



INTRINSIC CHARM IN THE PROTON

STEFANO FORTE UNIVERSITÀ DI MILANO & INFN



UNIVERSITÀ DEGLI STUDI DI MILANO

DIPARTIMENTO DI FISICA



MIAPP, AUGUST 19, 2022

HIGH PRECISION AT THE LHC

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INTRINSIC CHARM IN THE PROTON

THE NNPDF COLLABORATION

RICHARD D. BALL, ALESSANDRO CANDIDO, JUAN CRUZ MARTINEZ STEFANO FORTE, TOMMASO GIANI, FELIX HEKHORN, KIRILL KUDASHKIN, GIACOMO MAGNI AND JUAN ROJO

AMSTERDAM-EDINBURGH-INFN-MILAN-NIKHEF



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Evidence for intrinsic charm quarks in the proton

The NNPDF Collaboration

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Abstract

The theory of the strong force, quantum chromodynamics, describes the proton in terms of quarks and gluons. The proton is a state of two up quarks and one down quark bound by gluons, but quantum theory predicts that in addition there is an infinite number of quark– antiquark pairs. Both light and heavy quarks, whose mass is respectively smaller or bigger than the mass of the proton, are revealed inside the proton in high-energy collisions. However, it is unclear whether heavy quarks also exist as a part of the proton wavefunction, which is determined by non-perturbative dynamics and accordingly unknown: so-called intrinsic heavy quarks¹. It has been argued for a long time that the proton could have a sizable intrinsic component of the lightest heavy quark, the charm quark. Innumerable efforts to

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Benjamin Thompson & Nick Petrić Howe Nature Nature Podcast 17 Aug 2022

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Ramona Vogt Nature News & Views 17 Aug 2022



"INTRINSIC" CONSTITUENTS IN THE PROTON AT THE SSC (1984)

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INTRINSIC CHEVROLETS AT THE SSC

DE85 013896

Stanley J. Brodsky

Stanford Linear Accelerator Center, Stanford University, Stanford CA 94305

John C. Collins

Department of Physics, Illinois Institute of Technology, Chicago IL 60616 and

High Energy Physics Division, Argonne National Laboratory, Argonne IL 60439

Stephen D. Ellis

Department of Physics, FM-15, University of Washington, Seattle WA 98195

John F. Gunion

Department of Physics, University of California, Davis CA 95616

Alfred H. Mueller

Department of Physics, Columbia University, New York NY 10027

Summary

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The possibility of the production at high energy of heavy quarks, supersymmetric particles and other large mass colored systems via the intrinsic twist-six components in the proton wave function is discussed. While the existing data do not rule out the possible relevance of intrinsic charm production at present energies, the extrapolation of such intrinsic contributions to very high masses and energies suggests that they will not play an important role at the SSC. sufficiently large. The data from the EMC collaboration⁴ on deep-inelastic muon scattering could also be intepreted as suggesting an unexpectedly large charm structure function in the region z > 0.3.

The possible existence of such a new production mechanism is of great importance for design considerations at the SSC^{5,6}. An example of the importance of this issue is that, if intrinsic large x production is dominant, experiments and, perhaps, even the machine should be designed to focus on the forward "diffractive" regime⁵. The ques-



IS IT TRIVIAL?

• SEA PDFS AT HIGH SCALE ALL LOOK ALIKE

- IF $Q \gg m_c$, CHARM MASS NEGLIGIBLE: $\ln \frac{Q^2 + m_c^2}{m_c^2} \approx \ln \frac{Q^2}{m_c^2}$
- GLUON RADIATION IS FLAVOR BLIND

DECOUPLING

EVOLVE CHARM PDF ($N_f = 4$ SCHEME) DOWN TO $Q \sim m_c$



• IF $Q \sim m_c$ ($m_c = 1.51 \text{ GeV}$), CHARM QUARK DECOUPLES (Collins, Wilczek, Zee, 1978): $\ln \frac{Q^2 + m_c^2}{m_c^2} \approx \frac{m_c^2}{Q^2}$

- $N_f = 3$ active flavors in β function & evolution equations
- DECOUPLING VS $\overline{\mathrm{MS}} \Leftrightarrow$ DIFFERENT RENORMALIZATION & FACTORIZATION SCHEMES

MATCHING

• PDFS, α_s in $N_f = 3$ & $N_f = 4$ Related by matching conditions OME CONTRIBUTING TO THE CHARM PDF

SOLID \Rightarrow HEAVY; DASHED \Rightarrow LIGHT

M. Buza et al.: Charm

- DETERMINED BY COMPUTING
 - **OPERATOR MATRIX ELEMENTS**

IN EITHER SCHEME AND EQUATING:

NNLO (Buza, et al., 1998),

 $N^{3}LO$ (Ablinger, Blümlein et al, 2009-2017)



Fig. 2. $O(\alpha_s^2)$ contributions to the purely-singlet OME $A_{q'q}^{PS}$. Here q and q' are represented by the *dashed* and *solid lines* vertices q' = H these graphs contribute to the heavy-quark OME A_{Hq}^{PS}

PERTURBATIVE CHARM

- NO CHARM PDF IN $N_f = 3$ SCHEME
- IN $N_f = 4$ Scheme, charm determined by perturbative matching starting at NNLO (two loops) does not vanish at any scale (heavy quark loops)



NLO & NNLO CHARM PDF; $n_f = 3$ vs. $n_f = 4$

INTRINSIC CHARM

- $n_f = 3$ Scheme \Rightarrow charm PDF scale-independent (no collinear logs)
- PERTURBATIVE CHARM \Leftrightarrow VANISHING $n_f = 3$ CHARM
- CHARM PDF AT ALL SCALES FULLY DETERMINED BY MATCHING CONDITIONS

INTRINSIC CHARM \Leftrightarrow NONVANISHING STATIC $n_f = 3$ -SCHEME CHARM PDF

THE NNPDF4.0 CHARM PDF ($n_f = 4$ scheme)

- NNPDF4.0 CHARM PDF \Rightarrow DETERMINED FROM THE DATA ALONG WITH ALL OTHER PDFS:
 - MORE REALISTIC UNCERTAINTIES
 - STABLE UPON VARIATION OF m_c
 - INDEPENDENT OF MATCHING CONDITIONS
- **DIFFERS SIGNIFICANTLY** FROM PERTURBATIVE CHARM
- MATCHING CONDITIONS PERTURBATIVELY UNSTABLE!



INTRINSIC CHARM?

- INVERT MATCHING CONDITIONS $\Rightarrow N_f = 3$ Charm PDF
- COMPARE NNLO & N³LO INVERSION TO CONTROL MHOU

INTRINSIC CHARM! THE EKO CODE (Candido, Hekhorn, Magni, 2022)

- IMPLEMENTS DIRECT & INVERSE EVOLUTION & MATCHING
- N³LO MATCHING ALSO IMPLEMENTED



THE INTRINSIC ($N_f = 3$) CHARM PDF (NNLO)

- MHOU ESTIMATED FROM N³LO-NNLO DIFFERENCE
 - LARGE UNCERTAINTY AT SMALL x
 - NEGLIGIBLE UNCERTAINTY IN VALENCE REGION
- COMPATIBLE WITH ZERO AT SMALL x
- CLEAR EVIDENCE FOR "INTRINSIC" VALENCE PEAK

MODELS

- SHAPE OF INTRINSIC CHARM PREDICTED BY MODELS
- FOCK-SPACE WAVE FUNCTION (Brosky, Hoyer, Peterson, Sakai, 1980)
- MESON CLOUD (Hobbs, Londergan, Melnitchouk, 2014)



NNPDF4.0 INTRINSIC CHARM VS. MODELS

SURPRIZING AGREEMENT!

MORE DATA EMC 1983

- DIRECT MEASUREMENT OF THE CHARM STRUCTURE FUNCTION F_2^c
- EVIDENCE FOR INTRINSIC CHARM CLAIMED, BUT EXPERIMENT DISPUTED
- NOT INCLUDED IN DEFAULT NNPDF4.0



INTRINSIC CHARM WITH EMC DATA INCLUDED

COMPLETE CONSISTENCY!



MORE DATA LHCB 2021

Z

c

Z

c

INTRINSIC CHARM WITH LHCB DATA INCLUDED: COMPLETE CONSISTENCY





More than 3 σ evidence

SUMMARY

WE FITTED THE CHARM PDF IN ORDER TO GET

- REALISTIC ERROR ESTIMATE
- NO STRONG DEPENDENCE ON CHARM MASS
- \bullet no sensitivity to MHOU in matching condition

WE FOUND

- LARGE UNCERTAINTIES AND CHARM COMPATIBLE WITH ZERO AT SMALL \boldsymbol{x}
- THREE- σ EVIDENCE FOR AN INTRINSIC CHARM VALENCE PEAK



WHICH DATA DRIVE THE ANSWER?:

DATA SUBSETS THREE FLAVOR SCHEME DIS ONLY COLLIDER ONLY



NO LHCB







STABILITY:

CHARM MASS



STABILITY:

DATA SUBSETS FOUR FLAVOR SCHEME

