### Collinear PDFs for $M_W$ determination

 $\mathsf{GDR}\ \mathsf{QCD}\ \mathsf{workshop}\ \mathsf{on}\ W\ \mathsf{mass}$ 

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## PDFs and recent $M_W$ determinations



Adapted from arXiv:2412.13872

PDF uncertainty non-negligible on the final  $M_W$  uncertainty e.g. on the CMS result  $\delta_{tot} = 9.9$  MeV and  $\delta_{PDF} = 4.4$  MeV

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#### Making predictions with PDFs CT18 4000 • MSHT20 NNPDF3.1 3900 ABMP16 ATLASpdf21 3800F ♦ PDF4LHC15 $\sigma_{W^-}$ [pb] ✿ PDF4LHC21 NNPDF4.0 37003600 3500LHC 14 TeV, $2\sigma$ 4600 500048005200 $\sigma_{W^+}$ [pb]

Acta Phys.Polon.B 53 (2022) 12

Predictions obtained with different PDF sets display significantly different uncertainties Where do these come from (data, theory, methodology)?

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## 1. Data

## Available cross setions and their kinematic coverage





Increasing relevance of LHC data in PDF determination

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### Quark flavour separation





Perturbative charm alters the flavour decomposition and deteriorates the fit

 $\chi^2_{
m fitted\,charm} = 1.17 
ightarrow \chi^2_{
m pert.\,charm} = 1.19$  mainly due to a worsening

of the LHC W, Z and top pair data sets

fitting charm reduces the dependence from  $m_c$ 

[EPJ C76 (2016) 647; C77 (2017) 663; C82 (2022) 428]

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## Quark flavour separation



100

## The PDF4LHC21 benchmark [JPG 49 (2022) 080501]



Same reduced data set

Same heavy quark mass values:  $m_c = 1.4 \text{ GeV}$ ;  $m_b = 4.75 \text{ GeV}$ ;  $\alpha_s(M_Z) = 0.118$ 

Same strong coupling value:  $\alpha_s(M_Z) = 0.118$ 

No strangeness asymmetry at input scale:  $(s - \bar{s})(Q_0) = 0$ 

Charm perturbatively generated

Positive-definite quark distributions

No deuteron or nuclear corrections to analyse  $\nu N$  cross sections

Fixed branching ratio for charm hadrons to muons; NNLO corrections for dimuon data

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## The PDF4LHC21 benchmark [JPG 49 (2022) 080501]



Very good agreement within uncertainties; similar size of uncertainties in the data region Remaining differences reflect methodological choices

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## How well do different PDF sets generalise on unseen data?

Example: the CMS W rapidity measurement [PRD 102 (2020) 092012] The covariance matrix include exp, PDF, MHOU, and  $\alpha_s$  uncertainties [arXiv:2501.10359]

![](_page_9_Figure_2.jpeg)

![](_page_9_Figure_3.jpeg)

![](_page_9_Figure_4.jpeg)

# 2. Theory

## QED corrections in PDF determination: LuxQED photon

#### LUXQED [PRL 117 (2016) 242002]

View the  $ep \rightarrow e + X$  process as an electron scattering off the photon field of the proton

Consider a BSM process, e.g. production of a heavy supersymmetric lepton L in ep collision, write the cross section in terms of structure functions and of  $f_{\gamma}$ , and equate the two to obtain  $f_{\gamma}$ 

![](_page_11_Figure_4.jpeg)

$$\begin{split} \sigma &= c_0 \sum_a \int_x^1 \frac{dz}{z} \, \hat{\sigma}_a(z,\mu^2) \frac{M^2}{zs} f_{a/p} \left( \frac{M^2}{zs}, \mu^2 \right) & x f_{\gamma/p}(x,\mu^2) = \\ \sigma &= \frac{c_0}{2\pi} \int_x^{1-\frac{2\pi m_p}{2}} \frac{dz}{z} \int_{Q_{\min}^2}^{Q_{\max}^2} \frac{dQ^2}{Q^2} \alpha_{ph}^2(-Q^2) \left[ \left( 2-2z+z^2 & \frac{1}{2\pi\alpha(\mu^2)} \int_x^1 \frac{dz}{z} \left\{ \int_x^{\frac{\mu^2}{1-z}} \frac{dQ^2}{Q^2} \alpha^2(Q^2) + \frac{2x^2 m_p^2}{Q^2} + \frac{z^2 Q^2}{Q^2} - \frac{2x^2 Q^2 m_p^2}{M^4} \right\} F_2(x/z,Q^2) & \left[ \left( z p_{\gamma q}(z) + \frac{2x^2 m_p^2}{Q^2} \right) F_2(x/z,Q^2) - z^2 F_L\left(\frac{x}{z},Q^2\right) \right] \\ &+ \left( -z^2 - \frac{z^2 Q^2}{2M^2} + \frac{z^2 Q^4}{2M^4} \right) F_L(x/z,Q^2) \right] & - \alpha^2(\mu^2) z^2 F_2\left(\frac{x}{z},\mu^2\right) \end{split}$$

Iterate a QCD fit including  $f_{\gamma}$  in DGLAP and in the momentum sum rule

## QED corrections in PDF determination: LuxQED photon

![](_page_12_Figure_1.jpeg)

General agreement in  $\mathcal{L}_{\gamma\gamma}$  and  $\mathcal{L}_{\gamma g}$ ; differences in  $\mathcal{L}_{q\bar{q}}$  independent of photon PDF

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## Photon PDF: impact on $W^{\pm}$ differential cross sections

![](_page_13_Figure_1.jpeg)

PDF4LHC benchmark and combination ongoing

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## N<sup>3</sup>LO QCD corrections in PDF determination

NNLO is the precision frontier for PDF determination

N3LO is the precision frontier for partonic cross sections

Mismatch between perturbative order of partonic cross sections and accuracy of PDFs is becoming a significant source of uncertainty

$$\hat{\sigma} = \alpha_s^p \hat{\sigma}_0 + \alpha_s^{p+1} \hat{\sigma}_1 + \alpha_s^{p+2} \hat{\sigma}_2 + \mathcal{O}(\alpha_s^{p+3}) \qquad \delta(\text{PDF} - \text{TH}) = \frac{1}{2} \left| \frac{\sigma_{\text{NNLO-PDFs}}^{(2)} - \sigma_{\text{NLO-PDFs}}^{(2)}}{\sigma_{\text{NNLO-PDFs}}^{(2)}} \right|$$

![](_page_14_Figure_5.jpeg)

[JHEP 11 (2020) 143]

## N<sup>3</sup>LO QCD corrections in PDF determination

#### Splitting Functions

- Singlet  $(P_{qq}, P_{gg}, P_{gq}, P_{qg})$
- large- $n_f$  limit [NPB 915 (2017) 335; arXiv:2308.07958]
- small-x limit [JHEP 06 (2018) 145]
- large-x limit [NPB 832 (2010) 152; JHEP 04 (2020) 018; JHEP 09 (2022) 155]
- -5 (10) lowest Mellin moments [PLB 825 (2022) 136853; ibid. 842 (2023) 137944; ibid. 846 (2023) 138215]
- Non-singlet ( $P_{NS,v}$ ,  $P_{NS,+}$ ,  $P_{NS,-}$ )
- large- $n_f$  limit [NPB 915 (2017) 335; arXiv:2308.07958]
- small-x limit [JHEP 08 (2022) 135]
- large-x limit [JHEP 10 (2017) 041]
- 8 lowest Mellin moments [JHEP 06 (2018) 073]

#### DIS structure functions ( $F_L$ , $F_2$ , $F_3$ )

- DIS NC (massless) [NPB 492 (1997) 338; PLB 606 (2005) 123; NPB 724 (2005) 3]
- DIS CC (massless) [Nucl.Phys.B 813 (2009) 220]
- massive from parametrisation combining known limits and damping functions [NPB 864 (2012) 399]

PDF matching conditions

- all known except for  $a_{H,a}^3$  [NPB 820 (2009) 417; NPB 886 (2014) 733; JHEP 12 (2022) 134]

#### Coefficient functions for other processes

- DY (inclusive) [JHEP11 (2020) 143]; DY (y differential) [PRL 128 (2022) 052001]

### Theory uncertainties in PDF determination

Assuming that theory uncertainties are (a) Gaussian and (b) independent from experimental uncertainties, modify the figure of merit to account for theory errors

$$\chi^2 = \sum_{i,j}^{N_{\text{dat}}} (D_i - T_i) (\text{cov}_{\text{exp}} + \text{cov}_{\text{th}})_{ij}^{-1} (D_j - T_j); \ (\text{cov}_{\text{th}})_{ij} = \frac{1}{N} \sum_k^N \Delta_i^{(k)} \Delta_j^{(k)}; \ \Delta_i^{(k)} \equiv T_i^{(k)} - T_i$$

Problem reduced to estimate the th. cov. matrix, e.g. in terms of nuisance parameters

$$\Delta_i^{(k)} = T_i(\mu_R, \mu_F) - T_i(\mu_{R,0}, \mu_{F,0});$$
 vary scales in  $\frac{1}{2} \le \frac{\mu_F}{\mu_{F,0}}, \frac{\mu_R}{\mu_{R,0}} \le 2$ 

![](_page_16_Figure_5.jpeg)

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![](_page_17_Figure_5.jpeg)

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![](_page_18_Figure_4.jpeg)

## Impact of aN<sup>3</sup>LO corections on partonic luminositites

![](_page_19_Figure_1.jpeg)

Trend observed for NNPDF4.0 [EPJ C84 (2024) 659] similar to MSHT20 [EPJ C83 (2023) 185]

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## aN<sup>3</sup>LO PDFs: MSHT20 vs NNPDF4.0

![](_page_20_Figure_1.jpeg)

A 3% shift for  $M_X \sim 100$  GeV, already at NNLO, partly sensitive to higher moments

## aN<sup>3</sup>LO PDFs: impact on $W^{\pm}$ total cross sections

![](_page_21_Figure_1.jpeg)

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## Combining QED and QCD higher order corrections

![](_page_22_Figure_1.jpeg)

QED corrections slightly reduce discrepancies between the MSHT20 and NNPDF4.0

## Combining MSHT20 and NNPDF4.0 aN<sup>3</sup>LO QED PDFs

![](_page_23_Figure_1.jpeg)

Statistical combination performed as in PDF4LHC21 [JPG 52 (2025) 065002]

## Combining MSHT20 and NNPDF4.0 NNLO QED PDFs

![](_page_24_Figure_1.jpeg)

Percent impact of aN<sup>3</sup>LO corrections; comparatively small impact of QED corrections

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# 3. Methodology

### Validate PDF uncertainties: closure tests

Fit PDFs to pseudodata generated assuming a known underlying law

#### **NNPDF**

bias difference of central prediction and truth variance uncertainty of replica predictions

If model complexity optimal, generalisation error minimum  $(E[bias])^2 = variance$ 25 fits, 40 replicas each

#### <u>MSHT</u>

parameterisation flexible enough to give faithful description of global pseudodata

NNPDF *vs* MSHT comparison similar fit quality similar errors, with MSHT unit tolerance

![](_page_26_Figure_8.jpeg)

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### Inconsistent closure tests

Generate pseudodata with statistical and systematic uncertainties

 $C = C^{\rm stat} + C^{\rm syst} \qquad C^{\rm syst}_{ij} = \sum_k \Delta^k_i \Delta^k_j \qquad \Delta^k_i \text{ is the k-th sys. unc. for the i-th point}$ 

Assume systematic uncertainties are underestimated and perform a fit with

 $\Delta^k_i \to \lambda \Delta^k_i \qquad \lambda = 1 \text{ consistency} \qquad \lambda = 0 \text{ extreme inconsistency}$ 

#### PREDICTED UNCERTAINTY ON GENERATED DATA

CONSISTENT

EXTREME INCONSISTENCY

![](_page_27_Figure_8.jpeg)

![](_page_27_Figure_9.jpeg)

arXiv:2503.17447

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![](_page_28_Figure_0.jpeg)

The ML model corrects for inconsistency except in extreme cases [arXiv:2503.17447]

## Simultaneous determination of PDFs and SM parameters

Example:  $\alpha_s$  and PDFs are correlated, hence looking at the  $\chi^2$  profile as a function of  $\alpha_s$  with fixed PDFs results in a bias [EPJC80 (2020) 182]

Simultaneous fits of SM parameters and PDFs are costly

Correlated replica method [EPJ C78 (2018) 408]

Perform multiple fits of the same data replica, changing the value of  $\alpha_s$  in each fit, thereby correlating PDFs and  $\alpha_s$ 

<u>Theory covariance method</u> [EPJ C81 (2021)830] Perform a single fit with a theory covariance matrix encoding a prior  $\alpha_s$  distribution, and determine  $\alpha_s$  from the Bayesian posterior

![](_page_29_Figure_6.jpeg)

Closure test passed by both methodologies. Apply them to real data to get  $\alpha_s(M_Z)$  $\alpha_s(M_Z) = 0.1194^{+0.0007}_{-0.0014}$  at aN<sup>3</sup>LO<sub>QCD</sub> $\otimes$ NLO<sub>QED</sub> [arXiv:2506.13871]

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# 4. Conclusions

## Summary

A precise and accurate determination of PDFs is key to do precision phenomenology. LHC measurements are being instrumental to reduce PDF uncertainties to few percent. This is not enough.

The goal of achieving PDF determinations accurate to 1% opens up some challenges. Understand the interplay between data, theory, and methodology into PDF uncertainties.

Refine the theoretical accuracy of a PDF determination.

Represent theory uncertainties into PDF uncertainties.

Deploy a robust fitting methodology and good statistical tests of it.

Benchmark efforts may benefit from public releases of PDF codes and inputs.

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## Thank you