background
- single-production at LO:



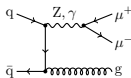
- but at NLO EW not (gauge-invariantly) separable:



→ ignore these datasets

- probably not too important, but see [E.R. Nocera, M. Ubiali, C. Voisey]

Z transverse momentum



$$\mu = M_Z \text{ vs. } \mu = \sqrt{M_Z^2 + (p_T^{\ell\bar{\ell}})^2}$$

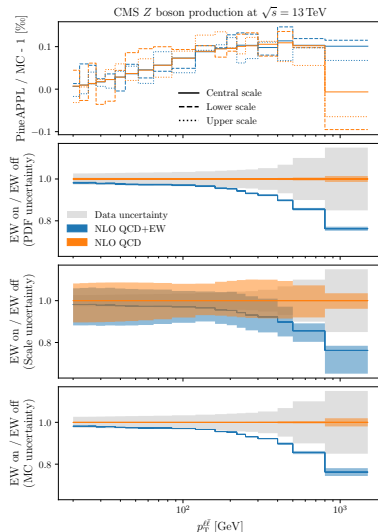
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static scale:

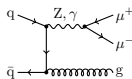
- accidental cancellation of NLO QCD correction \rightarrow uncertainty band shrinks
- NLO EW are artificially enhanced because of normalisation

dynamic scale:

- scale variation is stabilised
- still significant EW corrections, comparable to data uncertainty



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