



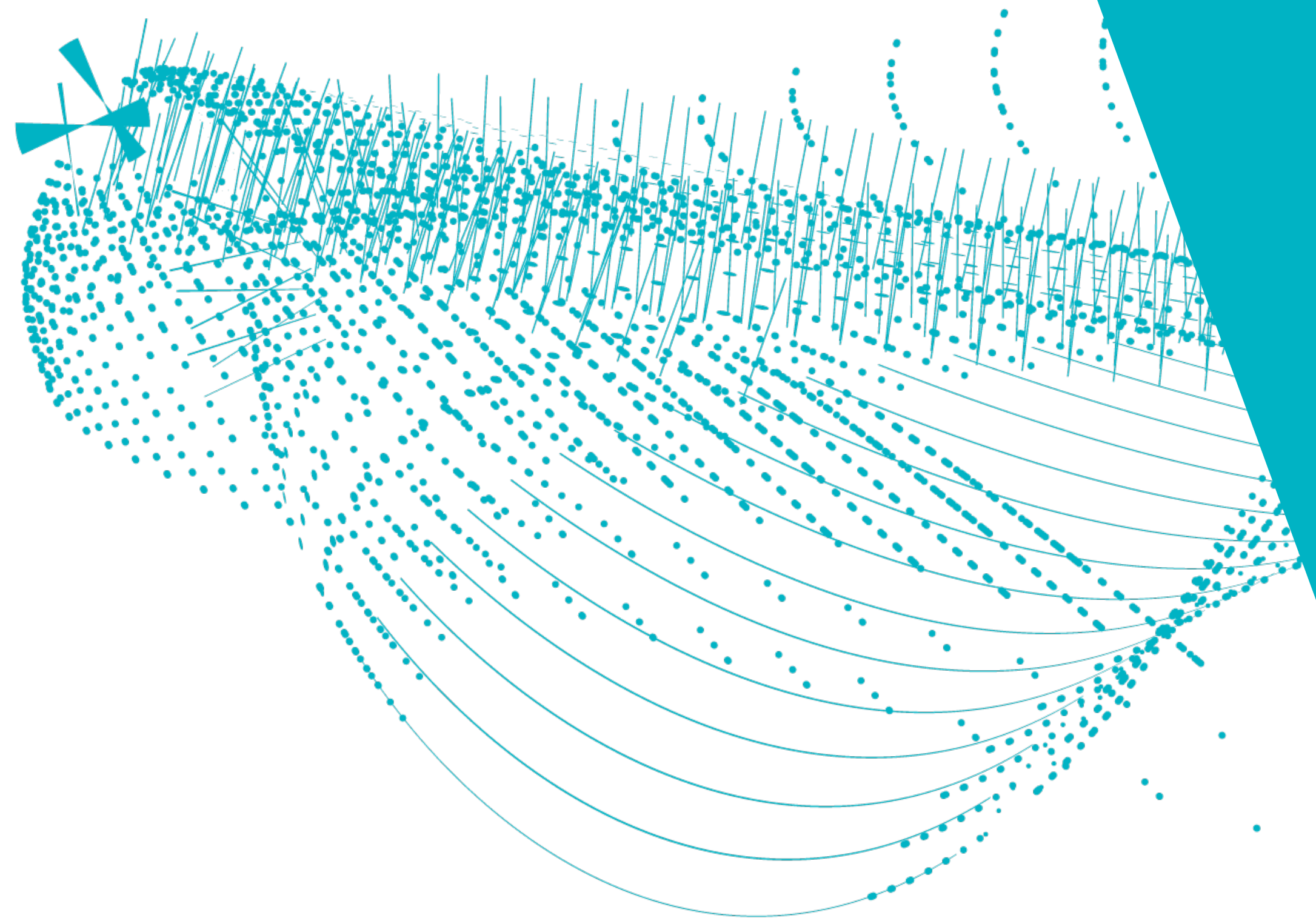
# NNPDF updates and the path towards NNPDF4.1: Data, theory, and Methodology



Tanjona R. Rabemananjara on behalf of the NNPDF collaboration  
DIS 2025, March 25th 2025  
Cape Town, South Africa

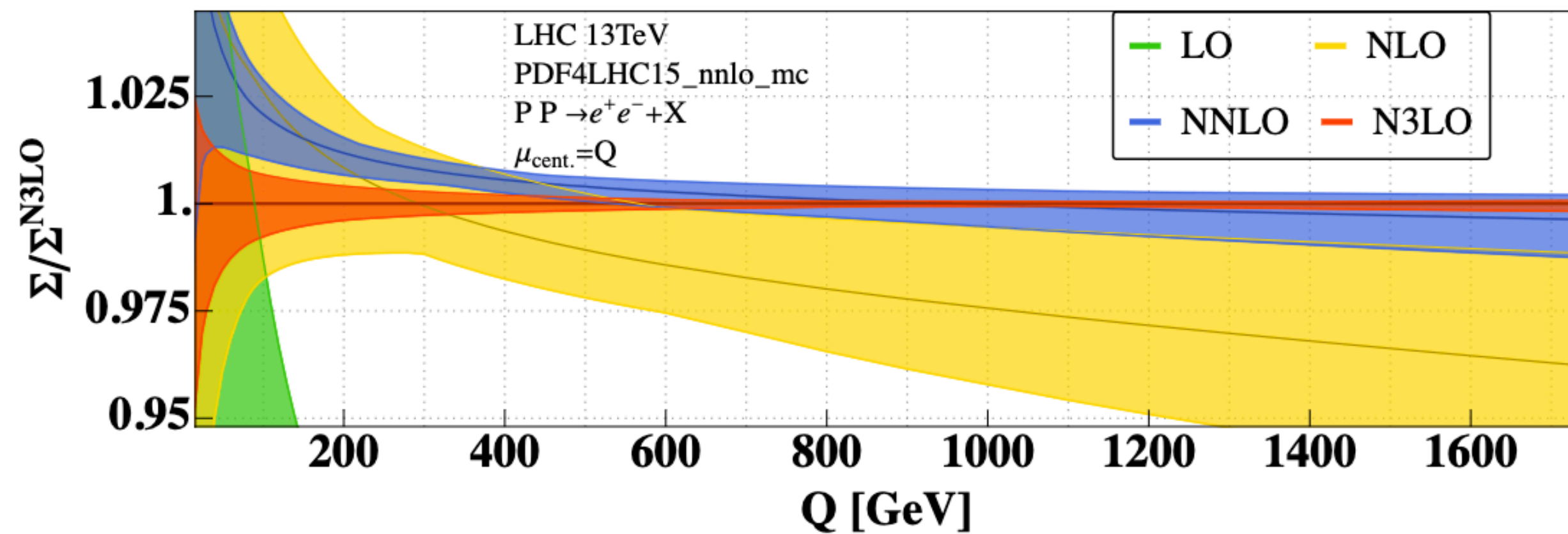


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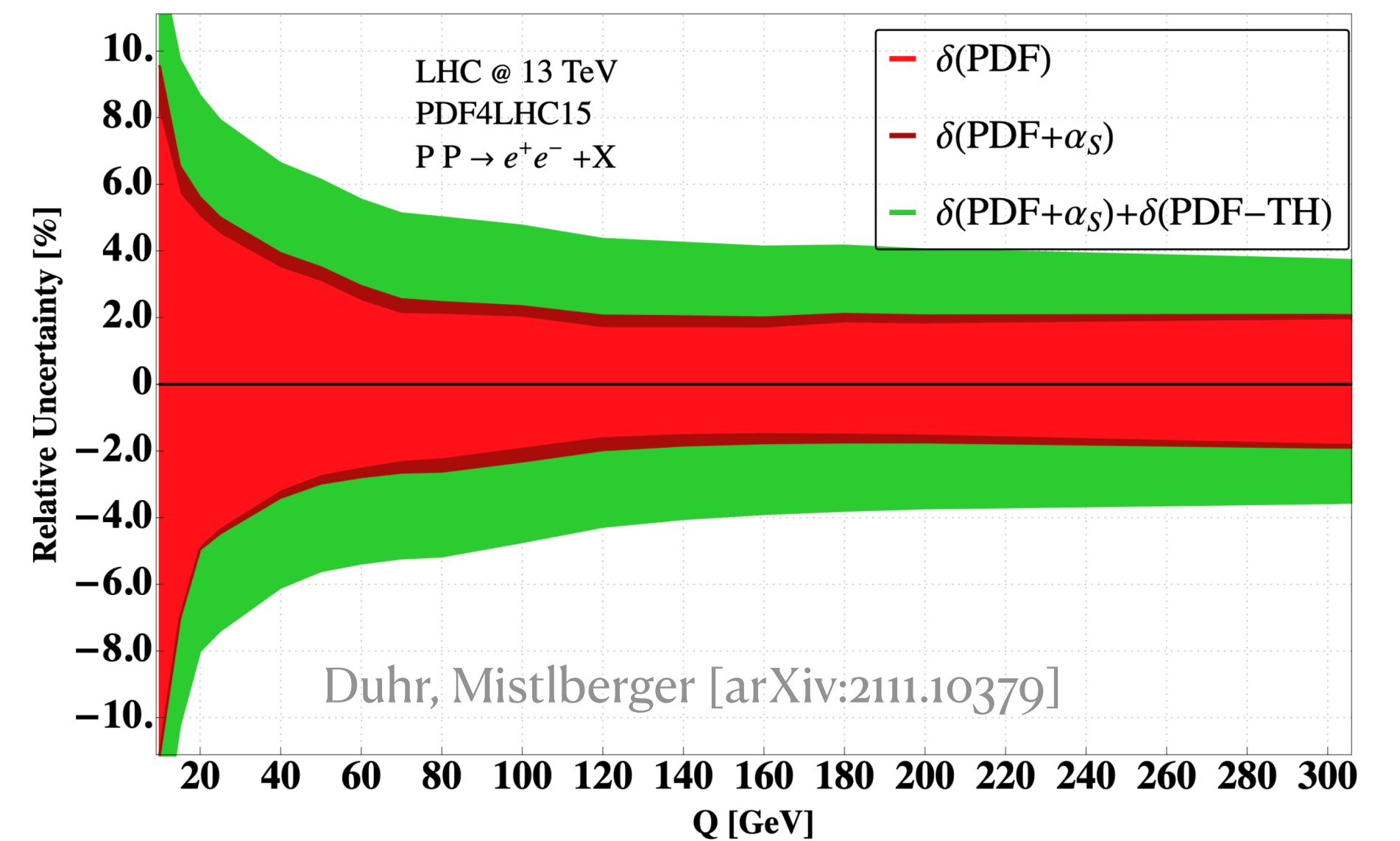
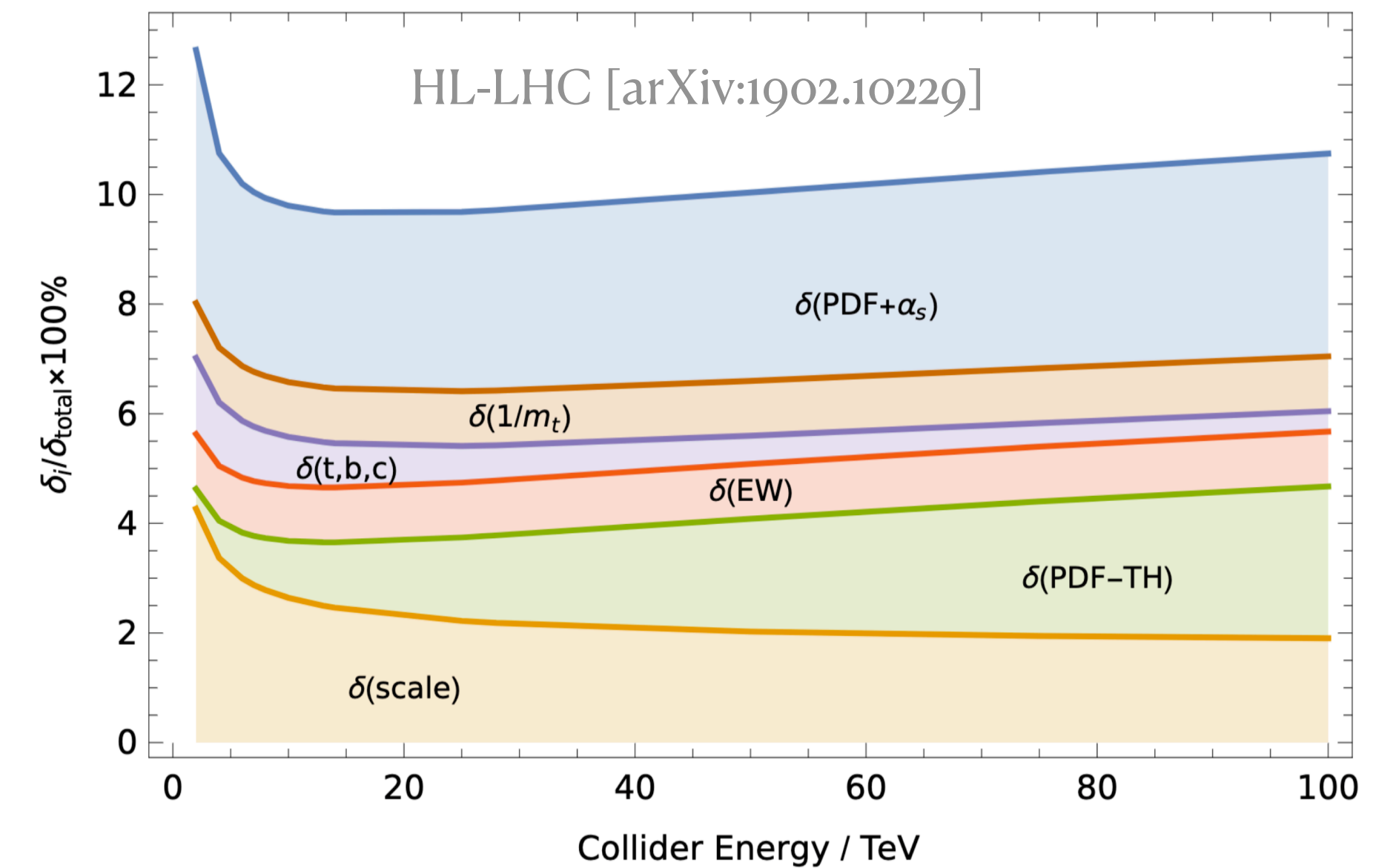


# Motivation: Why do we still care about PDFs?

- ❖ **PDFs are becoming a bottleneck for LHC precision calculations** with the largest uncertainties along with the incomplete knowledge of  $\alpha_s \iff$  **New Physics**
- ❖ **Progress on N3LO calculations** for Higgs (ggF, VBF, VH) & NC/CC DY processes requires **N3LO PDFs**
- ❖ **QED effects in PDFs** are **no longer negligible** as experimental measurements become more precise and determination of parton densities more accurate
- ❖ **Theoretical uncertainties** on PDFs are crucial to assess accuracy/uncertainties on MHOUs & IHOUs

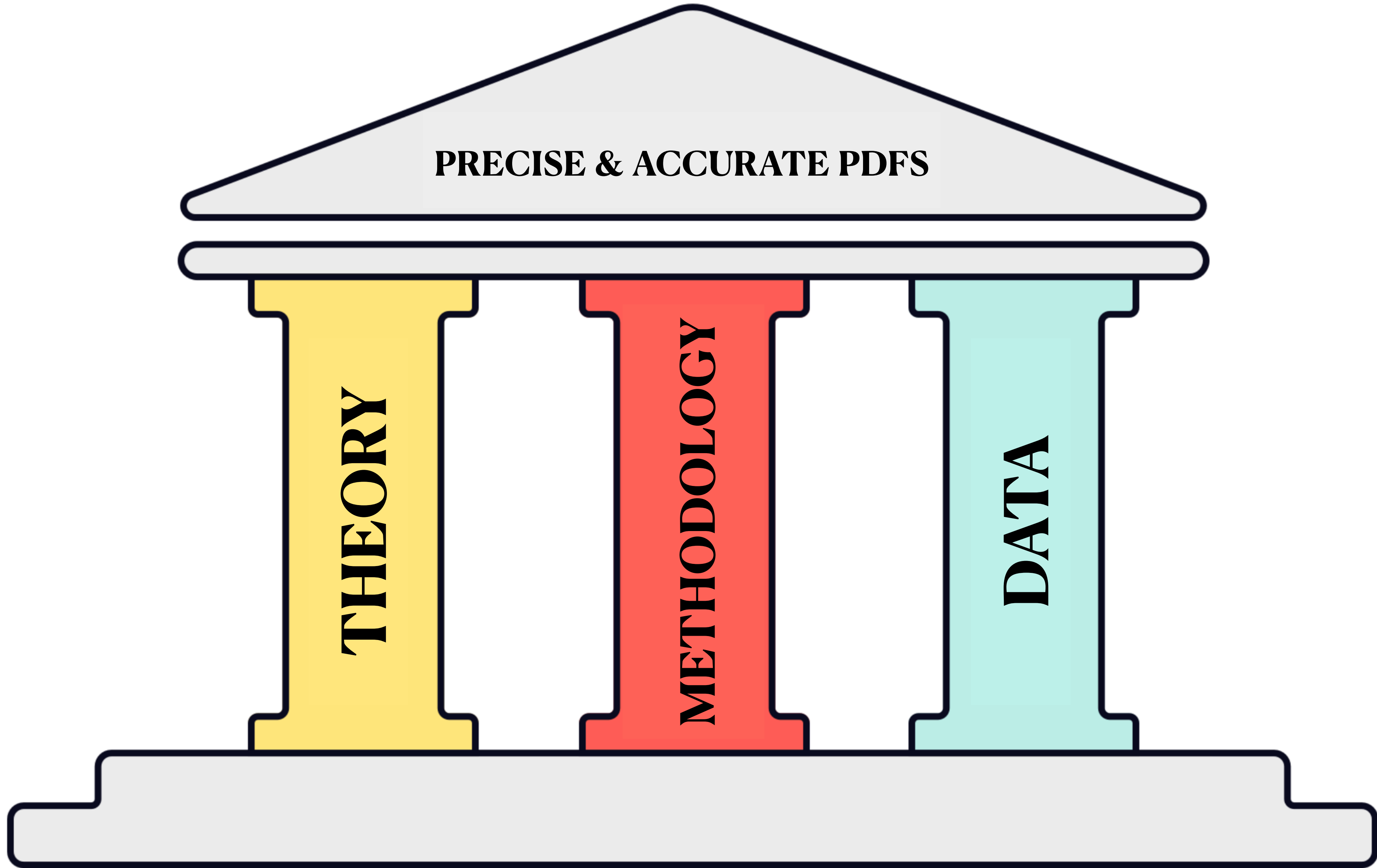


Duhr, Mistlberger [arXiv:2111.10379]



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# Outline of the Talk



# NNPDF Timeline

|         |   |                                       |
|---------|---|---------------------------------------|
| 2Q 2026 | NNPDF4.1  | [WIP]                                 |
| 3Q 2025 | PDFs with Higher-Twist corrections                            | [WIP]                                 |
| 2Q 2025 | Closure test with inconsistencies                             | [WIP]                                 |
| 2Q 2025 | alpha_s determination with aN3LO $\otimes$ QED $\otimes$ MHOU | [WIP] (see RS's talk this afternoon)  |
| 03/2025 | NNPDFpol2.0 - helicity dependent PDFs                         | [arXiv:2503.11814] (presented in WG5) |
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| 09/2022 | PDFs and New Physics (Afb asymmetry)                          | [arXiv:2209.08115]                    |
| 08/2022 | Intrinsic charm   | [arXiv:2208.08372]                    |
| 09/2021 | NNPDF4.0 (code paper)   | [arXiv:2109.02671]                    |
| 09/2021 | NNPDF4.0 (main release)                                       | [arXiv:2109.02653]                    |

# Part I: Theory

|         |  |                                       |
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# Missing Higher Order Uncertainties (MHOUs) @ (a)N{2,3}LO

For a given observable  $\mathcal{O}$ , **MHOUs** are commonly estimated by **varying the unphysical scales** in the **parton evolutions** and in the **partonic cross-sections**:

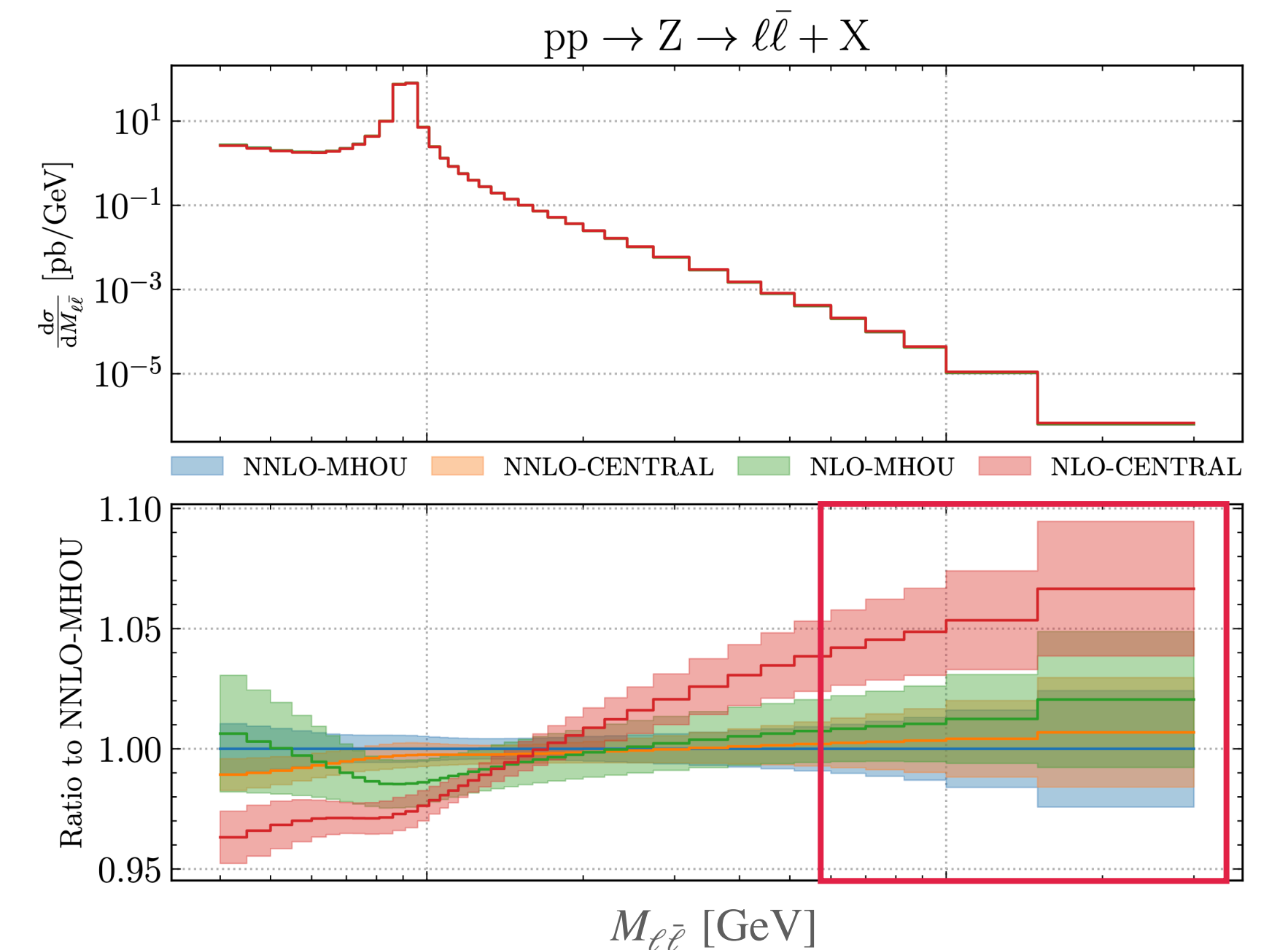
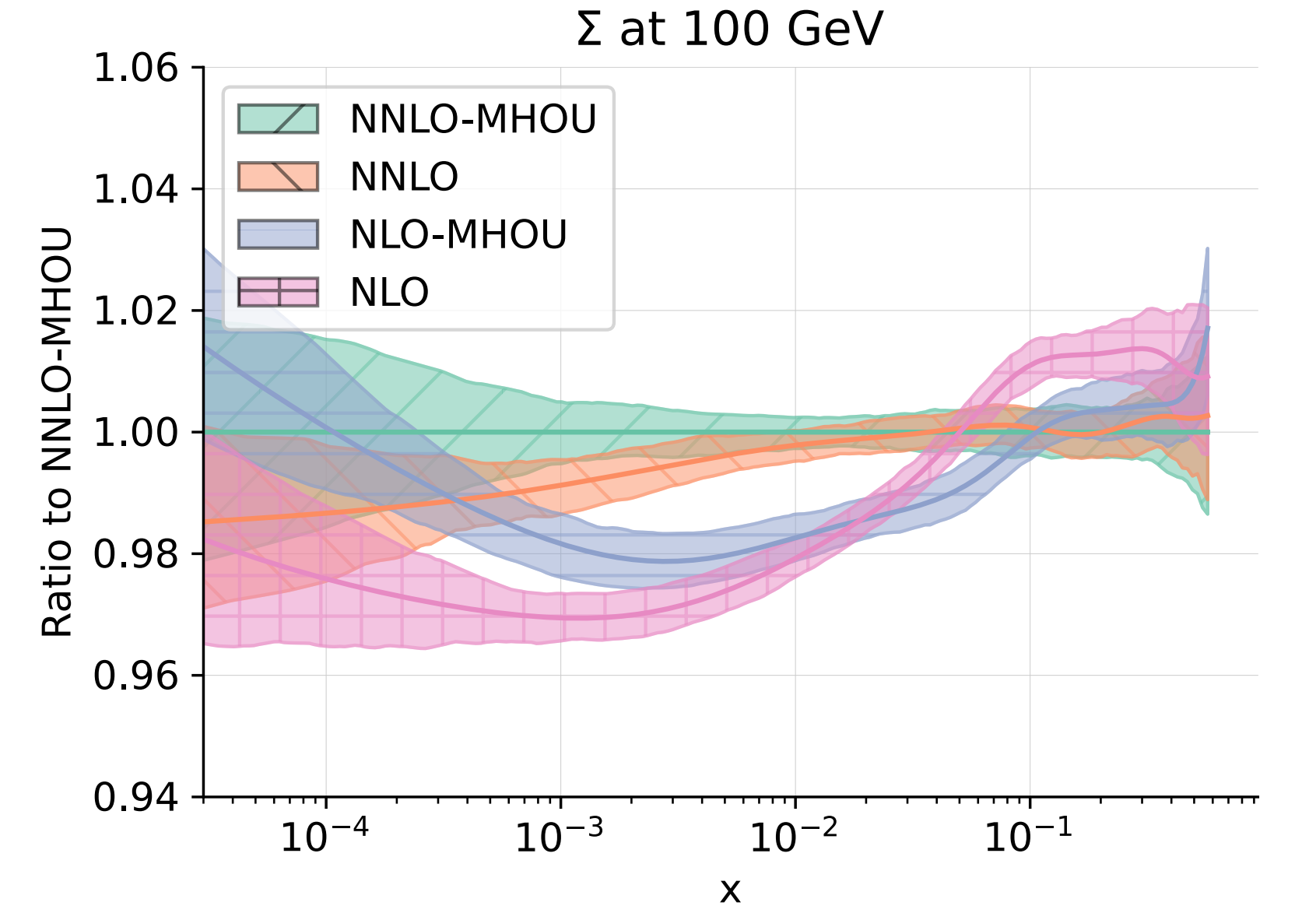
$$\mathcal{O}\left(\alpha_s(\mu^2), \frac{Q^2}{\mu_F^2}, \frac{Q^2}{\mu_R^2}\right) = \mathcal{L}\left(\alpha_s(\mu_F^2), \frac{Q^2}{\mu_F^2}\right) \mathcal{O}\left(\alpha_s(\mu_R^2), \frac{Q^2}{\mu_R^2}\right)$$

Variation of **Factorisation Scale**  $\kappa_F = Q^2/\mu_R^2$  estimates MHOUs from Anomalous Dimensions in the evolution while variation of **Renormalisation Scale**  $\kappa_R = Q^2/\mu_R^2$  estimates MHOUs from partonic cross-sections.

**MHOUs** can be added as a nuisance parameter to the Covariance Matrix [arXiv:1906.10698; arXiv:2105.05114]

$$\text{cov}_{i,j} = \text{cov}_{i,j}^{\text{exp}} + \text{cov}_{i,j}^{\text{MHOu}}, \quad \text{cov}_{i,j}^{\text{MHOu}} = \frac{1}{N_{\text{var}} - 1} \sum_{k=1}^{N_{\text{var}}} (S_{i,k} - \bar{S}_i) (S_{j,k} - \bar{S}_j)$$

**7-point scale variation prescription** is used. Points belonging to the same process are **CORRELATED** by  $\kappa_R$ -variation while  $\kappa_F$  correlates all the points.



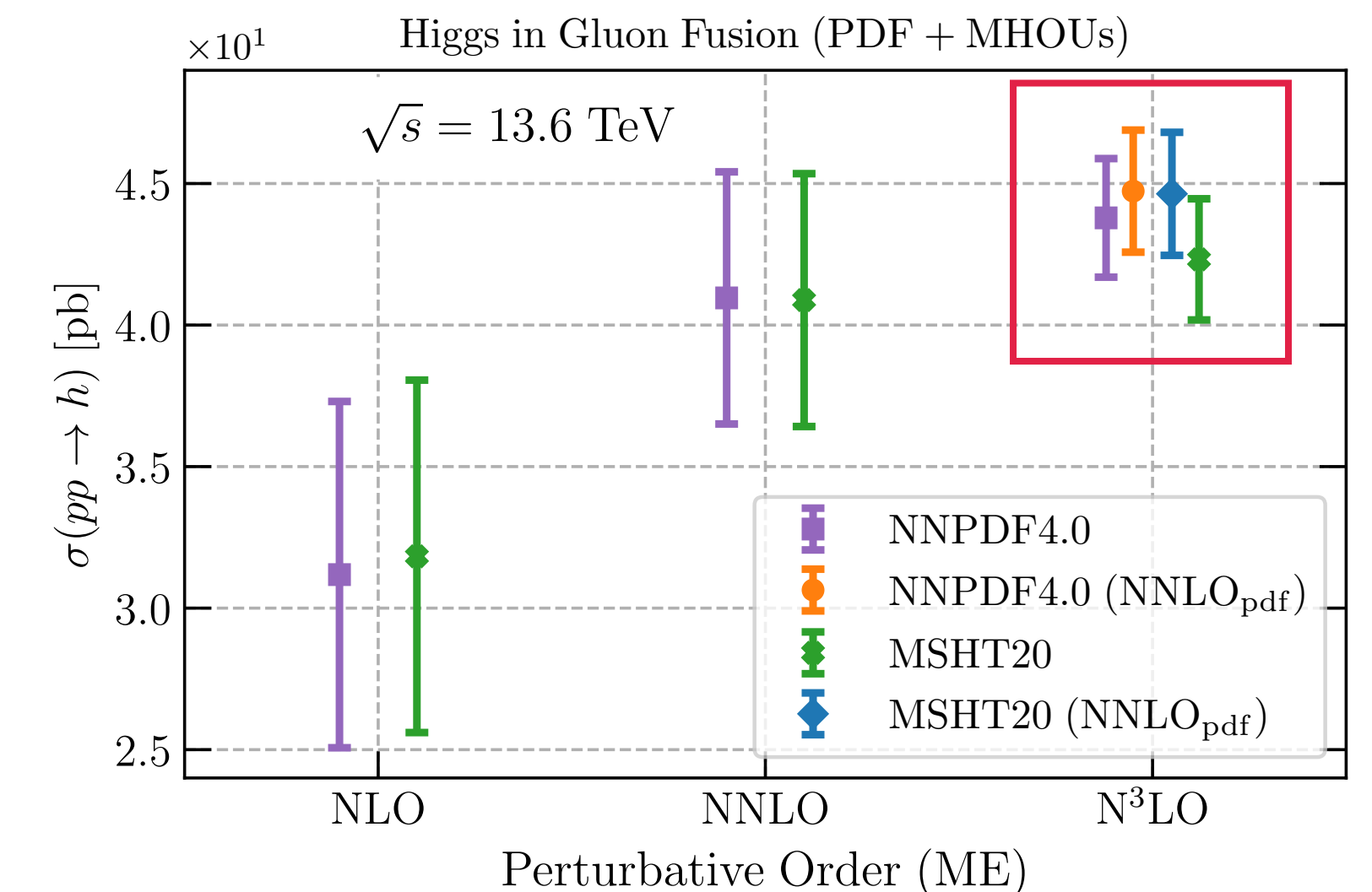
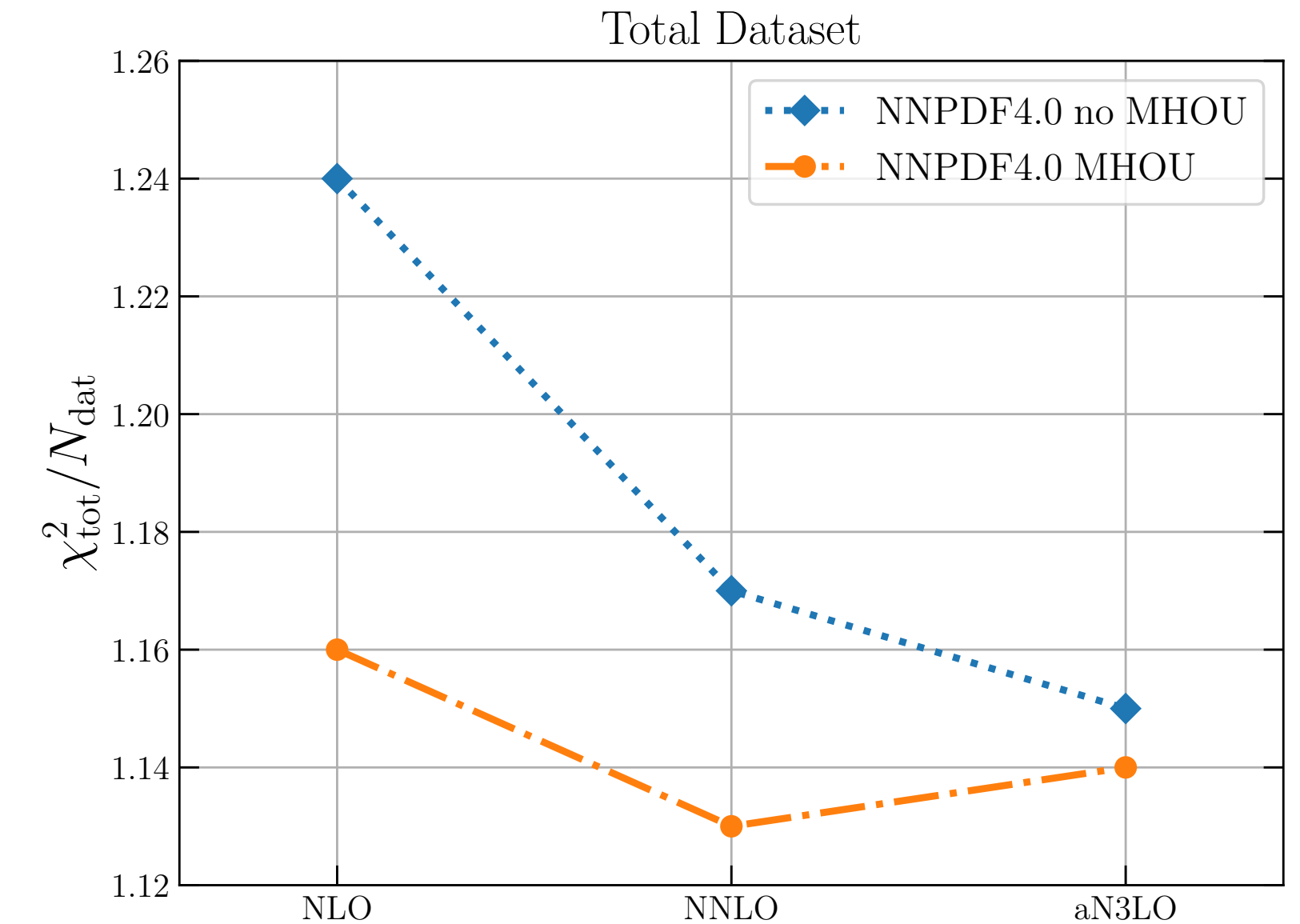
# Approximate N3LO (aN3LO) determination

NNPDF4.0 determination of aN3LO PDFs:

- ❖ **DGLAP Evolution:** accurate numerical approximations splitting functions (10 lowest moments, large- $x$  and small- $x$  limits)
- ❖ **Matching conditions:** all relevant terms are known (all exact,  $a_{Hg}^{(3)}$  parametrised)
- ❖ **DIS Coefficients Functions:** massless coefficients (both NC and CC) are known. Massive NC can be approximated
- ❖ **Hadronic coefficients:** some DY coefficients are known, but not yet available in a format suitable for PDFs fits, corrections to Jets and processes are still unknown

New N3LO terms that will be part of NNPDF4.1:

- ❖ Higher splitting function moments:  $P_{gg}^{(3)}, P_{gq}^{(3)}$  [arXiv:2404.09701; arXiv:2410.08089]
- ❖ Improved parametrisation for  $a_{H,g}^{(3)}$  **matching conditions** [arXiv:2403.00513]



**~1.5% suppression  
w.r.t. NNLO PDFs**

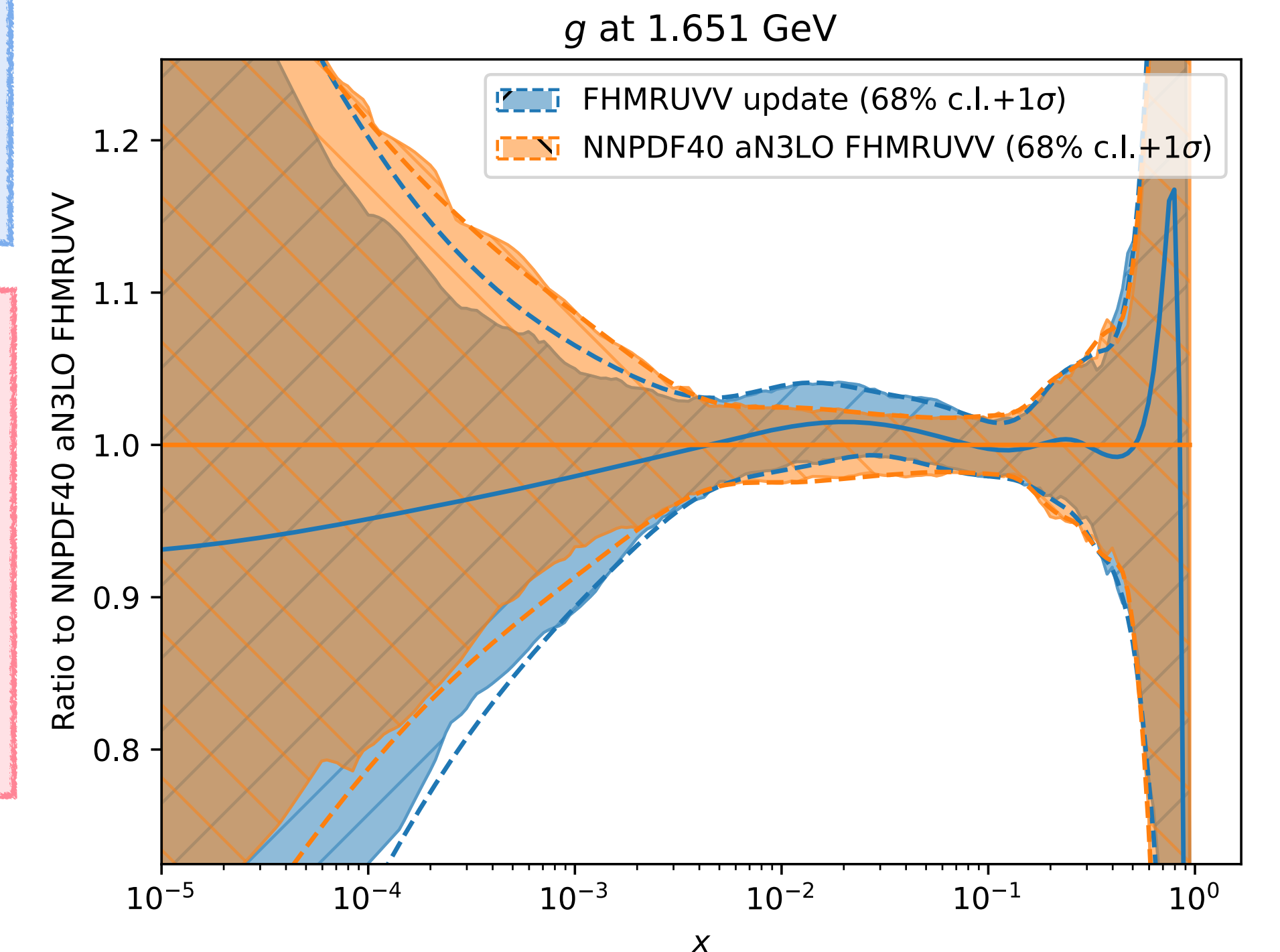
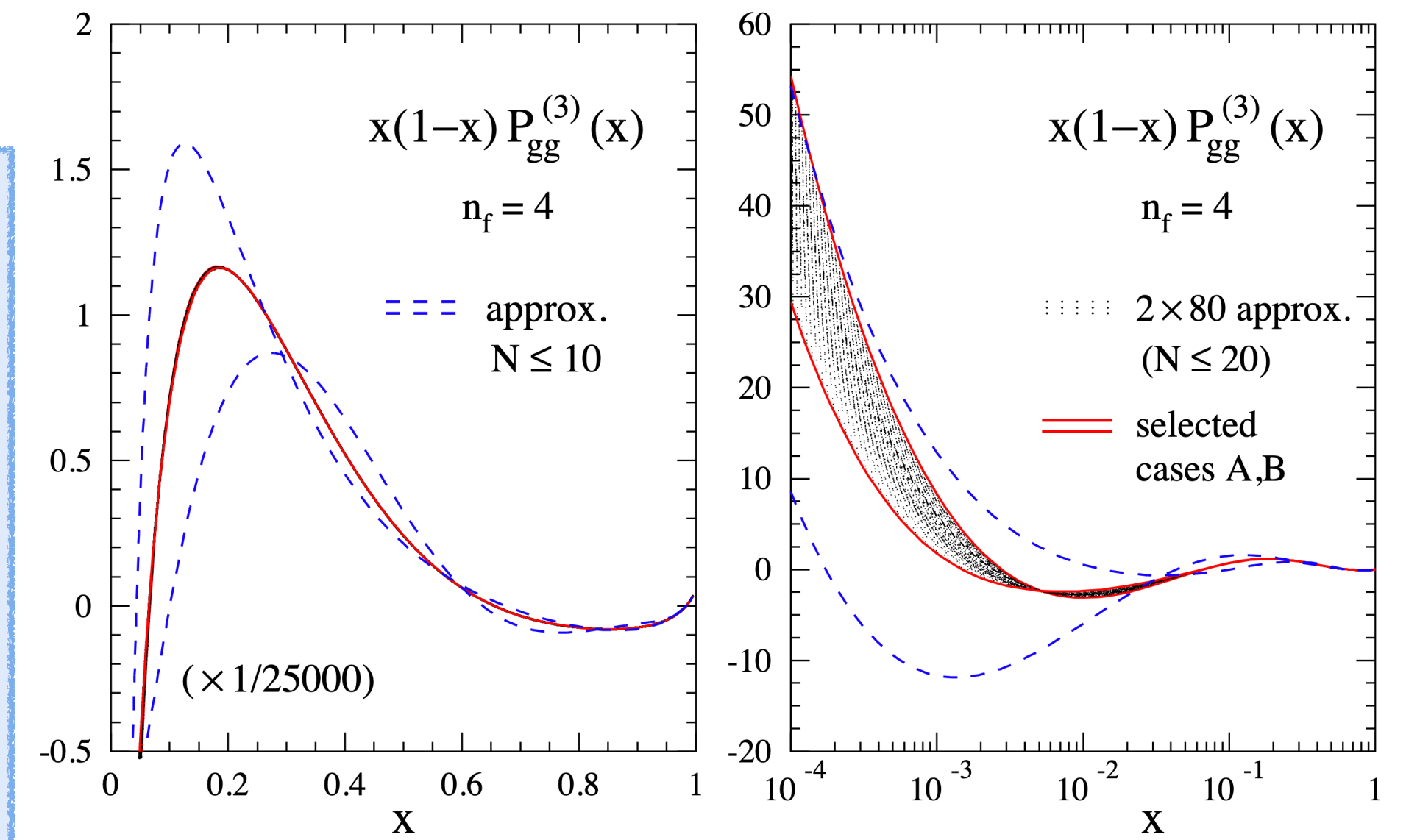
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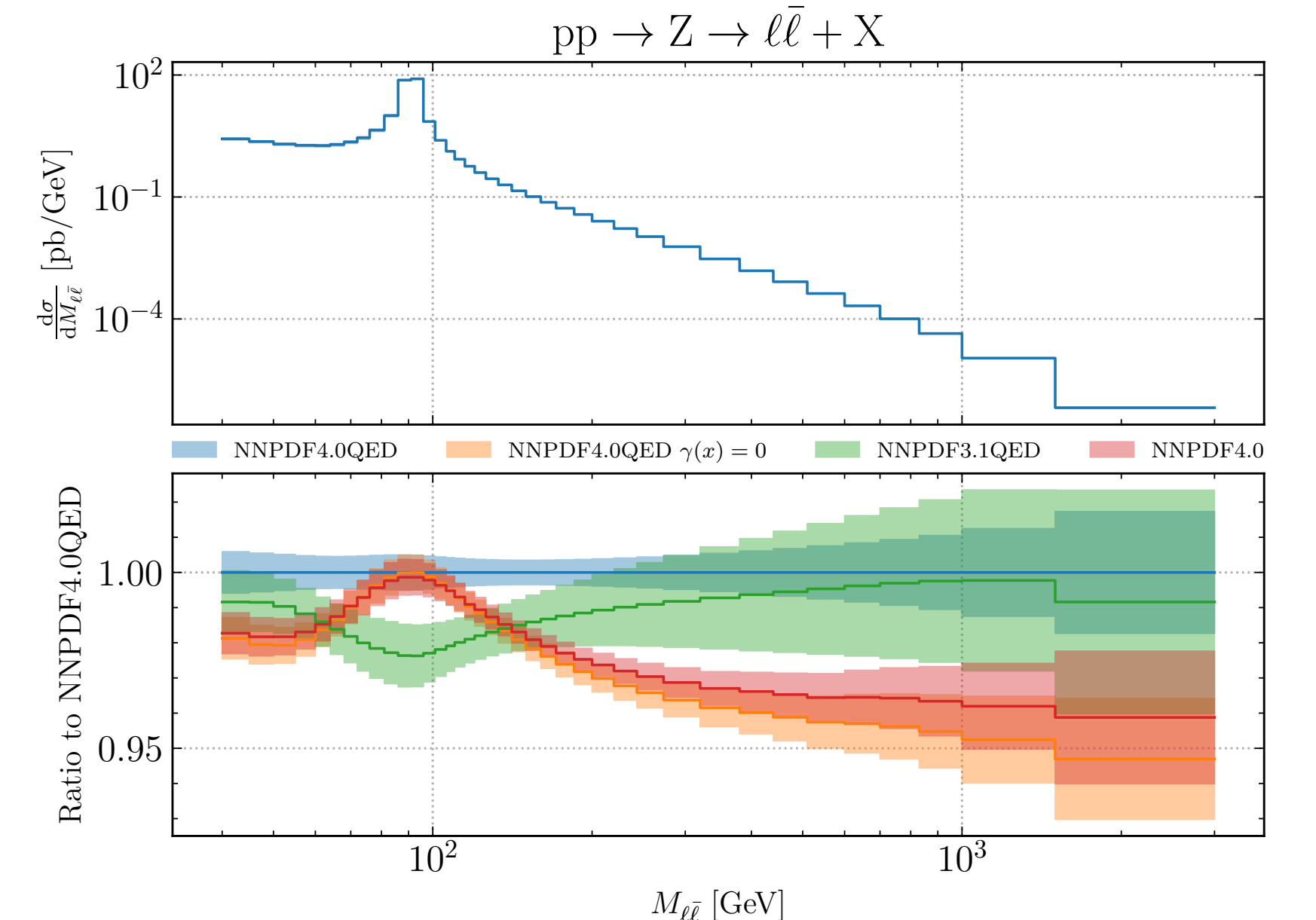
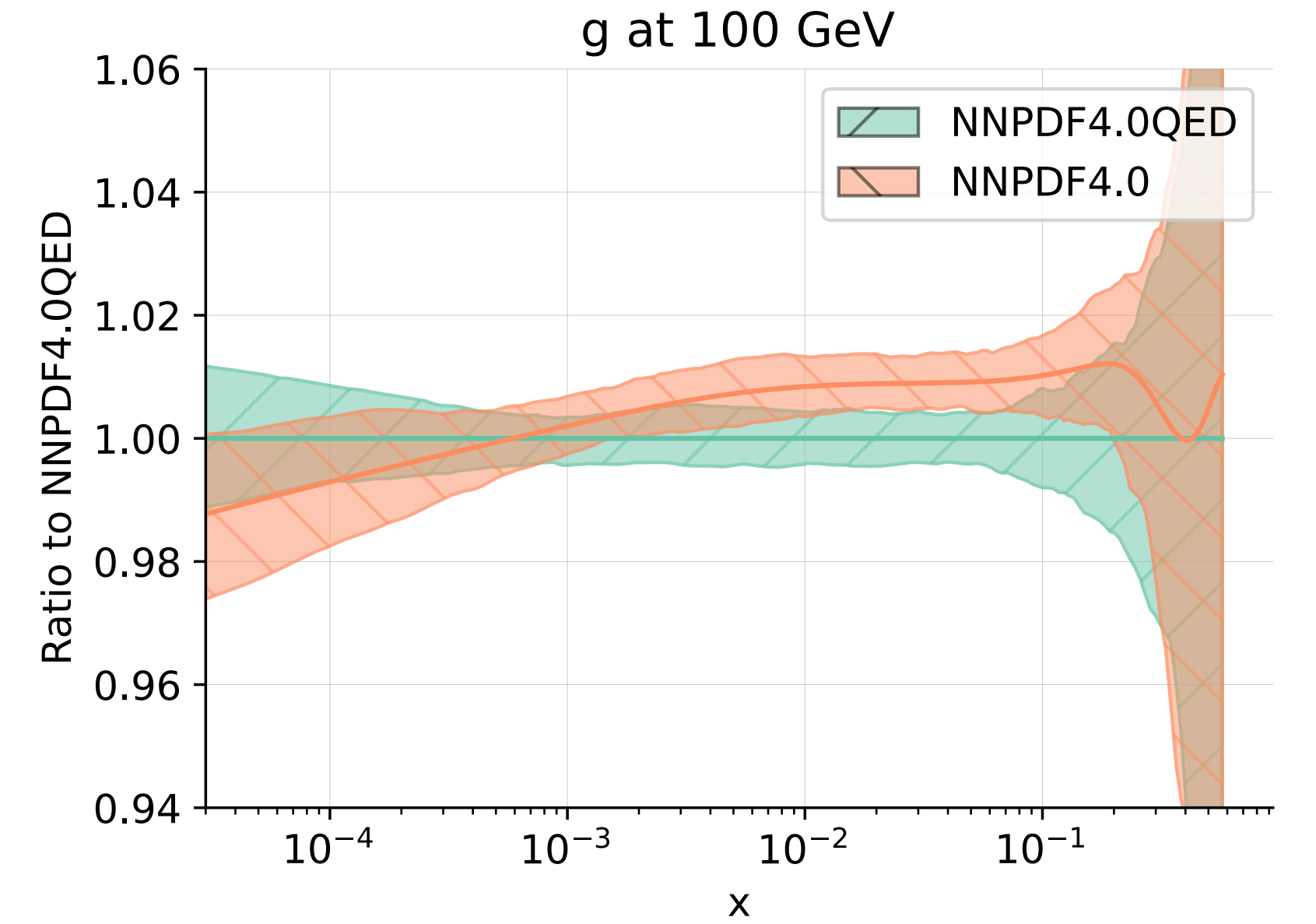
# QED corrections & Photon PDF

$\gamma$ -PDFs are computed from **DIS structure functions** [arXiv:1607.04266]:

$$x\gamma(x, \mu^2) = \frac{2}{a_{em}(\mu^2)} \int_x^1 \frac{dz}{z} \left\{ \int_{\frac{m_p^2 x^2}{(1-z)}}^{\frac{\mu^2}{(1-z)}} \frac{dQ^2}{Q^2} a_{em}^2(Q^2) \left[ -z^2 F_L(x/z, Q^2) \right. \right. \\ \left. \left. + \left( zP_{\gamma q}(z) + \frac{2x^2 m_p^2}{Q^2} \right) F_2(x/z, Q^2) \right] - a_{em}^2(\mu^2) z^2 F_2(x/z, \mu^2) \right\}$$

- ❖ Depending on the kinematic region the structure functions are computed from: **Elastic DIS**, **Resonance**, **Shallow Inelastic**, **DIS**
- ❖  $\gamma(x, Q^2)$  is computed iteratively during the fit
- ❖ Mixed QED $\otimes$ QCD DGLAP evolution: more difficult to diagonalise due to how  $\gamma$  couples differently to up-like and down-like quarks  $\implies$  **Unified Evolution Basis**
- ❖ While  $\gamma(x, Q^2)$  depends on the PDFs through the structure functions, it affects their determination during the iterative procedure.
- ❖ Additional **mixed QED $\otimes$ QCD Momentum Sum rules**:

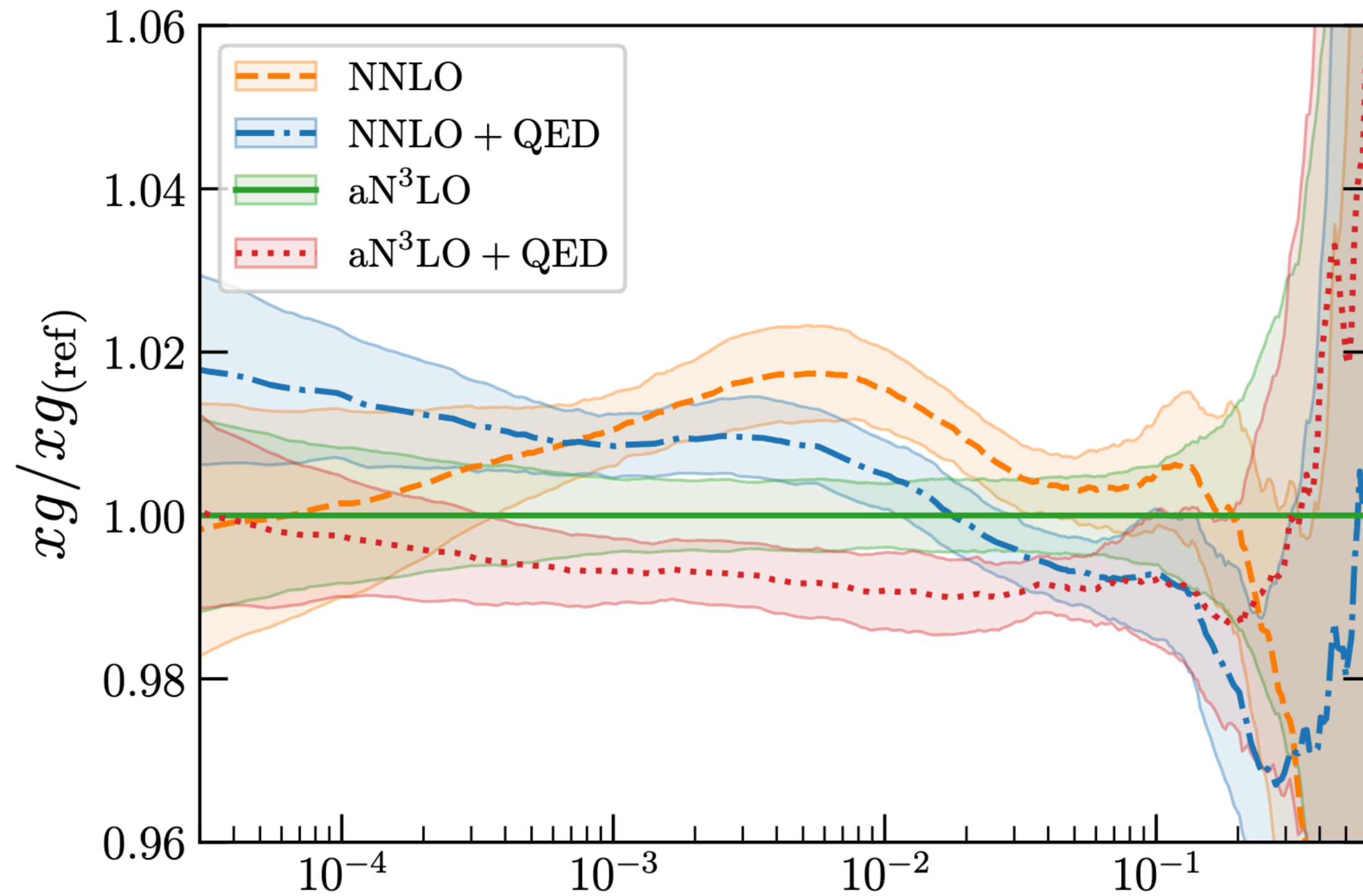
$$\int_0^1 dx (x\Sigma + xg + x\gamma)(x, Q^2) = 1$$



# State-of-the-Art: “aN3LO $\otimes$ QED $\otimes$ MHOU”

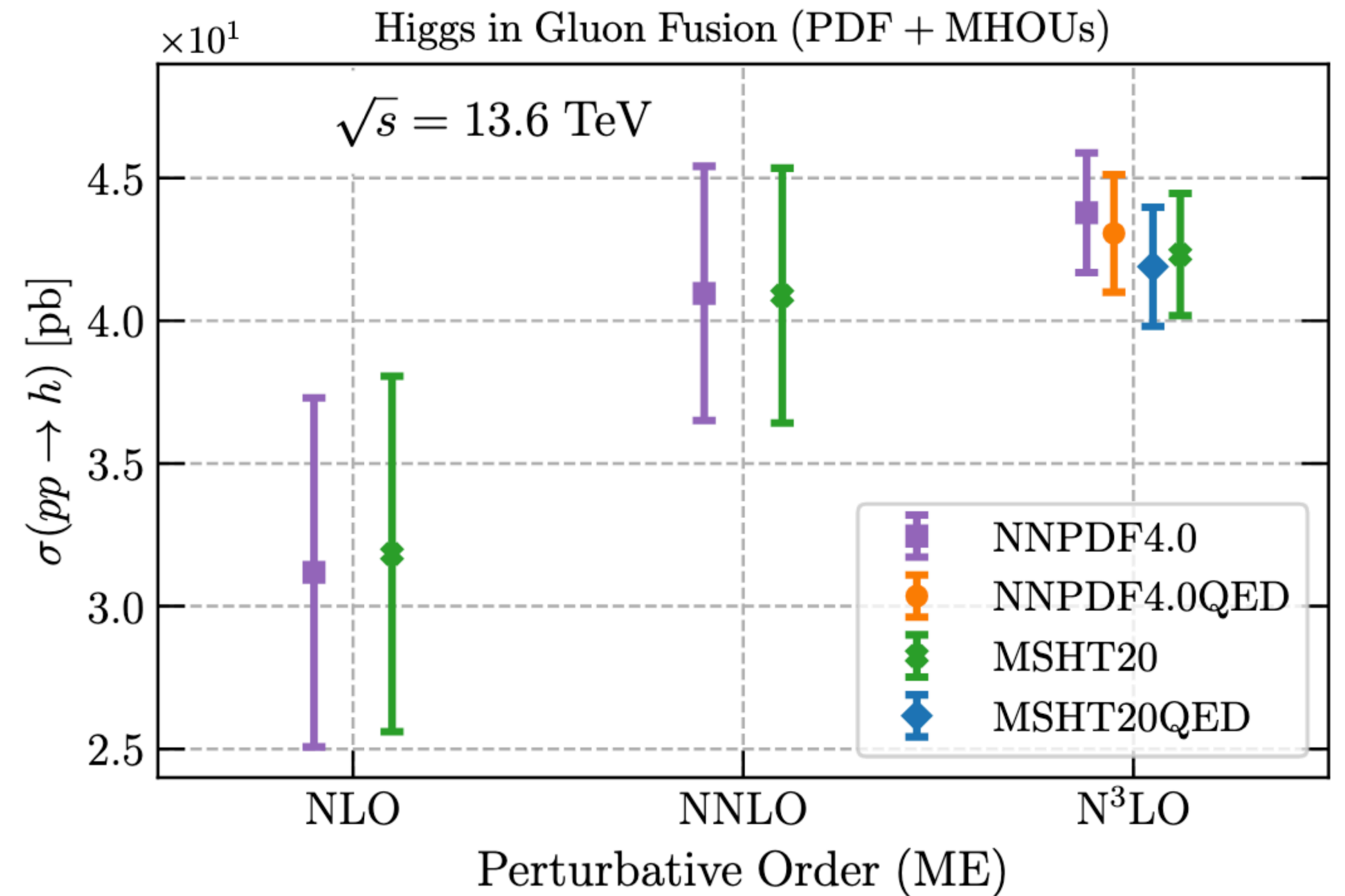
QED corrections &  $\gamma$ -PDF are key for LHC phenomenology:

- ❖ QED effects are of the same size as aN3LO
- ❖ Photon suppresses gluon momentum by up to 1%



Various LHC processes receive sizeable  $\gamma$ -initiated contributions:

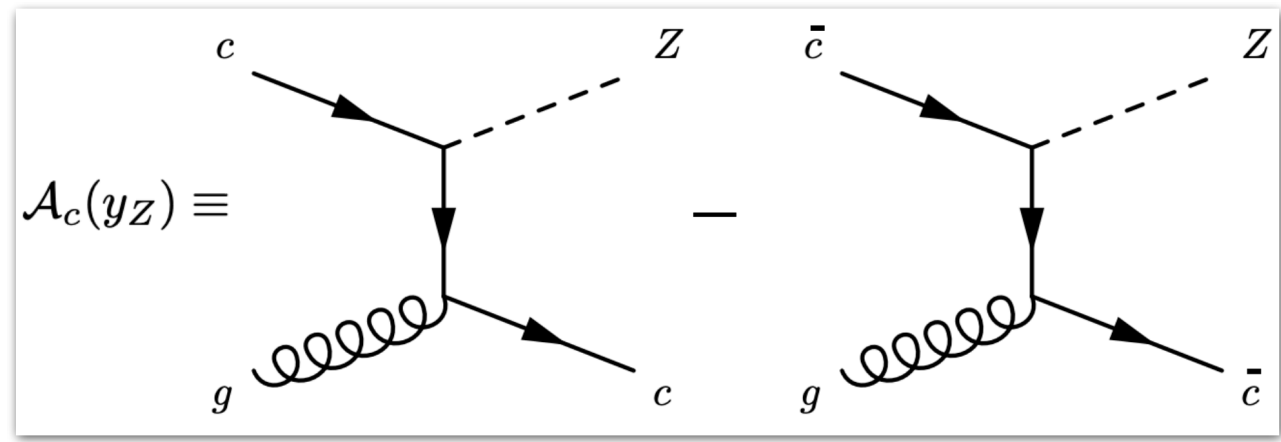
- ❖ aN3LO $\otimes$ QED result in a few percent suppression for ggH
- ❖ Difference between NNLO & N3LO predictions are reduced when using the appropriate PDFs at each order



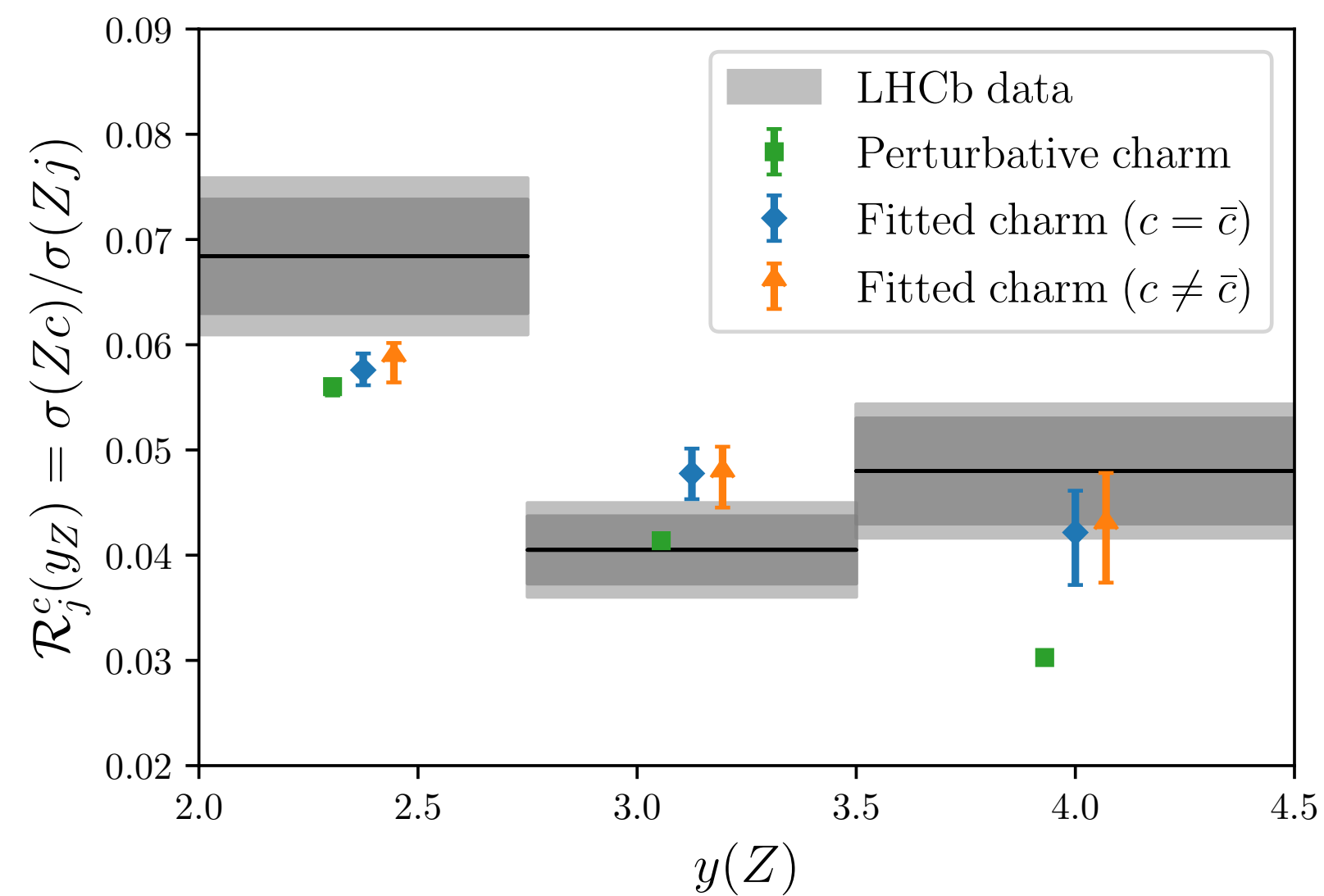
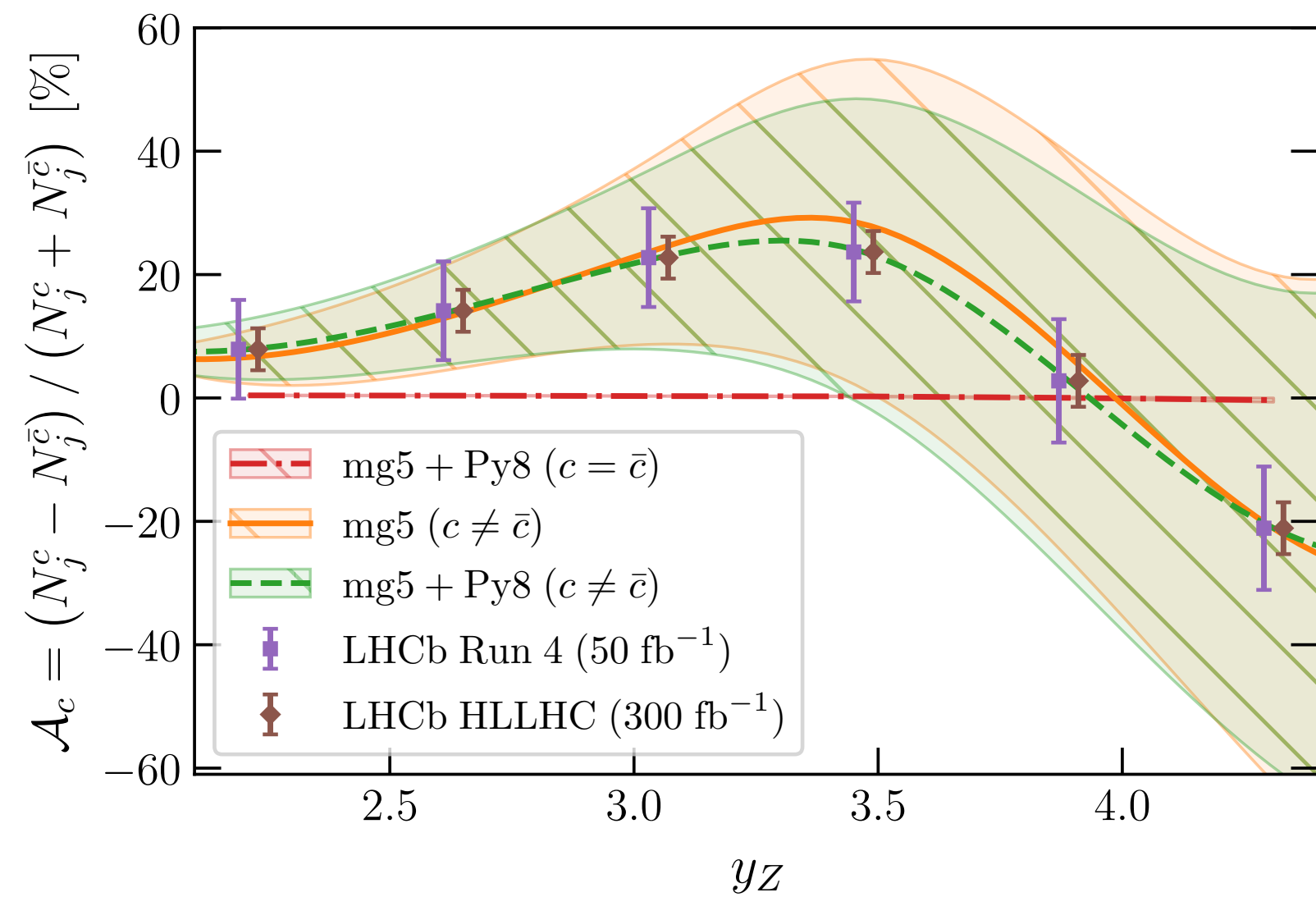
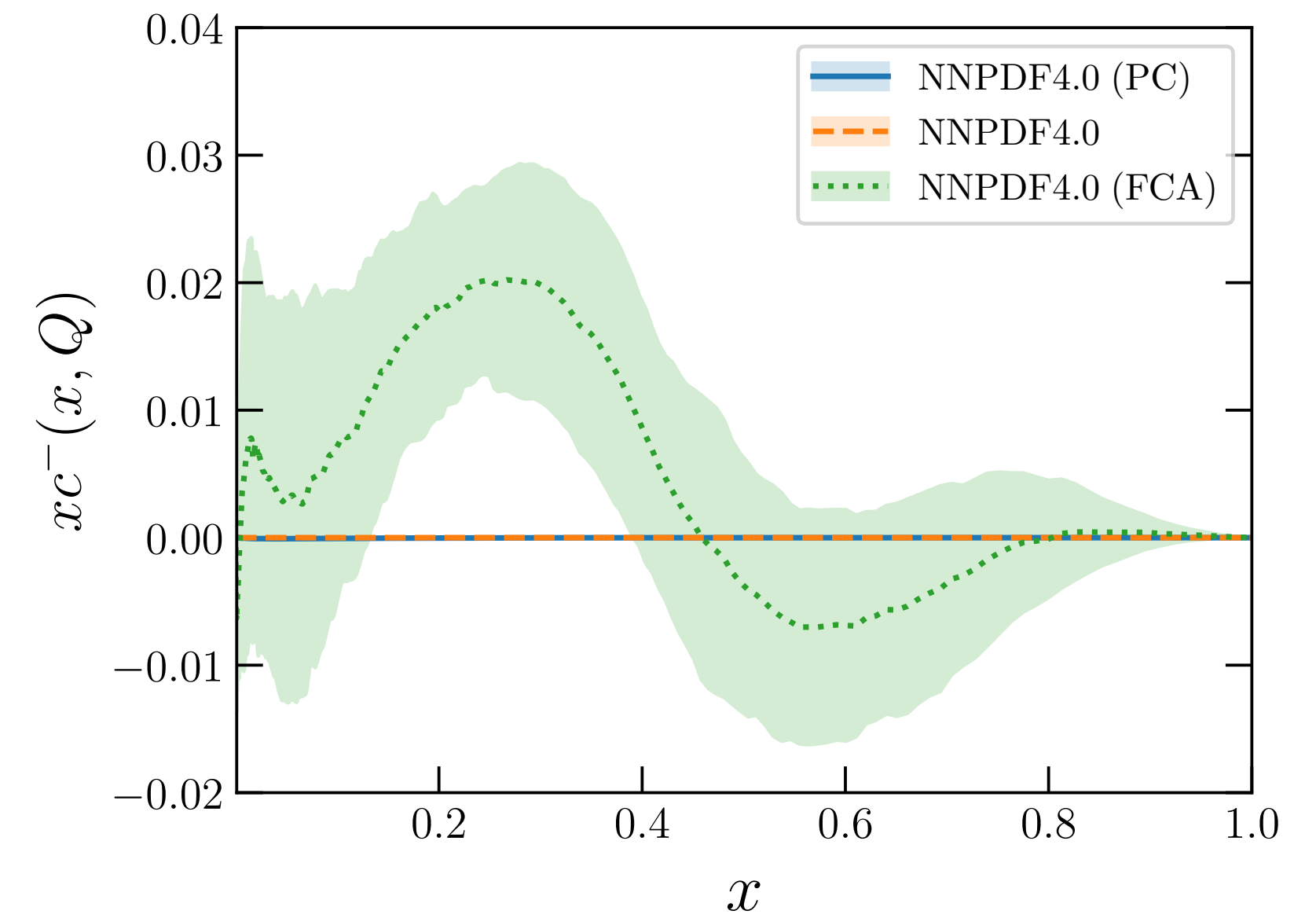
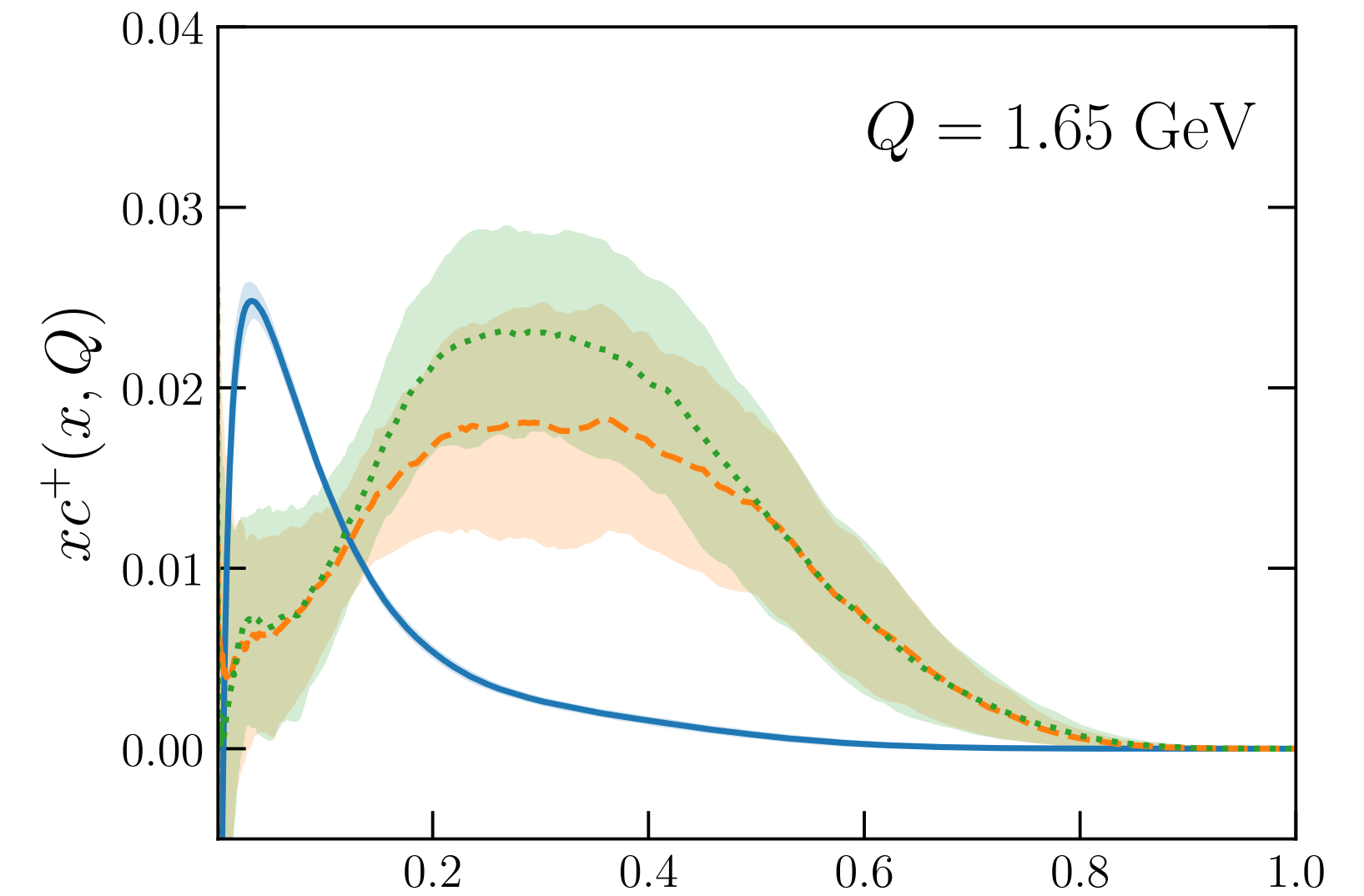
# Part II: Methodology

|                |  |                                       |
|----------------|--|---------------------------------------|
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# Separate determination of $c/\bar{c}$ PDFs



- ❖ There is no reason why intrinsic charm valence should vanish
- ❖ Projections for  $Z + c$ -jet at the LHCb favours non-vanishing valence charm
- ❖ Future LHC data will verify or falsify a non-zero charm valence.
- ❖ NNPDF4.1 will be based on a determination of a separate charm and anti-charm PDFs



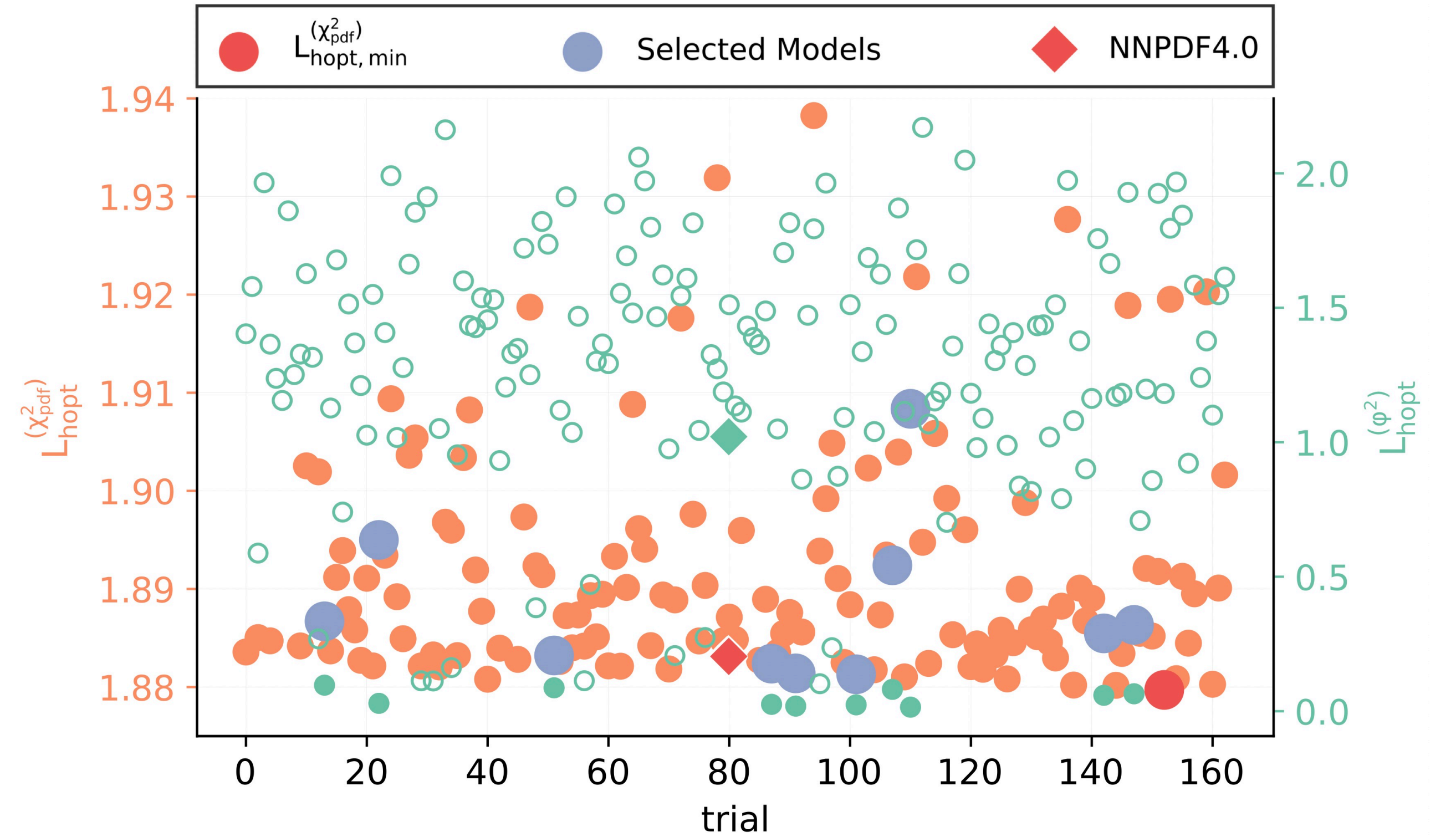
# Ensemble-based Hyperparameter Optimisation

- ❖ ML applications rely on a **large number of hyperparameters** with each combination defining a particular model
- ❖ Selection of the set of hyperparameters is crucial for a model to describe best the data & able to generalise  $\iff$  **Uncertainty quantifications**
- ❖ In NNPDF4.1, hyperparameter optimisation will be done **accounting for the full PDF distribution**
- ❖ The methodology still relies on the  $k$ -fold procedure but using a different  $L_{\text{hopt}}^{(\chi^2)}$  definition:

$$L_{\text{hopt}}^{(\chi_{\text{pdf}}^2)}(\hat{\theta}) = \frac{1}{n_{\text{fold}}} \sum_{p=1}^{n_{\text{fold}}} \min_{\theta \in \Theta} \left( \left\langle \chi_{\text{PDF},p}^2(\theta, \hat{\theta}) \right\rangle_{\text{rep}} \right)$$

- ❖ With an additional selection metric that maximises the generalisation power of the models:

$$\varphi_{\chi^2}^2 = \frac{1}{n_{\text{dat}}} \sum_{i,j=1}^{n_{\text{dat}}} (\text{cov})_{ij}^{-1} T_{ji}, \quad L_{\text{hopt}}^{(\varphi^2)}(\hat{\theta}) \equiv \left( \frac{1}{n_{\text{K}}} \sum_{p=1}^{n_{\text{K}}} \varphi_{\chi_p^2}^2(\hat{\theta}) \right)^{-1}$$



The outcome is **not a single Model**

Randomly sample over the complete population of acceptable hyperparameters displaying comparable performance

# Ensemble-based Hyperparameter Optimisation: Results

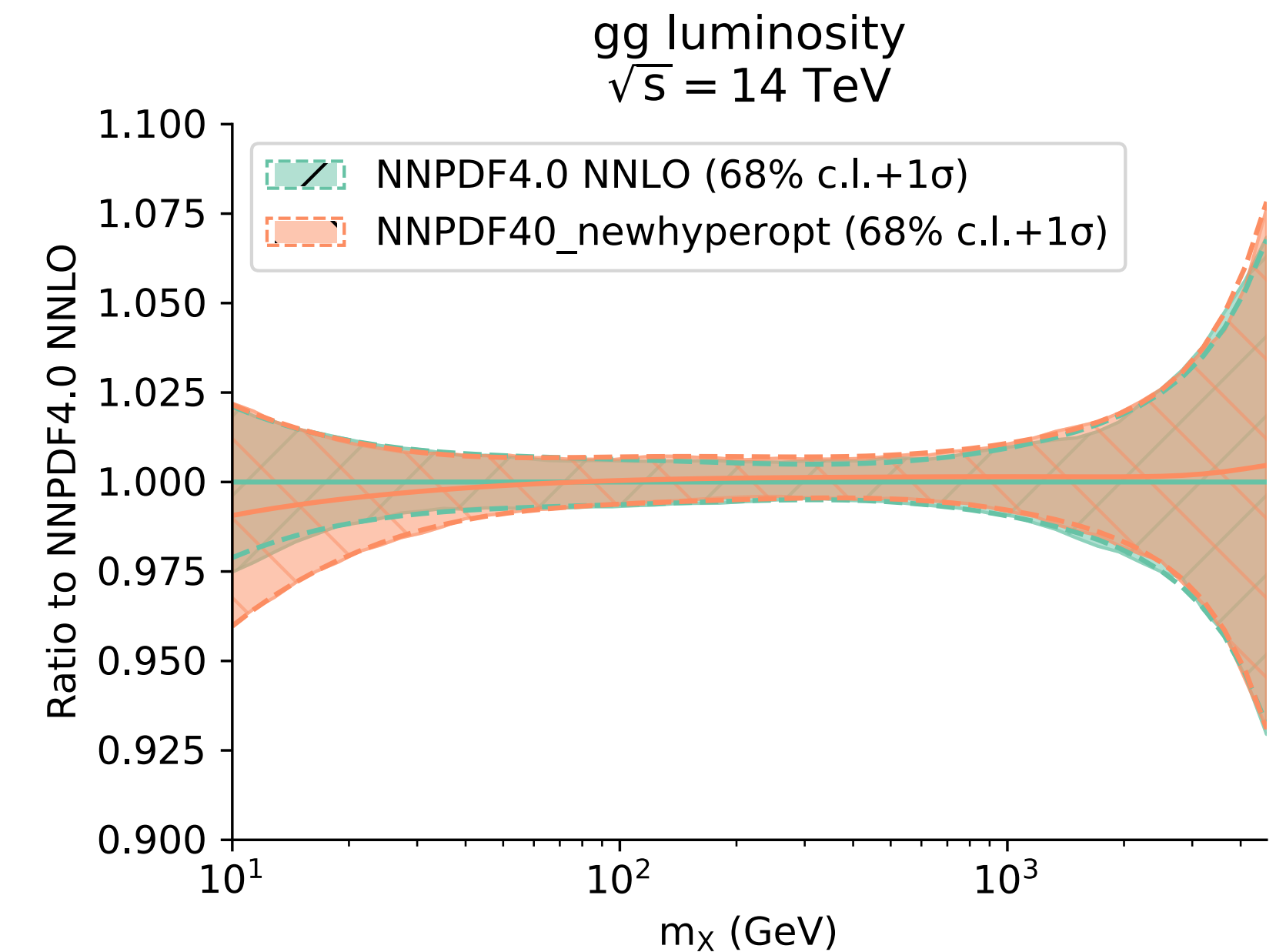
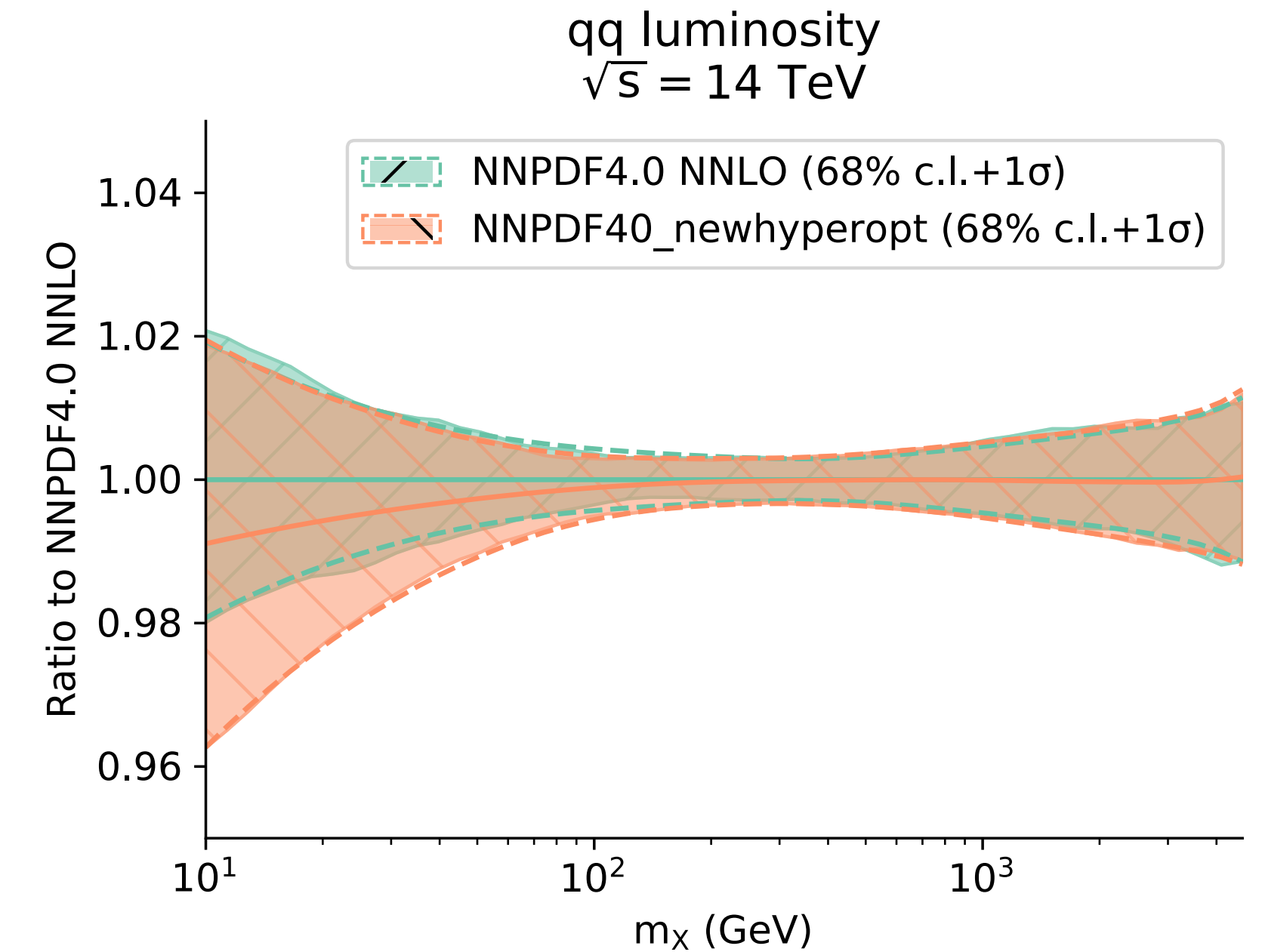
Ensemble-based hyperparameter optimisation is only possible with **hardware acceleration (GPUs)** and that provide various technical advantages:

| # Replicas       | 10   | 50  | 100 |
|------------------|------|-----|-----|
| Energy reduction | 78%  | 87% | 91% |
| Cost reduction   | -45% | 47% | 55% |

- ❖ Speed scales with the number of replicas (up to a factor of  $\sim 200$ )
- ❖ No significant increase of memory with the number of replicas
- ❖ Up to 90% energy reduction  $\iff$  More sustainable ML training

At the PDF level:

- ❖ **Excellent consistency with NNPDF4.0** with a moderate increase of the uncertainties in the extrapolation regions
- ❖ Non-trivial validation of the NNPDF methodology



# Part III: Data

|                |  |                                       |
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# How well do PDFs accommodate new data?

- ❖ Test PDF sets against new precise measurements from Run I/II using NNLO theories (w/o  $K$ -factors)
- ❖ Aim to assess how well PDF sets describe unseen data and whether these data will have effects on fits
- ❖ Agreement between data and theoretical predictions are quantified using the  $\chi^2$  definition:

$$\chi^2 = \frac{1}{n_{\text{dat}}} \sum_{i,j=1}^{n_{\text{dat}}} \left( T_i^{(0)} - D_i \right) (\text{cov}^{-1})_{ij} \left( T_j^{(0)} - D_j \right)$$

- ❖ Contributions to the covariance matrix include all sources of theoretical uncertainties:

$$\text{cov}_{ij} = \left( \text{cov}_{\text{exp}} \right)_{ij} + \left( \text{cov}_{\text{th}} \right)_{ij}$$

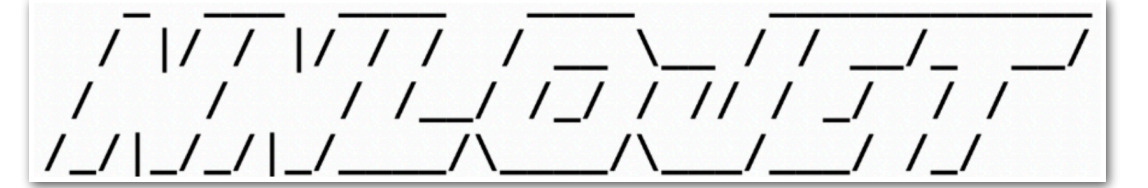
$$\left( \text{cov}_{\text{th}} \right)_{ij} = \left( \text{cov}_{\text{mho}} \right)_{ij} + \left( \text{cov}_{\text{pdf}} \right)_{ij} + \left( \text{cov}_{\text{as}} \right)_{ij}$$

$$\left( \text{cov}_{\text{pdf}}^{\text{HES}} \right)_{ij} = \sum_{k=1}^{n_{\text{eig}}} \left( T_i^{(k)} - T_i^{(0)} \right) \left( T_j^{(k)} - T_j^{(0)} \right)$$

$$\left( \text{cov}_{\text{pdf}}^{\text{MC}} \right)_{ij} = \frac{1}{n_{\text{rep}}} \sum_{k=1}^{n_{\text{rep}}} \left( T_i^{(k)} - \langle T_i \rangle_{\text{rep}} \right) \left( T_j^{(k)} - \langle T_j \rangle_{\text{rep}} \right)$$

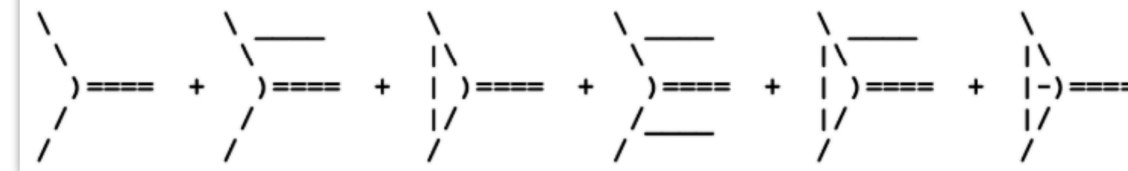
## PineAPPL

<https://github.com/NNPDF/pineappl>

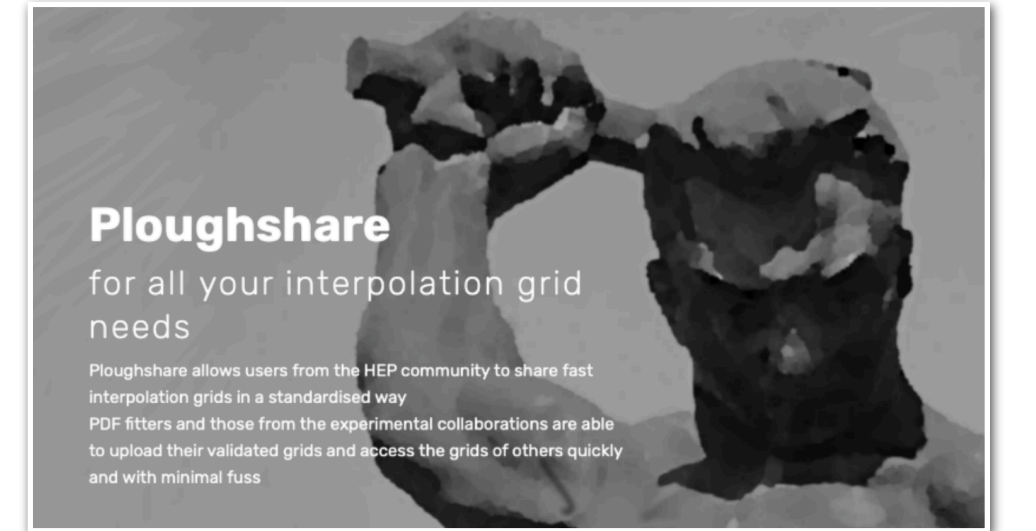


## MATRIX

Munich -- the Multi-channel Integrator at swiss (CH) precision --  
Automates qT-subtraction and Resummation to Integrate X-sections



<https://matrix.hepforge.org/>



<https://ploughshare.web.cern.ch/ploughshare/>

$$\left( \text{cov}_{\text{mho}} \right)_{ij} = \frac{1}{3} \left\{ \Delta_i^{+0} \Delta_j^{+0} + \Delta_i^{-0} \Delta_j^{-0} + \Delta_i^{0+} \Delta_j^{0+} + \Delta_i^{0-} \Delta_j^{0-} + \Delta_i^{++} \Delta_j^{++} + \Delta_i^{--} \Delta_j^{--} \right\}$$

$$\Delta_i(\kappa_R, \kappa_F) = T_i \left( \mu_R = \kappa_R \mu_R^{(0)}, \mu_F = \kappa_F \mu_F^{(0)} \right) - T_i \left( \mu_R^{(0)}, \mu_F^{(0)} \right)$$

$$\begin{aligned} \Delta_i^{+0} &= \Delta_i(2,1), & \Delta_i^{-0} &= \Delta_i(1/2,1), & \Delta_i^{0+} &= \Delta_i(1,1/2) \\ \Delta_i^{0-} &= \Delta_i(1,1/2), & \Delta_i^{++} &= \Delta_i(2,2), & \Delta_i^{--} &= \Delta_i(1/2,1/2) \end{aligned}$$

$$\left( \text{cov}_{\text{as}} \right)_{ij} = \frac{1}{2} \left\{ \Delta_{i,\alpha_s}^+ \Delta_{j,\alpha_s}^+ + \Delta_{i,\alpha_s}^- \Delta_{j,\alpha_s}^- \right\}$$

$$\Delta_{i,\alpha_s}^+ \equiv T_i(\alpha_s = 0.119) - T_i(\alpha_s = 0.118)$$

$$\Delta_{i,\alpha_s}^- \equiv T_i(\alpha_s = 0.118) - T_i(\alpha_s = 0.117)$$



# Experimental Data included in the study

| Process  | Experiment | Final State           | Observable  | $\sqrt{s}$ (TeV) | $\mathcal{L}$ (fb $^{-1}$ ) | $n_{\text{dat}}$ |
|----------|------------|-----------------------|---|------------------|-----------------------------|------------------|
| LHC W, Z | ATLAS      | Z $p_T$ spectrum      | $(\frac{1}{\sigma}) \frac{d\sigma}{dp_T^{\ell\bar{\ell}}}$          | 13               | 36.1                        | 38               |
|          | CMS        | W incl. prod.         | $\frac{d\sigma}{d \eta }$   | 13               | 35.9                        | 36               |
|          | LHCb       | Z incl. forward prod. | $\frac{d\sigma}{dy^Z}$  | 13               | 5.1                         | 17               |
|          | ATLAS      | Z incl. prod.         | $\frac{d\sigma}{d y }$  | 8                | 20.2                        | 7                |
| top-pair | ATLAS      | all-hadronic          | $(\frac{1}{\sigma}) \frac{d\sigma}{dm_{t\bar{t}}}$                  | 13               | 36.1                        | 9                |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{d y_{t\bar{t}} }$                | 13               | 36.1                        | 12               |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d^2\sigma}{d y_{t\bar{t}} dm_{t\bar{t}}}$ | 13               | 36.1                        | 11               |
|          | ATLAS      | $\ell$ +jets          | $(\frac{1}{\sigma}) \frac{d\sigma}{dm_{t\bar{t}}}$                  | 13               | 36.1                        | 9                |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{dp_T^{\ell}}$                    | 13               | 36.1                        | 8                |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{d y_{\ell} }$                    | 13               | 36.1                        | 5                |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{d y_{t\bar{t}} }$                | 13               | 36.1                        | 7                |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{dm_{t\bar{t}}}$                  | 13               | 137                         | 15               |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{dp_T^{\ell}}$                    | 13               | 137                         | 16               |
|          | CMS        | $\ell$ +jets          | $(\frac{1}{\sigma}) \frac{d\sigma}{d y_{t\bar{t}} }$                | 13               | 137                         | 10               |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d\sigma}{d y_{\ell} }$                    | 13               | 137                         | 11               |
|          |            |                       | $(\frac{1}{\sigma}) \frac{d^2\sigma}{d y_{t\bar{t}} dm_{t\bar{t}}}$ | 13               | 137                         | 35               |

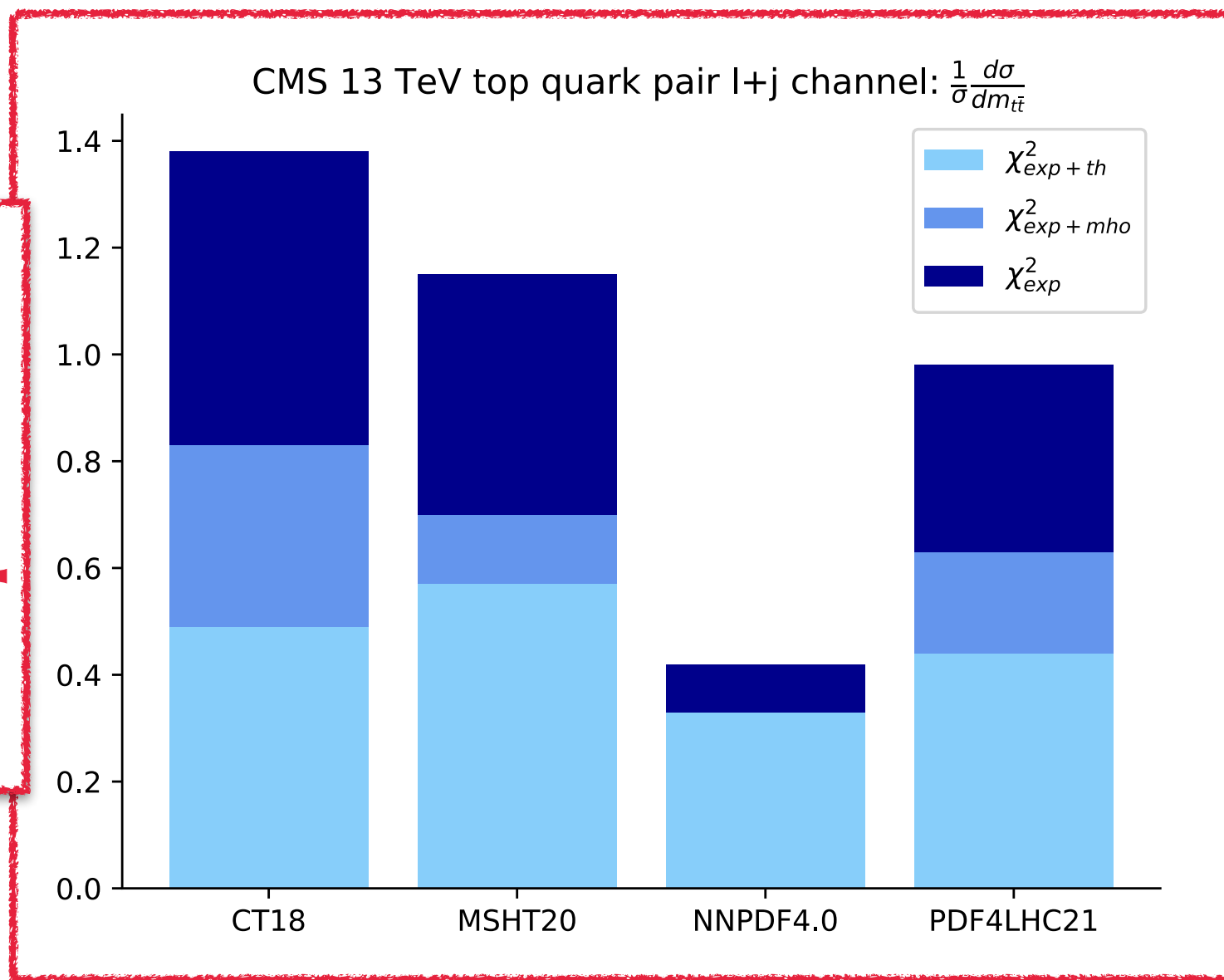
|           |       |                            |   |       |             |     |
|-----------|-------|----------------------------|---|-------|-------------|-----|
| LHC jets  | ATLAS | incl. jet $R = 0.6$        | $\frac{d^2\sigma}{dp_T d y }$                 | 13    | 3.2         | 177 |
|           | CMS   | incl. jets $R = 0.4$ (0.7) | $\frac{d^2\sigma}{dp_T d y }$                 | 13    | 36.3 (33.5) | 78  |
|           | ATLAS | di-jets $R = 0.6$          | $\frac{d^2\sigma}{dm_{jj} d y^* }$            | 13    | 3.2         | 136 |
| HERA jets | H1    | incl. jet (low $Q^2$ )     | $\frac{d^2\sigma}{dQ^2 dp_T}$                 | 0.319 | 0.29        | 48  |
|           | H1    | incl. jet (high $Q^2$ )    | $\frac{d^2\sigma}{dQ^2 dp_T}$                 | 0.319 | 0.351       | 24  |
|           | ZEUS  | incl. jet                  | $\frac{d^2\sigma}{dQ^2 dE_T}$                 | 0.300 | 0.038       | 30  |
|           | ZEUS  | incl. jet                  | $\frac{d^2\sigma}{dQ^2 dE_T}$                 | 0.319 | 0.082       | 30  |
|           | H1    | di-jets (low $Q^2$ )       | $\frac{d^2\sigma}{dQ^2 d\langle p_T \rangle}$ | 0.319 | 0.29        | 48  |
|           | H1    | di-jets (high $Q^2$ )      | $\frac{d^2\sigma}{dQ^2 d\langle p_T \rangle}$ | 0.319 | 0.351       | 24  |
|           | ZEUS  | di-jets                    | $\frac{d^2\sigma}{dQ^2 d\langle E_T \rangle}$ | 0.319 | 0.374       | 22  |

## Dataset selection criteria:

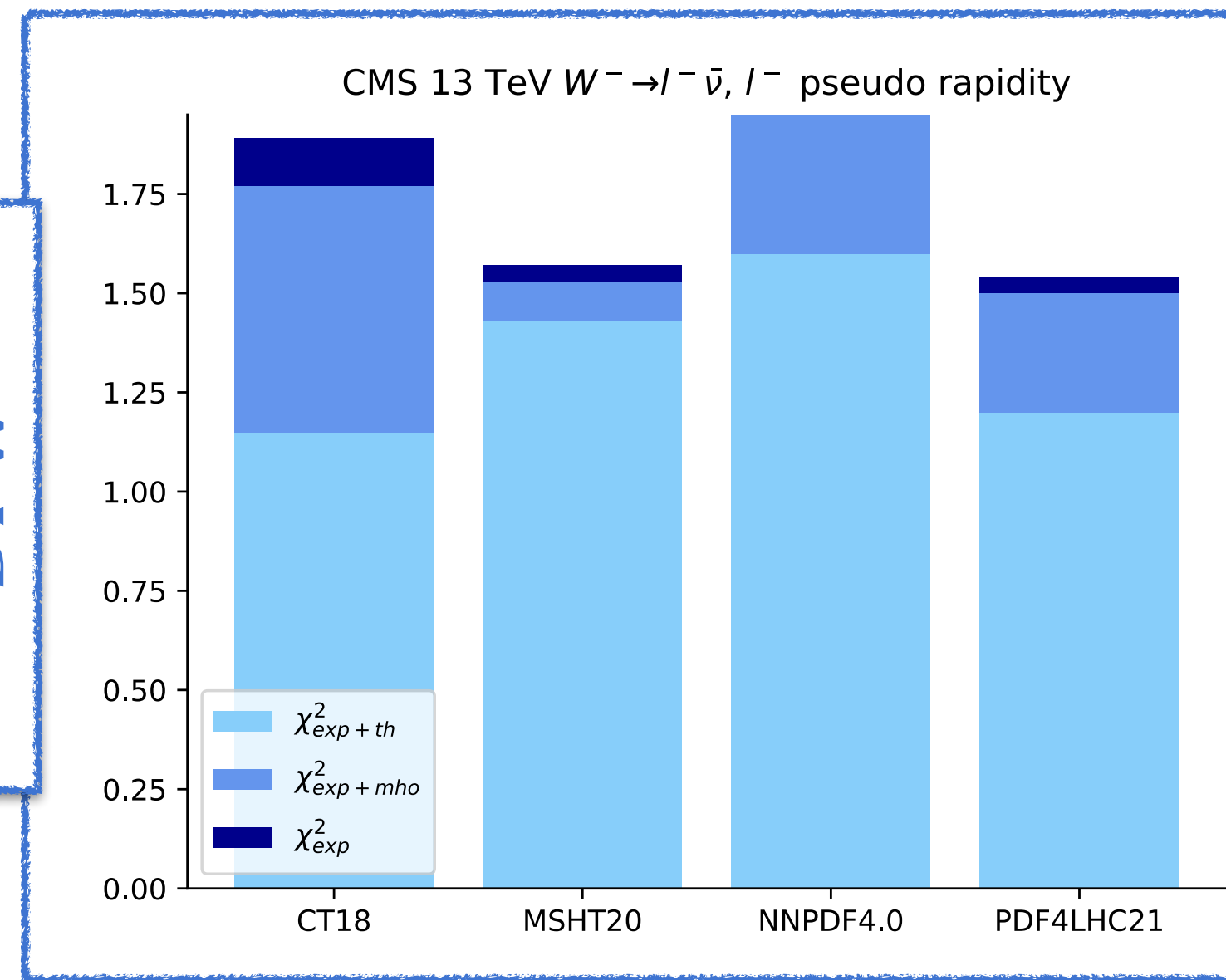
- ❖ Not included in NNPDF4.0 except for ATLAS Z @ 8 TeV
- ❖ Publicly available on HepData
- ❖ Provide info on PDFs of # partons & computable @ NNLO interfaced to PineAPPL fast interpolation grids

# $\chi^2$ comparisons for various PDFs

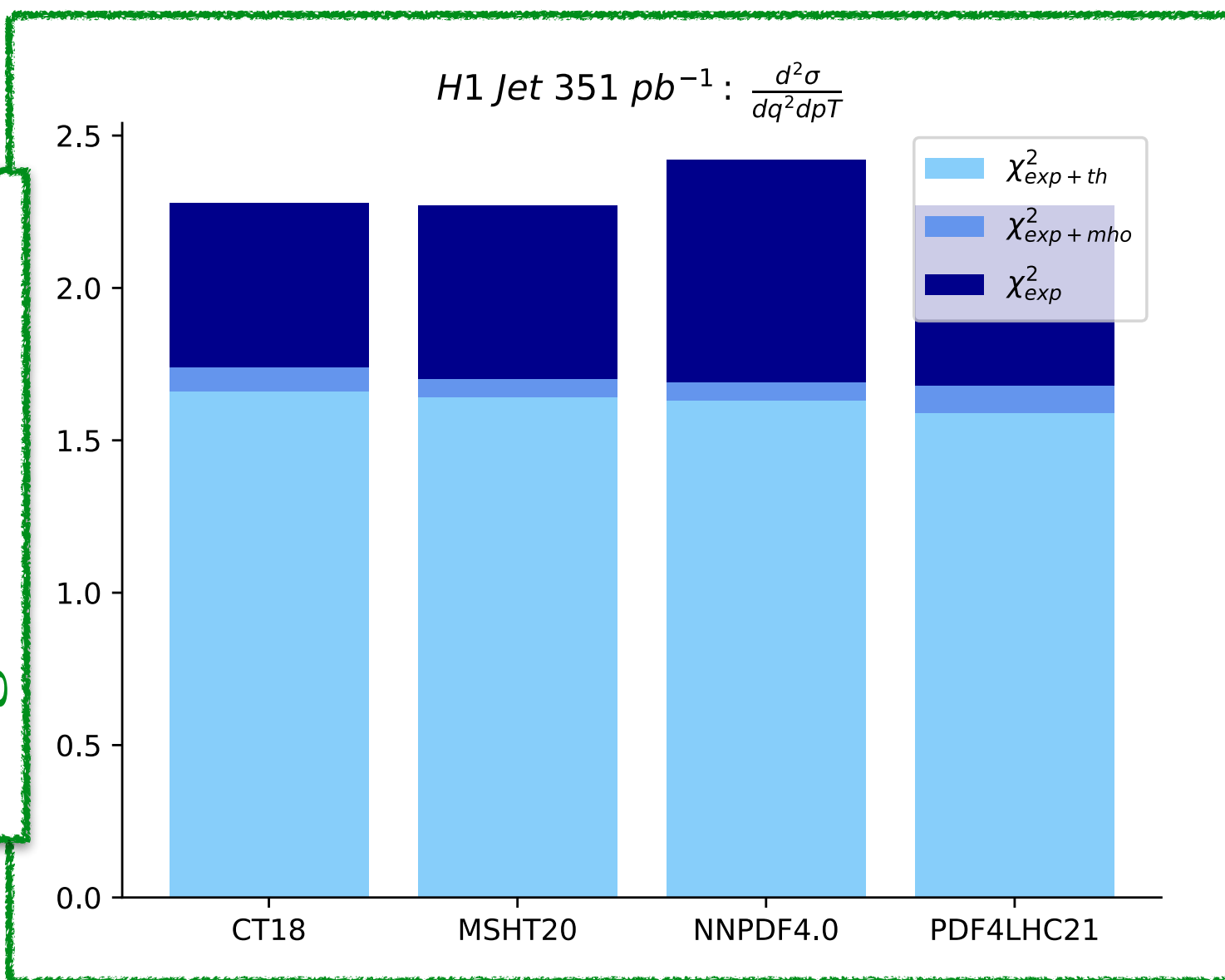
Top Sector



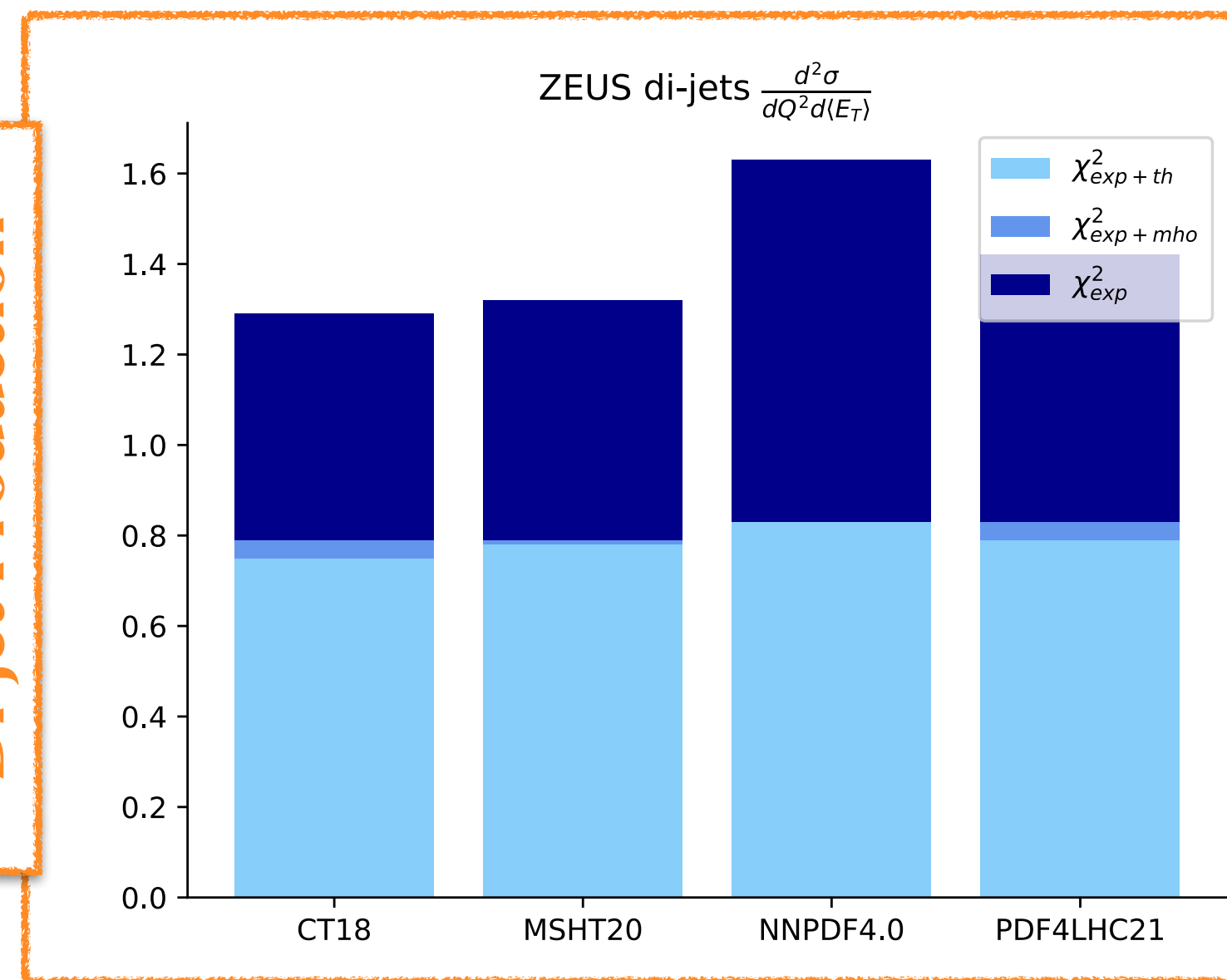
DY W



Single Inclusive Jet

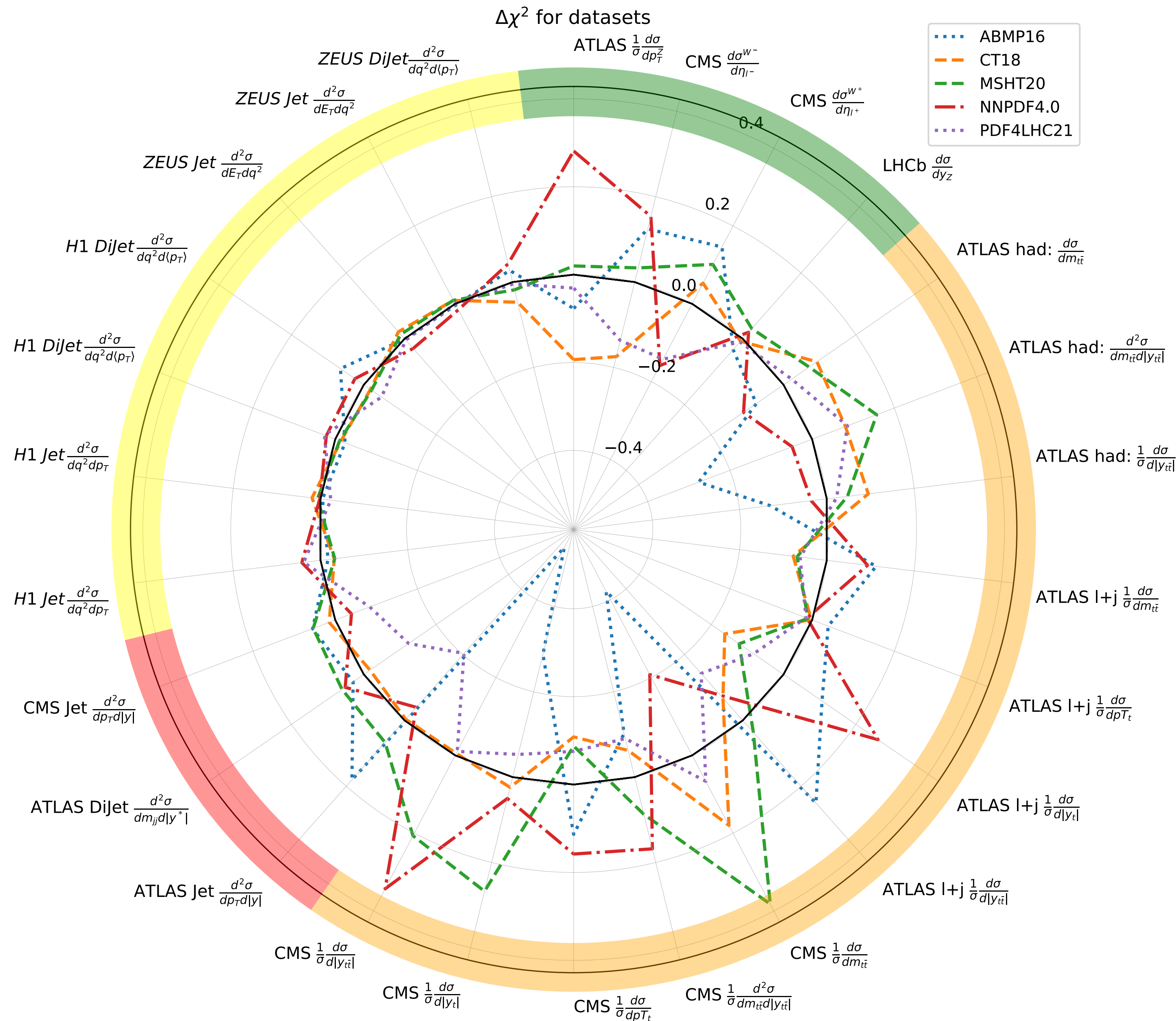


Di-jet Production



- ❖ All PDF sets have similar predictive power
- ❖ The inclusion of MHO uncertainties can have significant impact on the description of data
- ❖ By taking into account all possible sources of uncertainties (MHO, PDF,  $\alpha_s$ ), the differences at the cross section level dissipate
- ❖ LHC measurements do not strongly discriminate among PDF

# $\Delta\chi^2$ Results



- ❖ Relative change in the total  $\chi^2$  due to a change in the input PDF

$$\Delta\chi^{2(i)} = \frac{\chi_{\text{exp+th}}^{2(i)} - \langle \chi_{\text{exp+th}}^2 \rangle_{\text{pdfs}}}{\langle \chi_{\text{exp+th}}^2 \rangle_{\text{pdfs}}}$$

where

$$\langle \chi_{\text{exp+th}}^2 \rangle_{\text{pdfs}} = \frac{1}{n_{\text{pdfs}}} \sum_{i=1}^{n_{\text{pdfs}}} \chi_{\text{exp+th}}^{2(i)}$$

- ❖ **No systematic outlier** seen in the data description despite noticeable differences at the level of PDF
- ❖ As anticipated, PDF4LHC21 represents the **average** (with  $\Delta\chi^2 \sim 0$ )

# Conclusions & Outlook

- ❖ The precision era at the LHC requires precise & accurate PDFs and advancements are needed for the “three pillars”: theory, methodology, and data
- ❖ Significant progress in the NNPDF global analysis for a state-of-the-art determination at  $\alpha\text{N}^3\text{LO} \otimes \text{QED} \otimes \text{MHOU}$ s
- ❖ Significant improvements in the Hyperparameter Optimisation using ensemble-based methodology  $\iff$  **demonstrates the robustness of the NNPDF uncertainty estimate**
- ❖ A quantitative appraisal of PDF fits using precision LHC measurements show that **all PDF sets have similar predictive power** despite significant differences at the PDF level
- ❖ NNPDF4.1 will be based on the best theory (at the very least with pure NNLO hadronic predictions) and fitting methodology, as well as more precision LHC measurements

**THANKS FOR YOUR ATTENTION**



“Wanderer above the Sea of Fog” by Caspar David Friedrich